

Determination of the expansion rate from Type Ia Supernovae: a model-independent approach

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Type Ia Supernovae

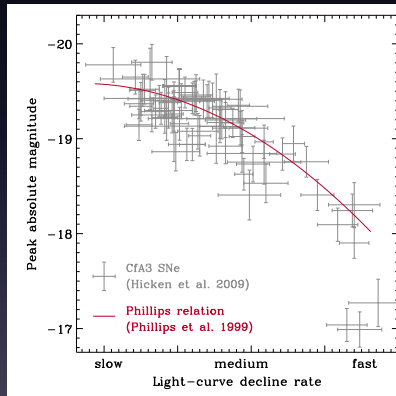
- ▶ Standardizable candles. Can be calibrated through the **width-luminosity relation** (*Phillips 1993*)

- ▶ **Brighter** events decline **slower**, fainter decay faster

- ▶ Distance modulus

$$\mu_B = m_B + \alpha X_1 - \beta c - M_B$$

$$\text{Luminosity Distance } D_L(z) = 10^{\mu/5+1}$$



(*S. Taubenberger*)

"Model-independent"?

- **Usual Approach:** Cosmological parameters are obtained from fits to observations and model predictions are validated

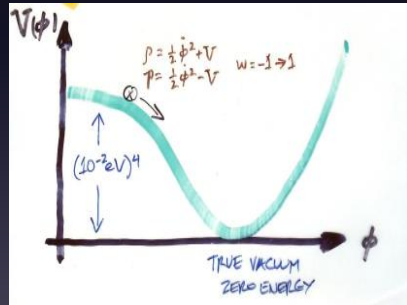
$$H^2(a) \equiv H_0^2 E^2(a) = H_0^2 \left[\frac{\Omega_{r0}}{a^4} + \frac{\Omega_{m0}}{a^3} - \frac{\Omega_{k0}}{a^2} + \Omega_{\Lambda 0} \right]$$

- SNe Ia directly test geometry. No need to define a Friedman Equation a priori
- **Direct reconstruction** of D_L and H
- Literature: *Starobinsky 1998; Huterer & Turner 1999, 2001; Shafieloo 2006, 2007, Ishida et al. 2011...*

Motivation

- **Geometrical framework** to compare different cosmological theories
- **vacuum energy**: $w = -1$
- dynamically evolving **scalar-field**:
 $w(a) = w_0 + w_a(1 - a)$
- **Modified Gravity**: terms on the left side of Einstein's equations

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^4}T_{\mu\nu}$$



The Method

From **Robertson-Walker metric** [$a = (1 + z)^{-1}$ and $e(a) = 1/E(a)$]

$$D_L(a) = \frac{c}{H_0} \frac{1}{a} \int_a^1 \frac{e(x)}{x^2} dx \xrightarrow{d/da} D'_L(a) = -\frac{1}{a^2} \int_a^1 \frac{e(x)}{x^2} dx - \frac{e(a)}{a^3}$$

Volterra Integral Equation

$$e(a) = f(a) + \lambda \int_1^x K(x, t) e(t) dt$$

$$e(a) = -a^3 D'_L + a \int_1^a x^{-2} e(x) dx$$

Neumann Series: (Arfken & Weber 1995)

$$e(a) = e_0 + \sum_{n=1}^{\infty} \lambda^n e_n(a)$$

$$e_0 = -a^3 D'_L(a)$$

$$e_n = \int_1^a x^{-2} e_{n-1}(t) dt$$

The Method

Problem: equation involves D_L' !

- First: **smooth the data**

$$D_L(a) = \sum_{j=0}^J c_j p_j(a) \rightarrow D_L'(a) = \sum_{j=0}^J c_j p_j'(a)$$

- Approximate this derivative to the derivative of the real data
- **Principal Component Analysis** to derive an *optimal* basis
- The expansion coefficients are found minimizing

$$\chi^2 = \sum_i^{\text{SNe}} \frac{[D_L|_{\text{obs},i} - D_L(a_i, c_j, H_0)]^2}{\sigma_i^2}$$

PC Approach

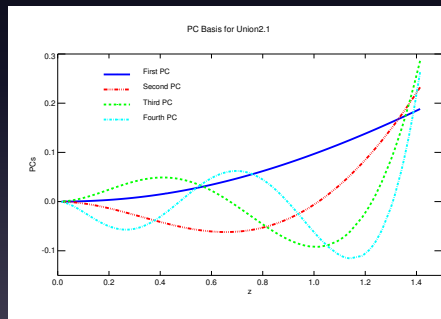
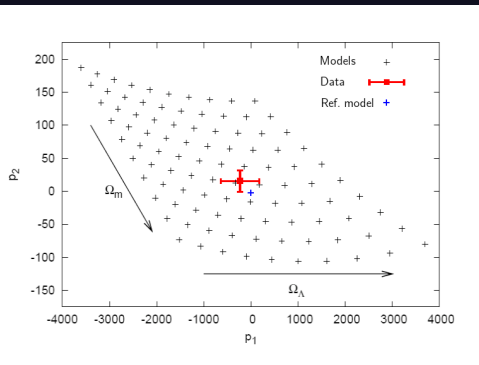
- 1 Construct **Training set** t_i that samples all possible behaviours of the observable
- 2 Define a **Reference vector** $\bar{t}$ which will be the origin of the new parameter space
- 3 Construct **Scatter Matrix** (analogous to a covariance matrix): $S = \Delta\Delta^T$ with $\Delta = (t_1 - \bar{t}, t_2 - \bar{t}, \dots, t_M - \bar{t})$ containing the differences between each training vector t_i and $\bar{t}$
- 4 Principal Components, w_i , calculated by solving the eigenvalue problem: $w_i = \lambda_i S w_i$
- 5 Optimal basis for that training set $\rightarrow$ concentrates information about deviations from the reference model in a few features (by definition)

The PC Basis

Initial Matrix:

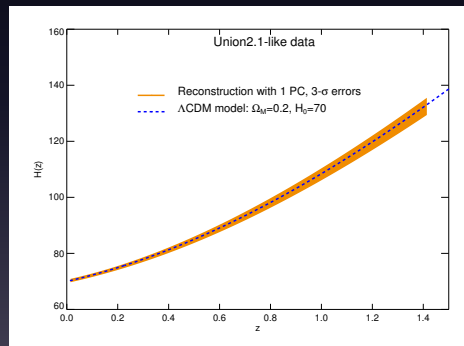
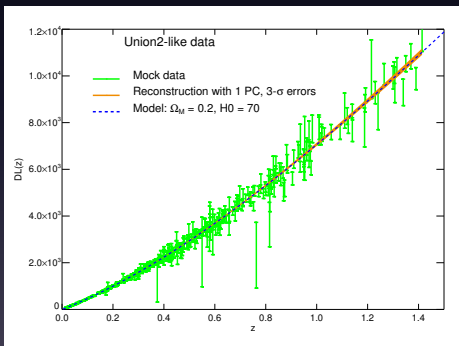
$$0.1 < \Omega_m < 0.5 ; 0.5 < \Omega_\Lambda < 0.9$$

PCs for the Union2.1 catalog



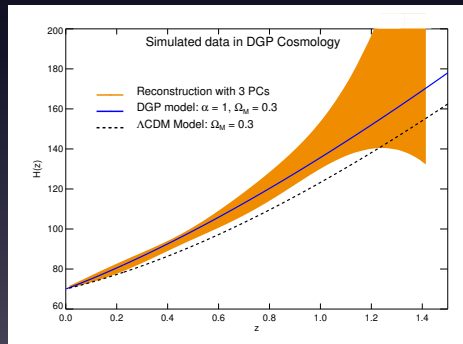
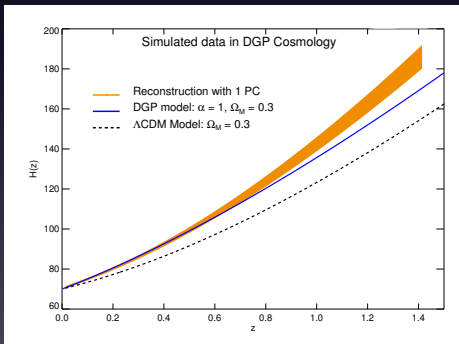
Tests with Mock Data

- Λ CDM: Several realizations varying Ω_m



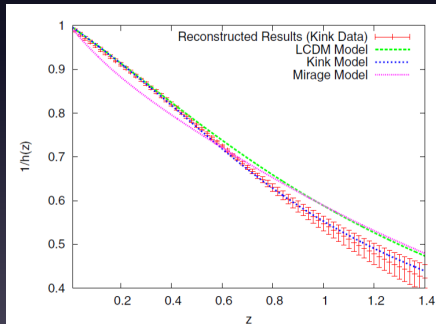
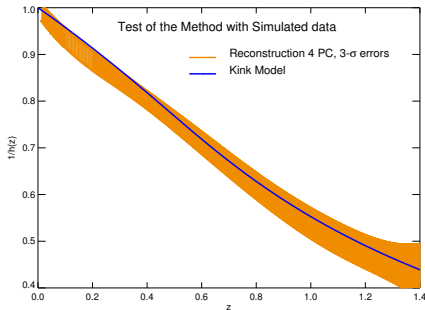
Tests with Mock Data

- ▶ Dvali-Gabadadze-Porrati (DGP) scenario
- ▶ IF higher number of PCs need to be calculated, indication that underlying cosmology might not be Λ CDM



The Kink Model

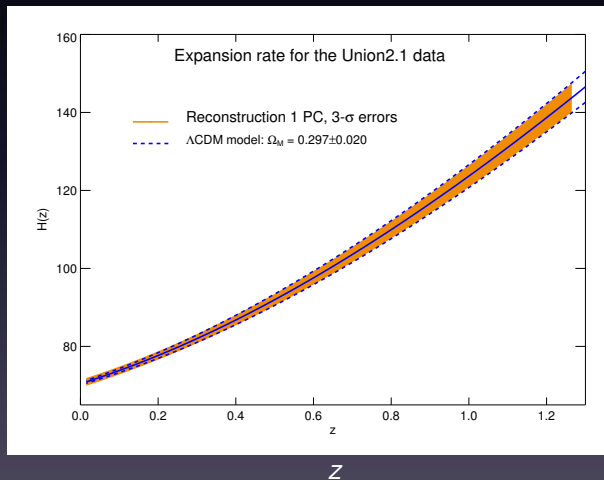
- ▶ Extreme model with rapidly varying equation of state
- ▶ Reconstruction consistent with other techniques: Shafieloo & Kim 2011



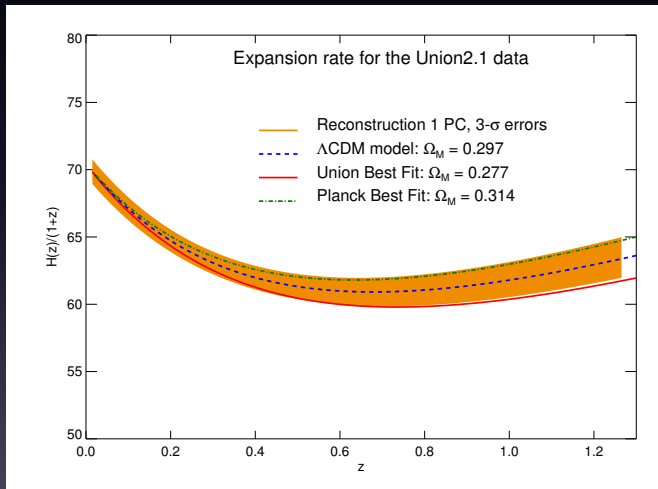
The Union2.1 Sample

- From the behaviour of $H(z) \rightarrow \Omega_m = 0.297 \pm 0.020$

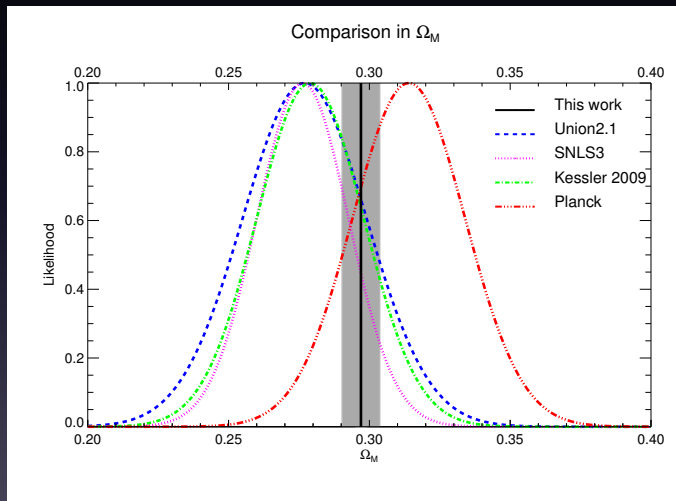
$H(z)$



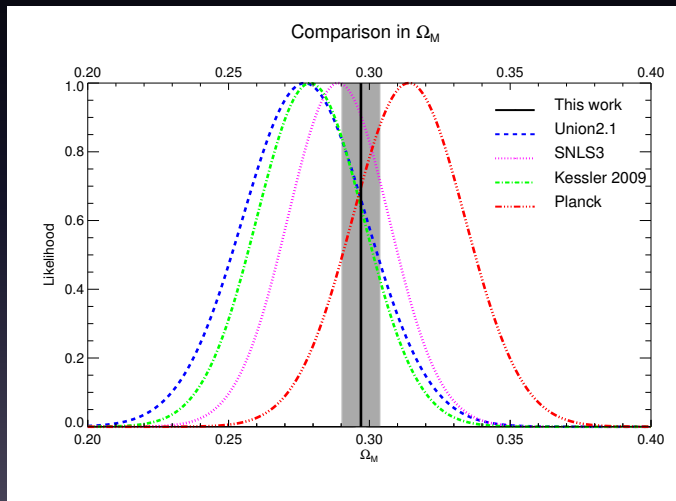
The Union2.1 Sample



Cosmological Constraints



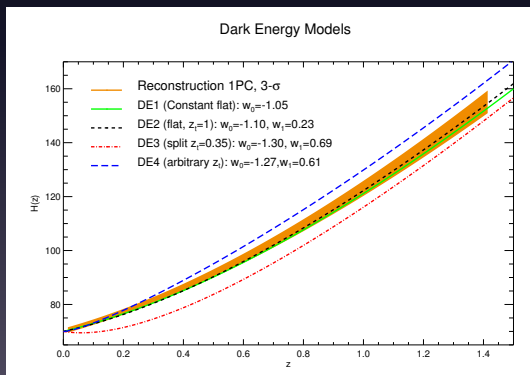
Cosmological Constraints



Other Cosmologies

1. **Dark Energy models:** extension of Chevallier+Polarski parameterization with varying Z_t (*Rapetti et al. 2005*)

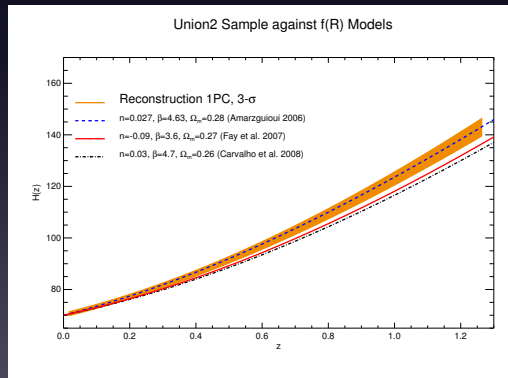
$$w(a) = \frac{w_{\text{et}}z + w_0z_t}{z + z_t} = \frac{w_{\text{et}}(1-a)a_t + w_0(1-a_t)a}{a(1-2a_t) + a_t}$$



Other Cosmologies

2. $f(R)$ models (*Carvalho et al. 2008*): $f(R) = R - \beta/R^n$

$$H^2 = H_0^2 \left[\frac{3\Omega_{m0}(1+z)^3 + f/H_0^2}{6f'\xi^2} \right]$$



Other Cosmologies

3. Kinematic models (*Guimaraes et al. 2009*)

Several parameterizations of the deceleration parameter $q(z)$

$$M0 : q(z) = q_0$$

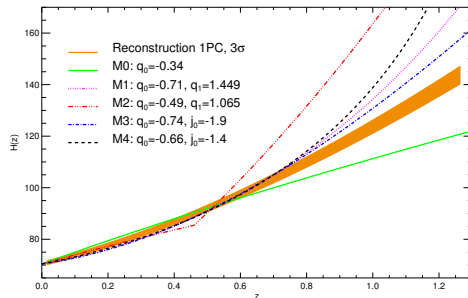
$$M1 : q(z) = q_0 + q_1 z$$

$$M2 : q(z) = q_0 \text{ for } z < z_t;$$

$$q(z) = q_1 \text{ for } z > z_t$$

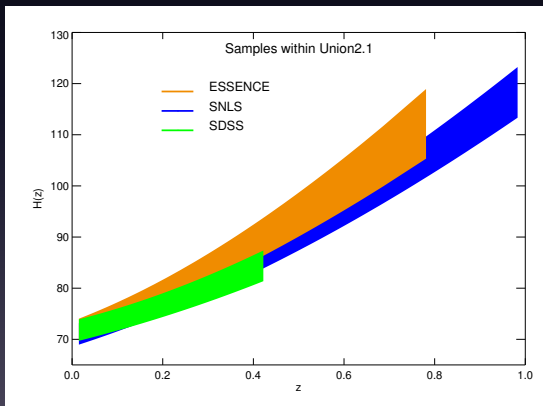
$$M3 : j(z) = j_0$$

Union2 Sample Vs. Kinematic Models



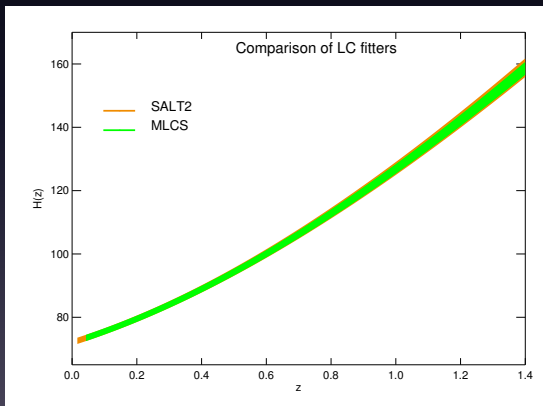
Systematic Errors

- Do different subsamples provide the same expansion history?



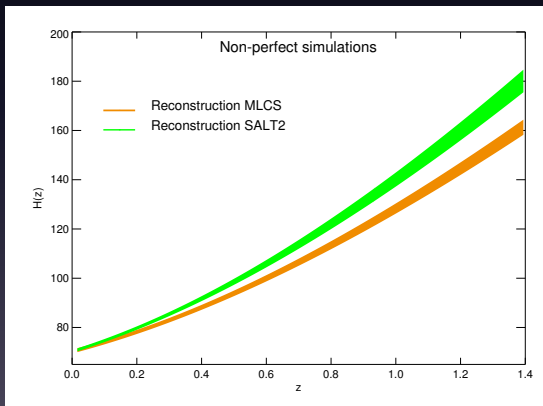
Systematic Errors

- ▶ Impact of light curve fitters in reconstructed $H(z)$
- ▶ Perfect simulations with **SNANA**



Systematic Errors

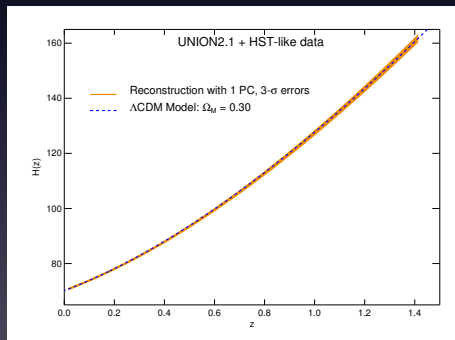
- Impact of light curve fitters in reconstructed $H(z)$
- Non-perfect simulations



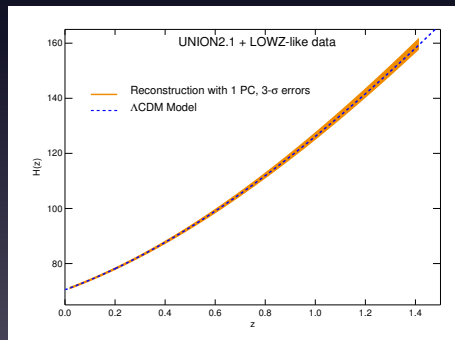
Testing Redshift Ranges

- Design a survey that is most sensitive to the expansion history in some specific redshift interval

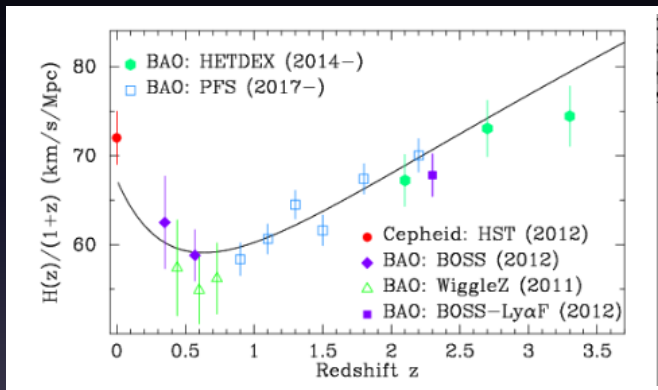
Union2.1 + High redshift



Union2.1 + Low redshift



Other Cosmological Probes



(Kim et al. 2013)

- Include **BAOs**: $D_L = (1+z)^2 D_A$

Summary

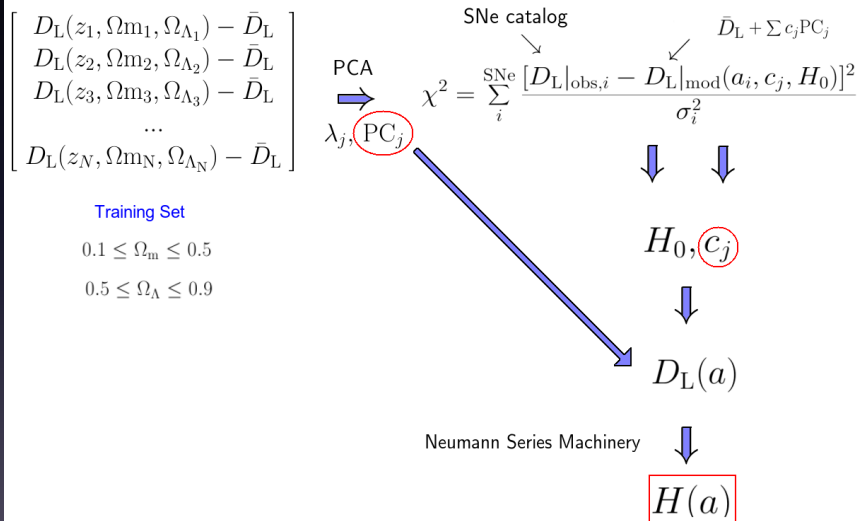
- 1) **Model-independent** approach to recover $H(z)$
- 2) Based on inversion of the distance-redshift relationship + **PCA** to build an *optimal basis* for the observable D_L
- 3) Framework to compare and constrain cosmological scenarios
- 4) Value for Ω_m agrees with new SNe constraints and Planck results

Looking ahead:

- ▶ BAO data and forecasts from future surveys, e.g. eBOSS
- ▶ Inhomogeneity: forecasts from LSST for inhomogeneous models

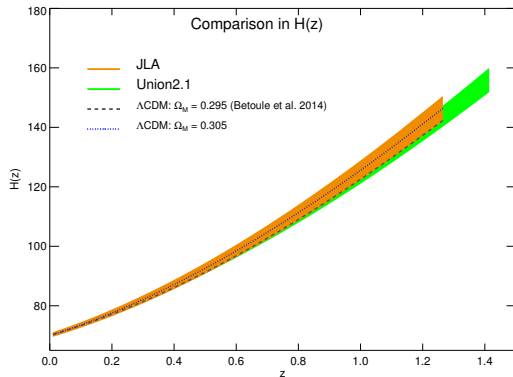
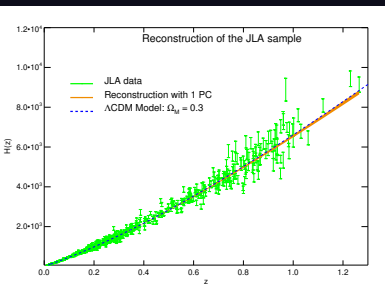
THANK YOU
FOR YOUR ATTENTION!

Method Summary



APPENDIX

The JLA sample



Uncertainties

- **Fisher matrix** of the χ^2 function

$$F_{ij} \equiv \left\langle \frac{\partial^2 \chi^2}{\partial c_i \partial c_j} \right\rangle \rightarrow F_{ij} = \sum_{k=0}^{\text{SNe}} \frac{p_i(a_k) p_j(a_k)}{\sigma_k^2}$$

- Errors satisfy $(\Delta c_i)^2 = (F^{-1})_{ii}$ and propagate into estimate of $\Delta D_L(a)$ and $\Delta e(a)$
- For D_L :

$$[\Delta D_L(a)]^2 = \sum_{j=0}^J (\Delta c_j)^2 [p_j(a)]^2$$

- For $e(a)$:

$$[\Delta e(a)]^2 = \sum_{j=0}^J (\Delta c_j)^2 [e_n^j(a)]^2$$

- Error in determination of PCs: $\Delta\lambda = \frac{1}{\sqrt{\lambda}}$

- Total error budget

$$\Delta_{\text{T}}^2 = \Delta_{\text{coeff}}^2 + \Delta_{H_0}^2 + \Delta_{\lambda}^2$$

Equations

Hilbert-Einstein action

$$S = \frac{-1}{16\pi G} \int d^4x R \sqrt{-g}$$

Friedmann Equation:

$$H^2(a) = H_0^2 \left[\frac{\Omega_{r0}}{a^4} + \frac{\Omega_{m0}}{a^3} - \frac{\Omega_{k0}}{a^2} + \Omega_{de0} F(a) \right]$$

$$F(a) = \exp \left[-3 \int_a^1 [1 + w(a')] \frac{da'}{a'} \right]$$

Equations II

Dark Energy

$$w(a) = \frac{w_{\text{et}}z + w_0 z_t}{z + z_t} = \frac{w_{\text{et}}(1-a)a_t + w_0(1-a_t)a}{a(1-2a_t) + a_t}$$

DGP

$$H^2 - \frac{H^\alpha}{r_c^{2-\alpha}} = H_0^2 \left[\frac{\Omega_m}{a^3} \right]; \quad r_c = (1 - \Omega_m)^{\frac{1}{\alpha-2}} H_0^{-1}$$

$f(R)$

$$H^2 = H_0^2 \left[\frac{3\Omega_{m0}(1+z)^3 + f/H_0^2}{6f'\xi^2} \right]$$

where ξ^2 depends on first and second derivatives of $f(R)$

KINEMATIC MODELS

$$M0 : q(z) = q_0$$

$$M1 : q(z) = q_0 + q_1 z$$

$$M2 : q(z) = q_0 \text{ for } z < z_t; q(z) = q_1 \text{ for } z > z_t$$

$$M3 : j(z) = j_0$$

Light Curve fitter models

SALT2

$$F(SN, p, \lambda) = x_0 \times [M_0(p, \lambda) + x_1 \times M_1(p, \lambda)] \times \exp[c \times CL(\lambda)]$$

$$\mu_B = m_B + \alpha x_1 - \beta c - M_B$$

MLCS

$$m_{\text{model}}^{e,f} = M^{e,f'} + p^{e,f'} \Delta + q^{e,f'} \Delta^2 + x_{\text{host}}^{e,f'} + K_{ff'}^e + \mu + x_{\text{MW}}^{e,f'}$$

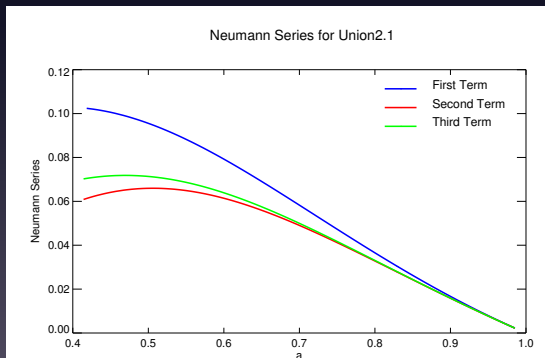
where $f \equiv$ observer frame filters; $f' \equiv$ rest-frame filters.

$M^{e,f'} \equiv$ absolute mag. for a SN with $\Delta = 0$

Neumann Series

$$e(a) = -a^3 D'_L + a \int_1^{a_{\min}} x^{-2} e_{n-1}(x) dx$$

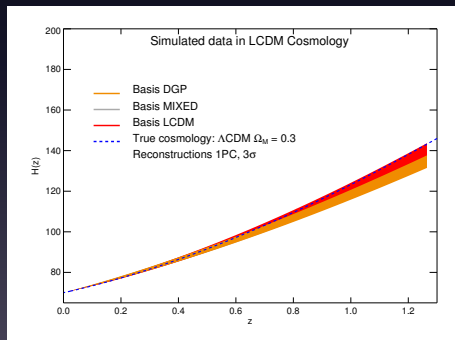
$$e(a) = \sum_{j=1}^J c_j e^{(j)}(a) = \sum_{j=1}^J c_j \left[-a^3 P C_j + \sum_{n=1}^3 a^n \int_1^{a_{\min}} x^{-2} e_{n-1}(x) dx \right]$$



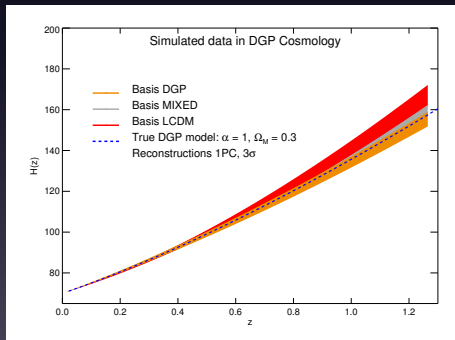
Extending the Training Set

- ▶ Training Set composed by Λ CDM behaviours only
- ▶ Make this initial matrix as general as possible
- ▶ Three different basis: Λ CDM, DGP and mixed

Λ CDM simulated data



DGP simulated data

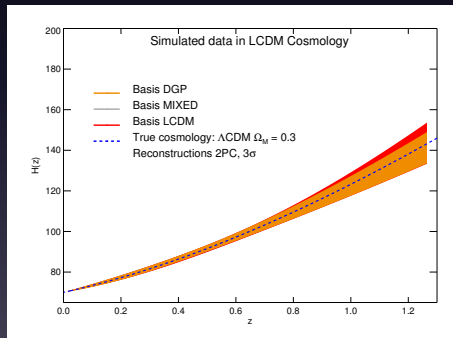


(Benitez-Herrera et al. in prep)

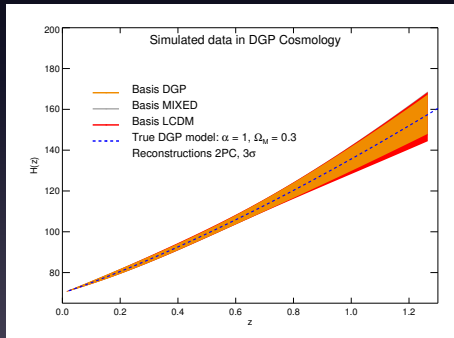
Extending the Training Set

- With 2PCs, the dependance with basis is negligible

Λ CDM simulated data



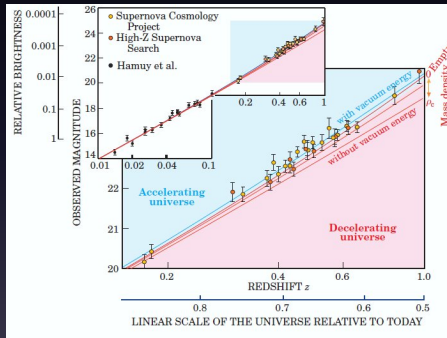
DGP simulated data



(Benitez-Herrera et al. in prep)

Cosmology with SNe Ia

- ▶ SNe beyond $z \sim 0.5$ appear **fainter** than expected for an empty universe
- ▶ First evidence of **accelerated expansion**
- ▶ Later confirmed by CMB and BAO measurements
- ▶ **Nobel Prize** in Physics 2011
S. Perlmutter, B. Schmidt, A. Riess



(Perlmutter et al. 1999)

- For D_L :

$$[\Delta D_L(a)]^2 = \sum_{j=0}^J (\Delta c_j)^2 [p_j(a)]^2$$

- For $e(a)$:

$$[\Delta e(a)]^2 = \sum_{j=0}^J (\Delta c_j)^2 [e_n^j(a)]^2$$

Mocks

In 10 redshifts bins, z_i :

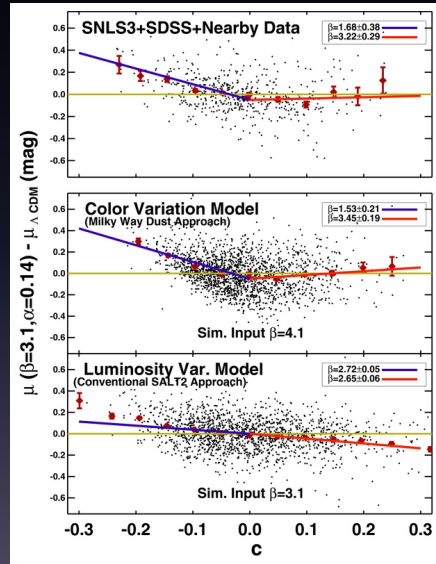
$$DL = DL_{model} + e$$

$$e = m + \sigma * randomnumber$$

with $\mu = 0, \sigma = 1$

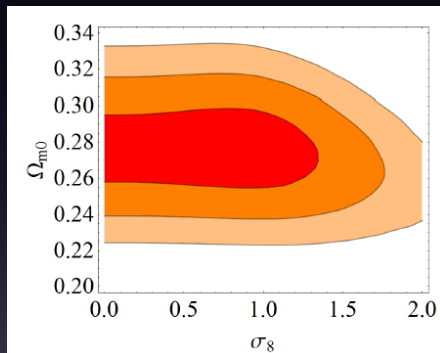
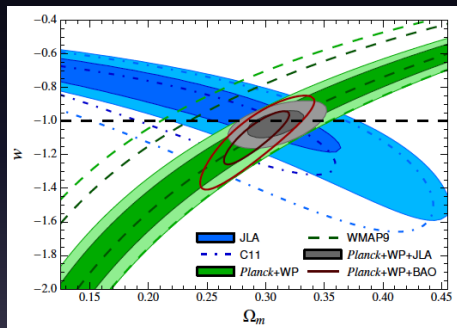
Systematic Errors

- ▶ SNANA-simulated samples with different systematics
- ▶ Test if detectable with PC algorithm



Current SN constraints

► Hubble diagram

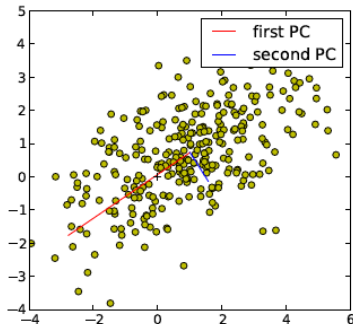


► $w = -1.027 \pm 0.055$
(Betoule et al. 2014)

► $\sigma_8 = 0.84^{+0.28}_{-0.65}$
(Castro & Quartin 2014)

Principal Component Analysis

- ▶ Statistical tool used in many astrophysical problems: e.g. reduction of data dimensionality



- ▶ Rotation of the parameter space: **new variables are uncorrelated and orthogonal**
- ▶ First components account for the largest variance in the data: directions of maximum clustering

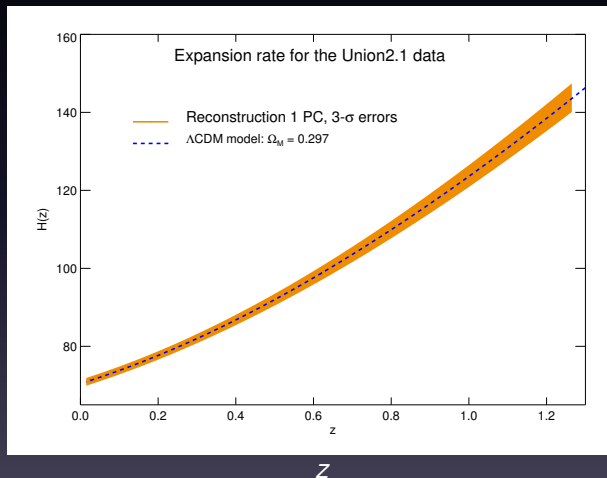
Principal Component Analysis

- ▶ X_p^k where k is the number of data, p parameters
- ▶ Compute the covariance matrix (symmetric):

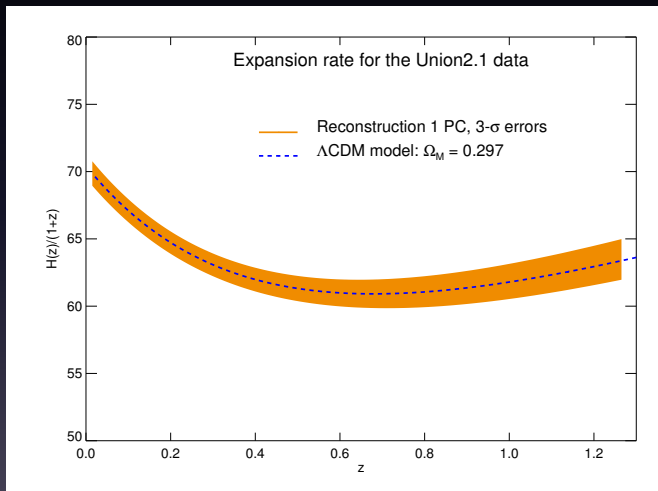
$$\Sigma_{ij} = \frac{\sum_k (X_i^k - \bar{X}_i)(X_j^k - \bar{X}_j)}{k}$$

- ▶ Diagonalize and sort the eigenvalues:
 $\lambda_1 > \dots > \lambda_p \rightarrow w_1(\lambda_1), \dots, w_p(\lambda_p)$
- ▶ Reproduce the initial data in a lower-dimensional space

The Reconstructed Expansion Rate

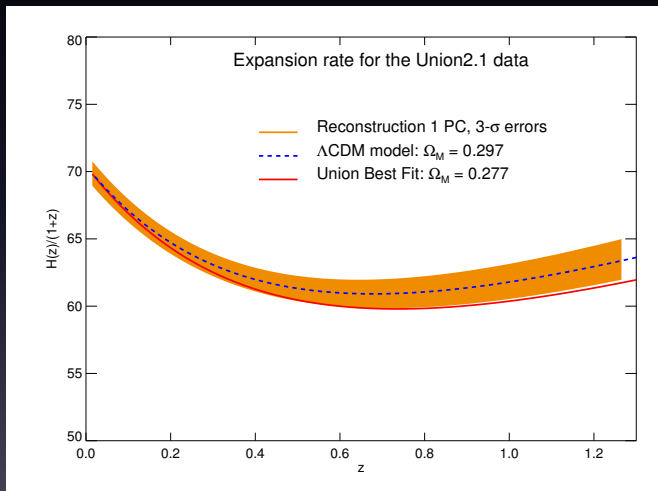
 $H(z)$ 

The Reconstructed Expansion Rate



(Benitez-Herrera et al. in prep)

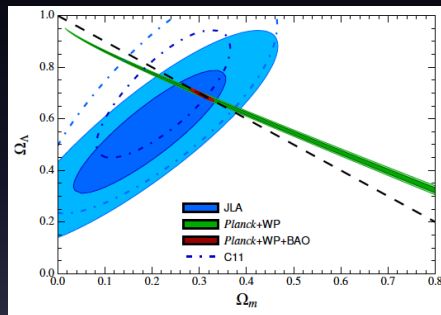
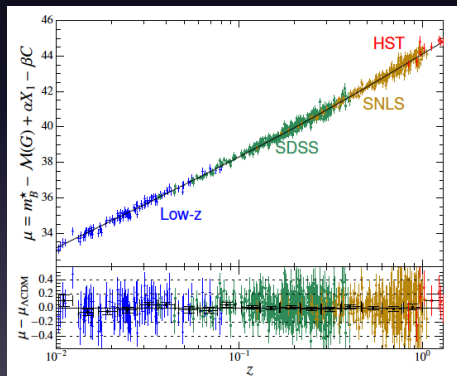
The Reconstructed Expansion Rate



(Benitez-Herrera et al. in prep)

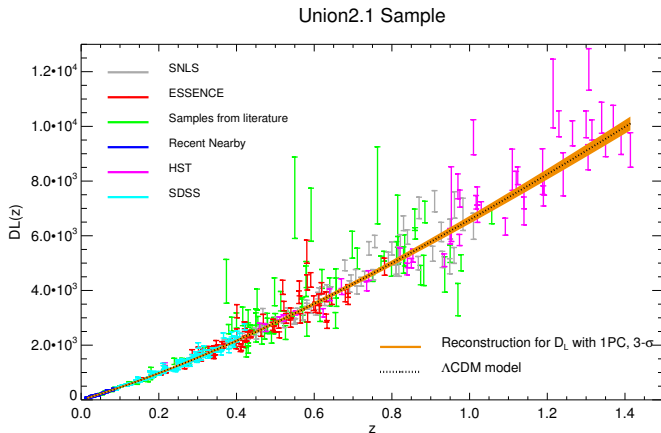
Current SN constraints

Hubble diagram: Intrinsic dispersion $\sigma_m \sim 0.15$ mag

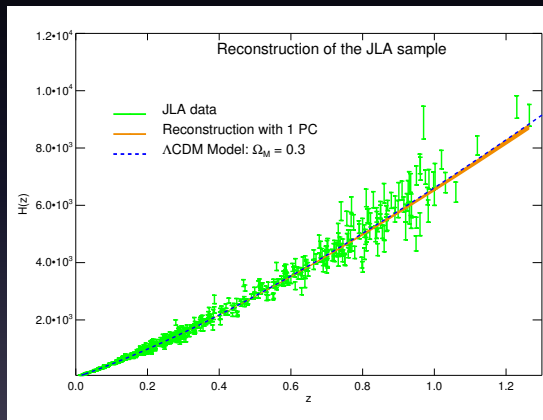
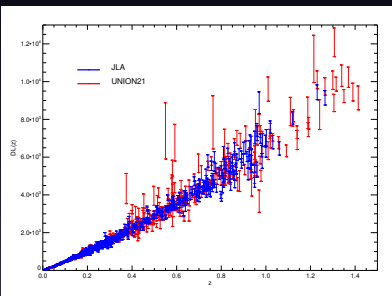


Consistent with a Λ CDM with
 $\Omega_m = 0.295 \pm 0.034$
(Betoule et al. 2014)

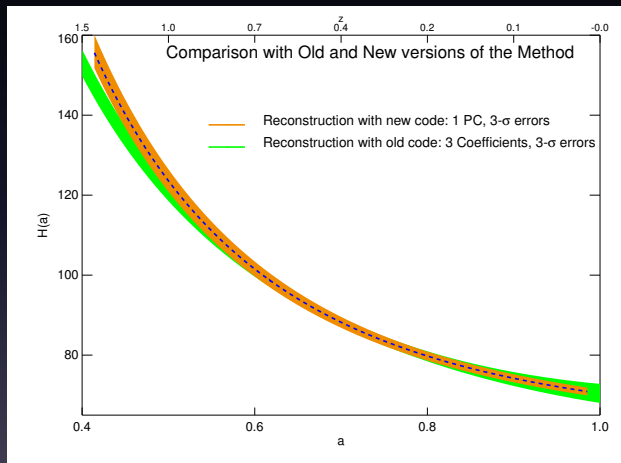
Union2.1 Sample *(Suzuki et al. 2012)*

 $D_L(z)$

 z *(Benitez-Herrera et al. 2013)*

The JLA sample



The Reconstructed $H(a)$

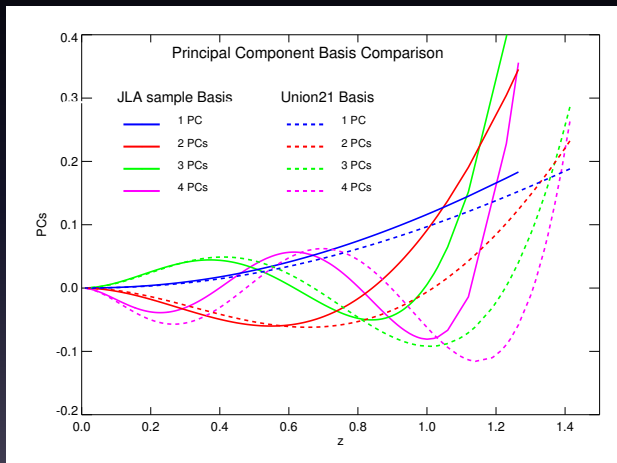
 $H(a)$  a

- Cumulative percentage of total variance

$$t_M = \left(\frac{\sum_{i=1}^L \lambda_i}{\sum_{j=1}^{N_{\text{PC}}} \lambda_j} \right) \times 100$$

- **Example:** Union2.1 data
1st PC accounts for **99.90%** of t_M

PC basis for real data



Systematic Errors II

Source	Error on w
Vega	0.033
All instrument calibration	0.030
Color correction	0.020
Mass correction	0.016
Contamination	0.016
Intergalactic extinction	0.013
Galactic extinction	0.010
Light curve shape	0.006
Quadrature sum	0.061