

# Records from Primordial Gravitational Waves and Cosmic Acceleration in CMB polarization

Carlo Baccigalupi

SISSA, Trieste

VI Challenges of New Physics in Space,  
Campos de Jordao, Brazil, May 25-29, 2015



# Records from Primordial Gravitational Waves and Cosmic Acceleration in CMB polarization

Part I: physical processes and cosmological relevance

Part II: status of observations

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# Outline

- Introduction to SISSA
- CMB anisotropies
- Initial conditions and Gravitational Waves
- Large Scale Structure and Dark Energy
- Concluding remarks, I

# The International School for Advanced Studies, Trieste, Italy

- Main goal: training PhD students to research
- Interdisciplinary:
  - physics (Astrophysics, Astroparticle, Elementary Particles, Condensed Matters, Statistical Physics),
  - Mathematics,
  - Neuro-sciences
- About 100 professors, 100 Post-Docs, 200 PhD students
- Located in Trieste, Italy, one of the places with the highest density of scientific institutions: the international center for theoretical physics, the University of Trieste, the Astronomical observatory in Trieste, research area laboratories



# The SISSA cosmology group: data analysis for CMB experiments

- PolarBear (Fabbian, Puglisi):
  - Map-making
  - Foreground control and cleaning, B mode power spectrum measurements
- Planck:
  - INAF-OATs and SISSA run the Low Frequency Instrument Data Processing Centre,
  - Component Separation for CMB and characterization (Basak, Bielewicz, Castex, Paci)
- EBEX (Castex, Puglisi)
  - Analysis of maps, support for foreground cleaning and power spectrum estimation
- CMB satellite/ground-based network proposals
  - Data analysis simulations



# The SISSA cosmology group: CMB-LSS Cross-Correlations

- Euclid:
  - CMB-N-body lensing simulations: geodesics (Calabrese, Fabbian)
  - CMB-N-body lensing simulations of cross-correlations (Bianchini, Calabrese)
- Herschel:
  - Galaxy formation and evolution (Lapi et al. for the Galaxy Formation group)
  - Cross-correlation with Planck (Bianchini, Bielewicz, Lapi et al. for the Galaxy formation group)



# CMB anisotropies

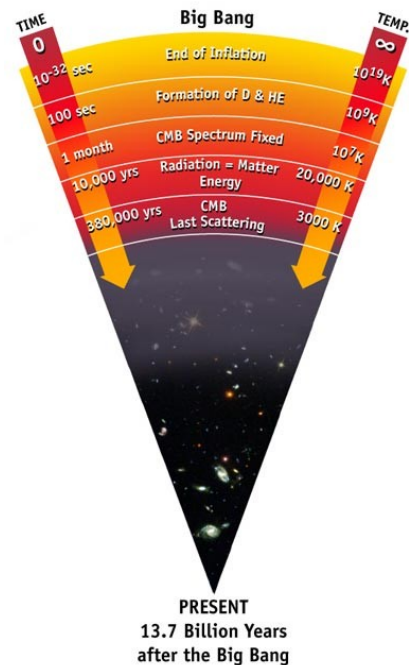
# CMB: where and when?

- Opacity:

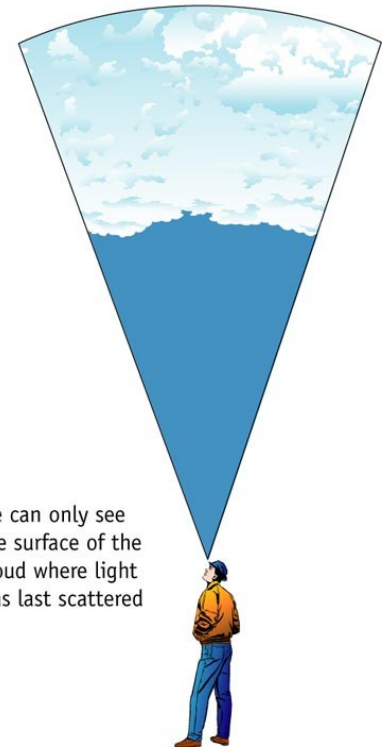
$$\lambda = 1/n\sigma \ll \text{horizon}$$

where the horizon is the distance at which information get at each time, inverse of the Hubble expansion rate

- Decoupling:  $\lambda \approx \text{horizon}$
- Free streaming:  $\lambda \gg \text{horizon}$
- Cosmological expansion, Thomson cross section and electron abundance conspire to activate decoupling about 380000 years after the Big Bang, at about 3000 K CMB photon temperature

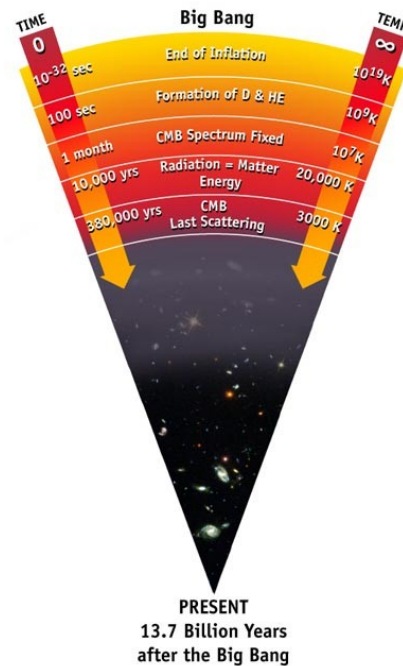


The cosmic microwave background Radiation's "surface of last scatter" is analogous to the light coming through the clouds to our eye on a cloudy day.

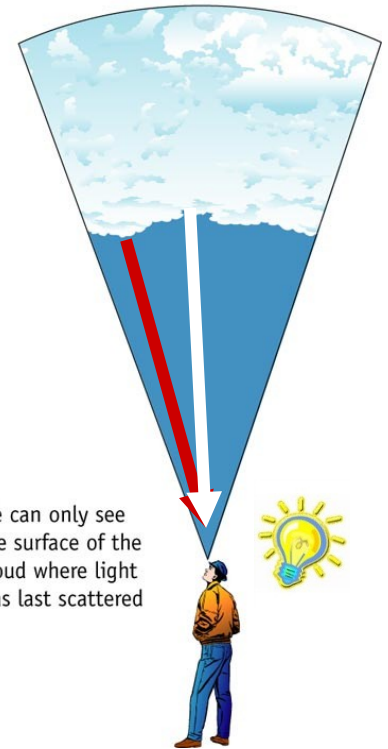


# A postcard from the big bang

- From the Stephan Boltzmann law, regions at high temperature should carry high density
- The latter is activated by perturbations which are intrinsic of the fluid as well as of spacetime
- Thus, the maps of the CMB temperature is a kind of snapshot of primordial cosmological perturbations

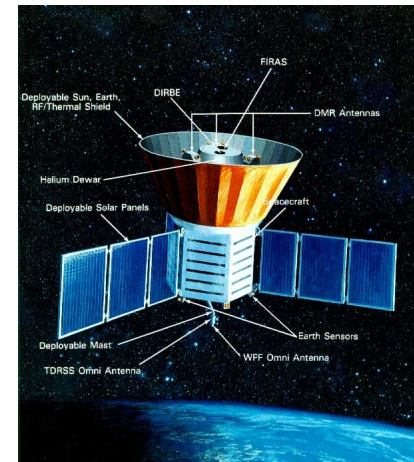


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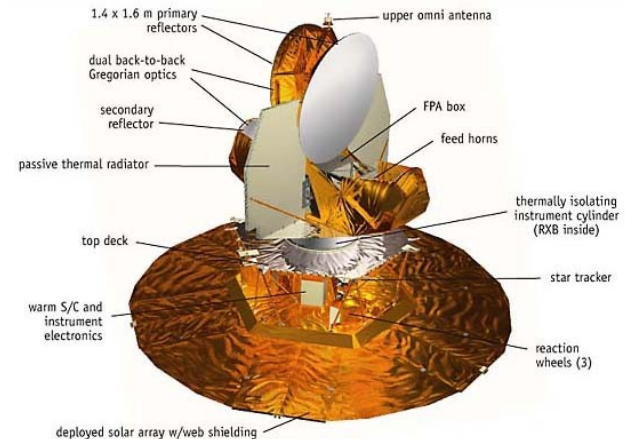
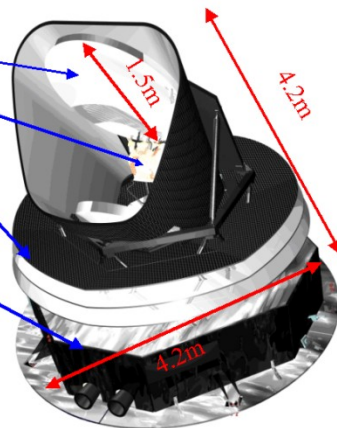


# From COBE to the Wilkinson Microwave Anisotropy Probe (WMAP) to Planck

- About 40 years of scientific and technological progresses
- Lots of experiments, people
- See [lambda.gsfc.nasa.gov](http://lambda.gsfc.nasa.gov)



- PLM (charge Utile)
  - Télescope (1.5m)
  - Instruments plans focaux
  - Cryo-structure (radiateur, baffles, V-grooves)
- SVM
  - Avionique (Contrôle (data, attitude), gestion puissance, Télécom+ Électroniques Instruments)
- Satellite:
  - 2000kg, 1600W, 4.2mx 4.2m
  - Durée de vie 21 mois
  - Satellite en rotation 1tour/min



# CMB physics: Boltzmann equation

$$\frac{d \text{ photons}}{dt} = \text{gravity} + \text{Compton scattering}$$

$$\frac{d \text{ baryons+leptons}}{dt} = \text{gravity} + \text{Compton scattering}$$

# CMB physics: Boltzmann equation

$d \text{ neutrinos}$

$dt$

= gravity + weak interaction

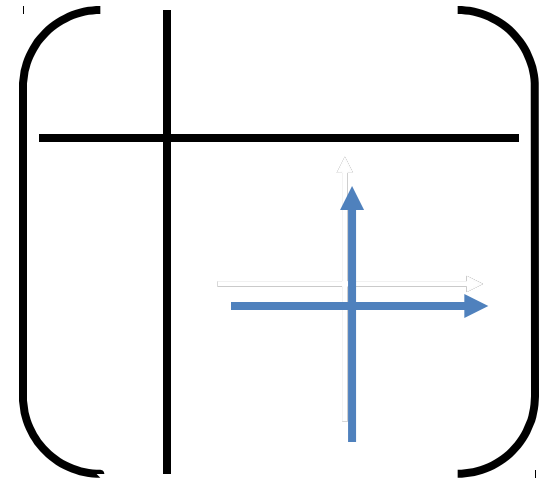
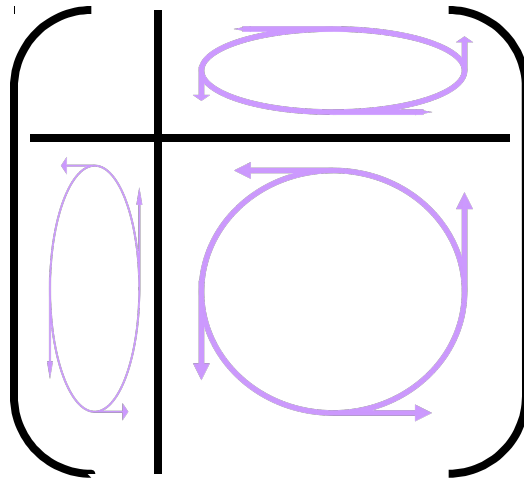
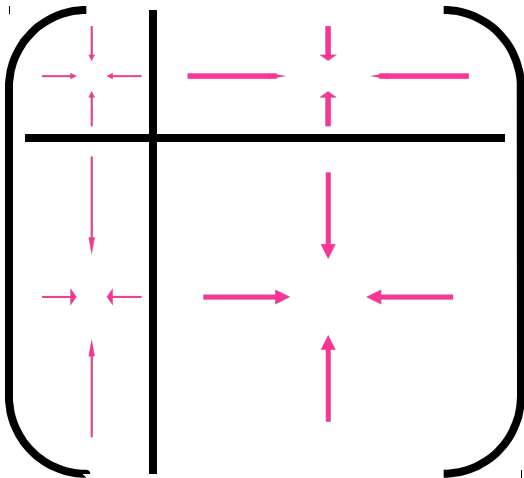
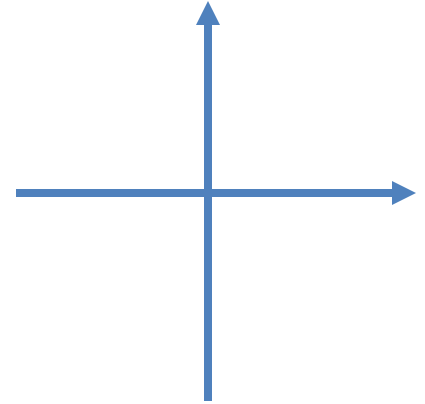
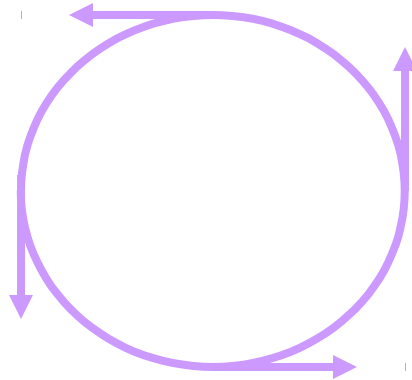
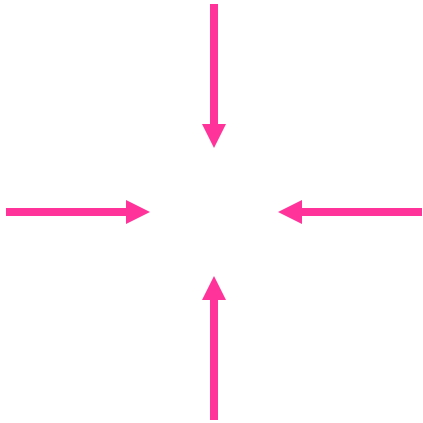
$d \text{ dark matter}$

$dt$

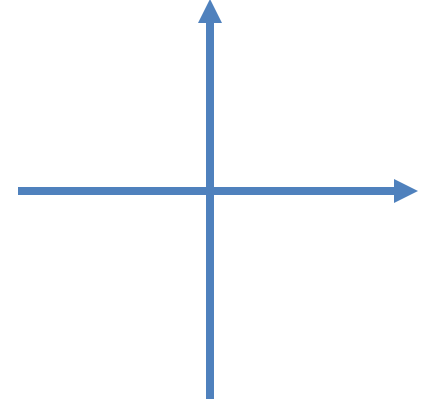
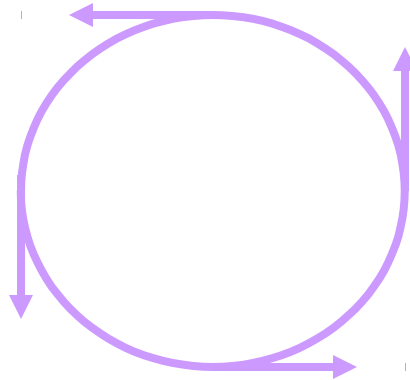
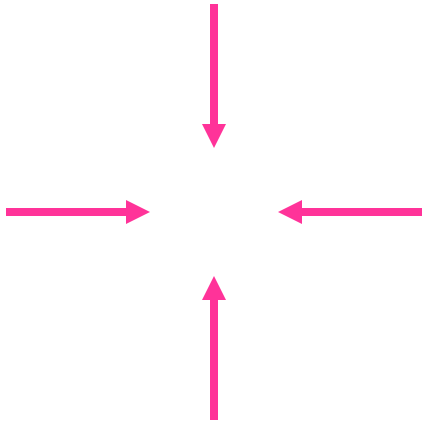
= gravity + weak interaction (?)

gravity = photons + neutrinos + baryons + leptons + dark matter

# CMB physics: gravity



# CMB physics: gravity



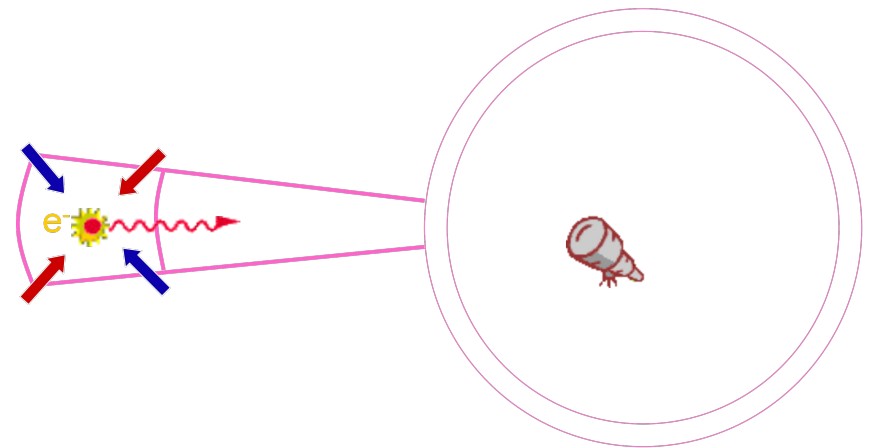
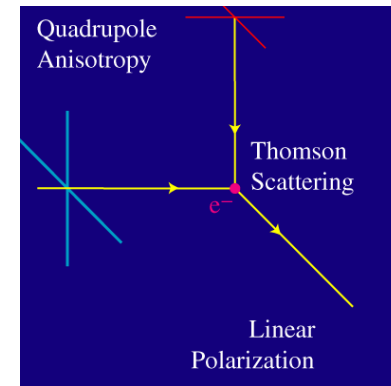
$$\begin{pmatrix} \Phi \\ \Psi \end{pmatrix}$$

$$\begin{pmatrix} B^{\pm 2} \\ B_{\pm 2} \end{pmatrix}$$

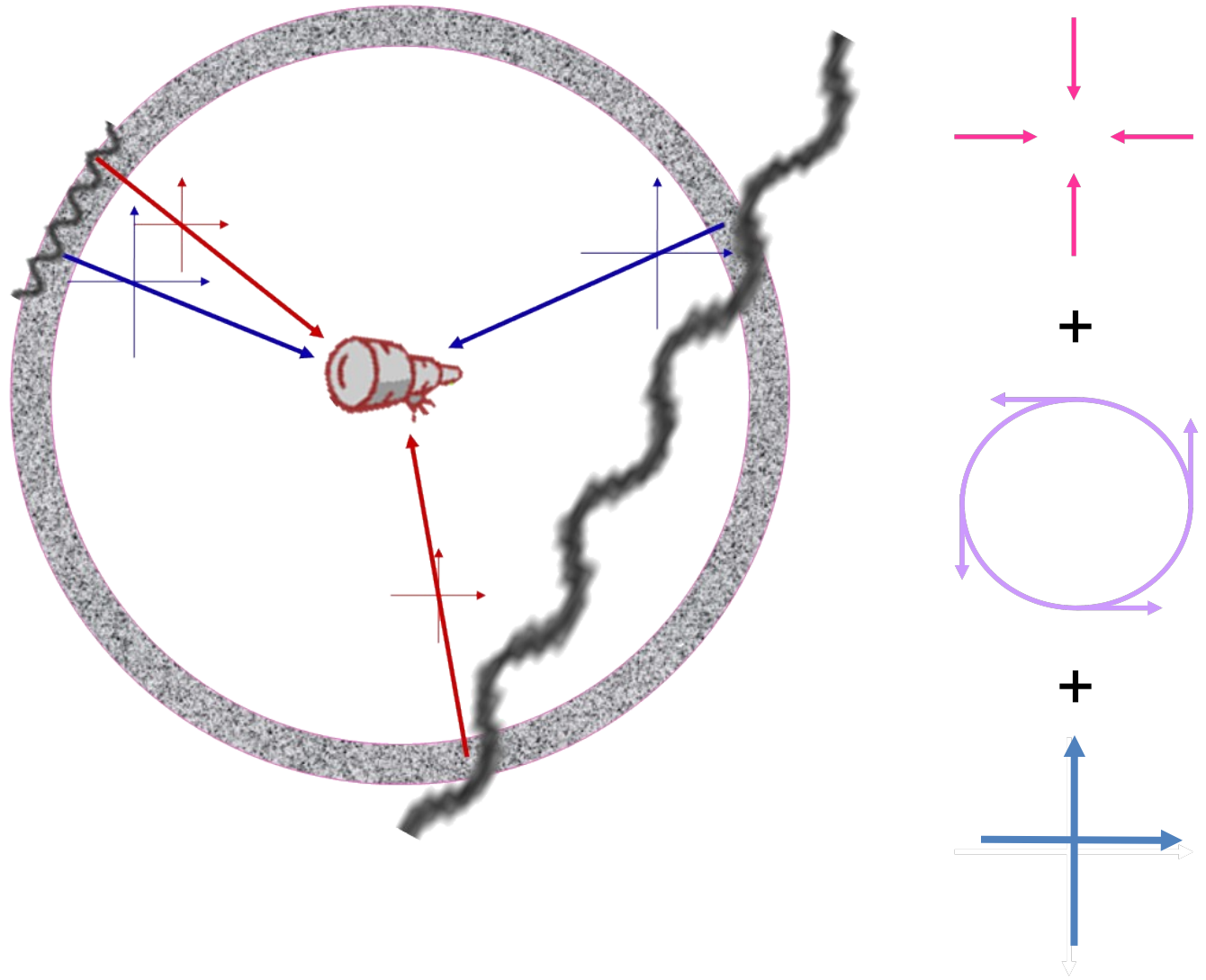
$$\begin{pmatrix} H_T^{\pm 2} \end{pmatrix}$$

# CMB Physics: Compton scattering

- Compton scattering is anisotropic
- An anisotropic incident intensity determines a linear polarization in the outgoing radiation
- At decoupling that happens due to the finite width of last scattering and the cosmological local quadrupole

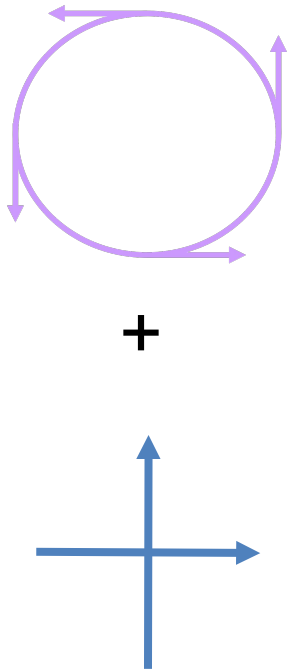


# CMB anisotropy: Total Intensity

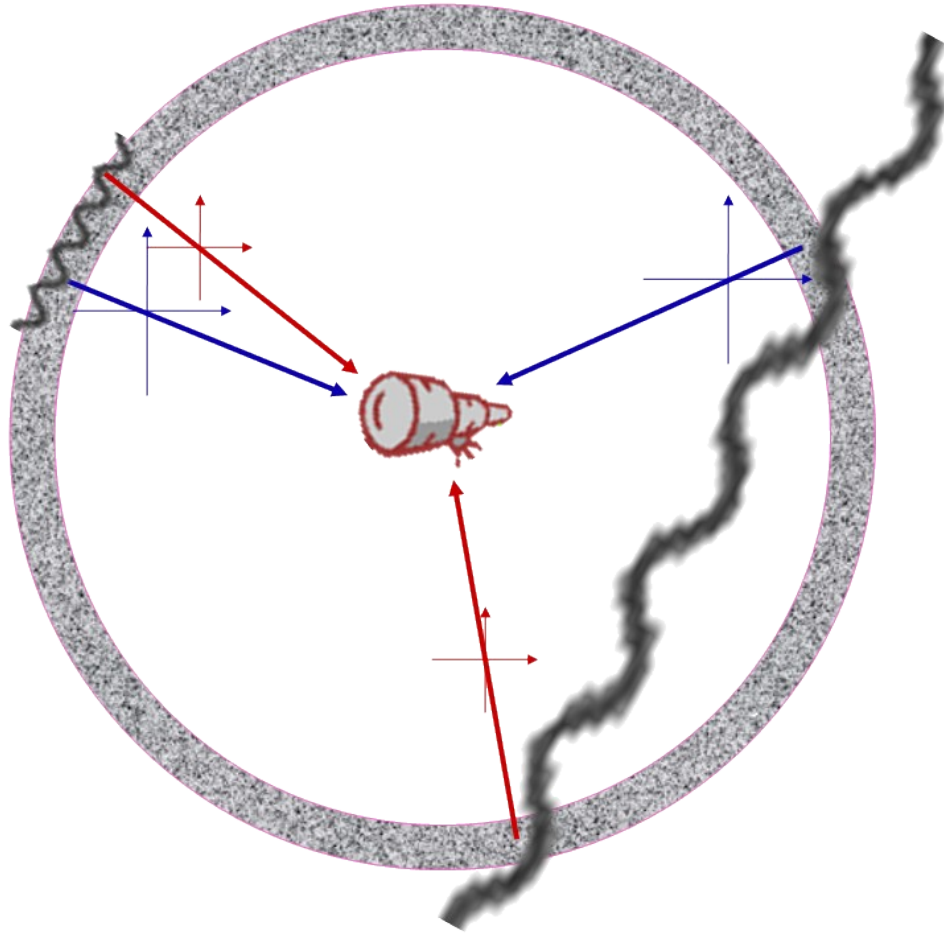
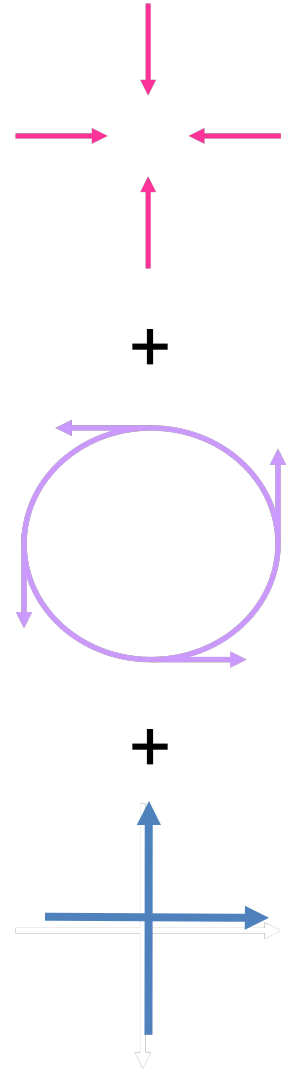


# CMB anisotropy: polarization

Curl (B):



Gradient (E):



# Anisotropies

$$T(\theta, \varphi), Q(\theta, \varphi), U(\theta, \varphi), V(\theta, \varphi)$$

spherical  
harmonics

$$X(\theta, \varphi) = \sum_{lm} a_{lm}^X Y_{lm}^s(\theta, \varphi)$$

$$X = T, E, B$$

$$s=0 \text{ for } T, 2 \text{ for } Q \text{ and } U$$

E and B modes have opposite parity

# Angular power spectrum

$$T(\theta, \varphi), Q(\theta, \varphi), U(\theta, \varphi), V(\theta, \varphi)$$

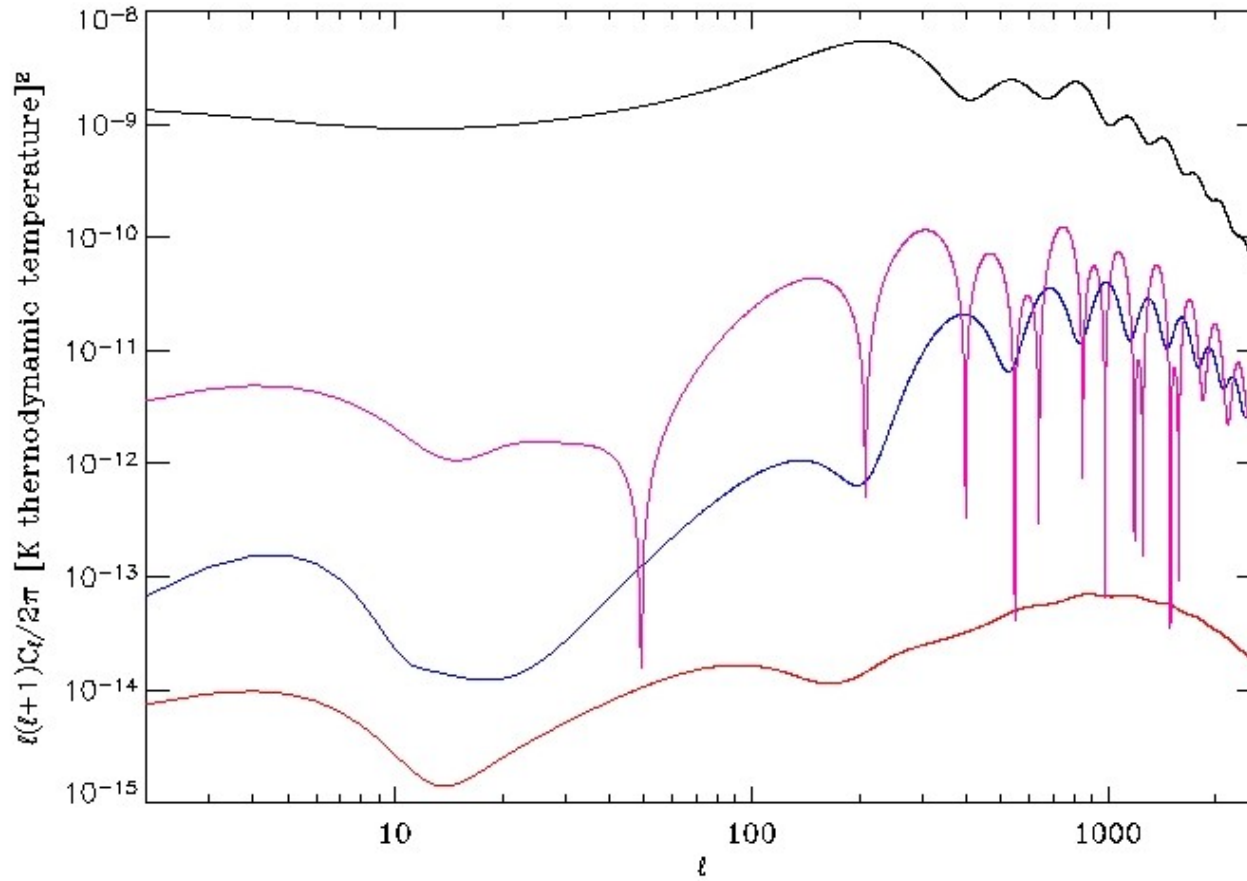
spherical  
harmonics

$$a_{lm}^X, X=T, E, B$$

information  
compression

$$C_l = \sum_m [(a_{lm}^X)(a_{lm}^Y)^*]/(2l+1)$$

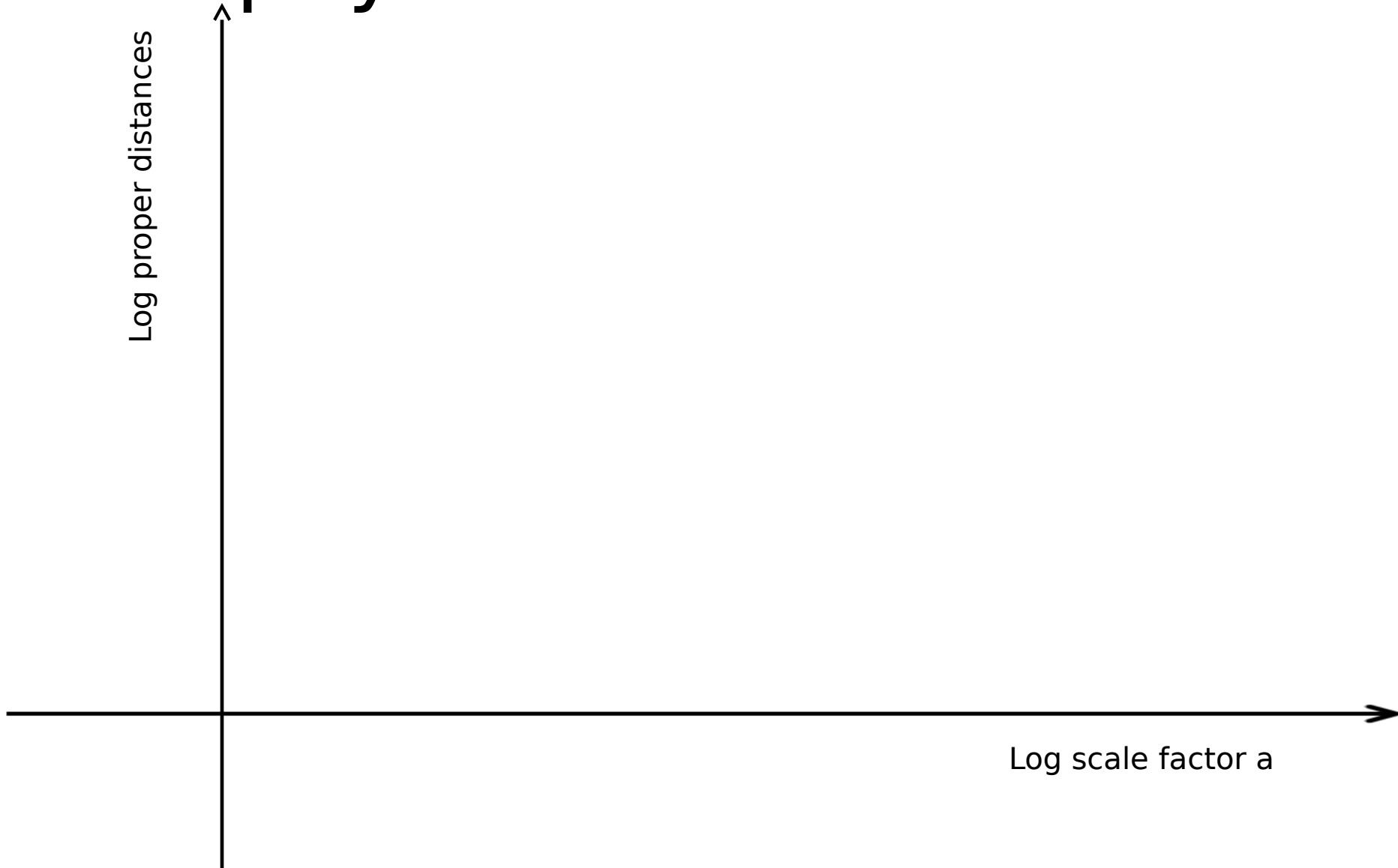
# CMB angular power spectrum



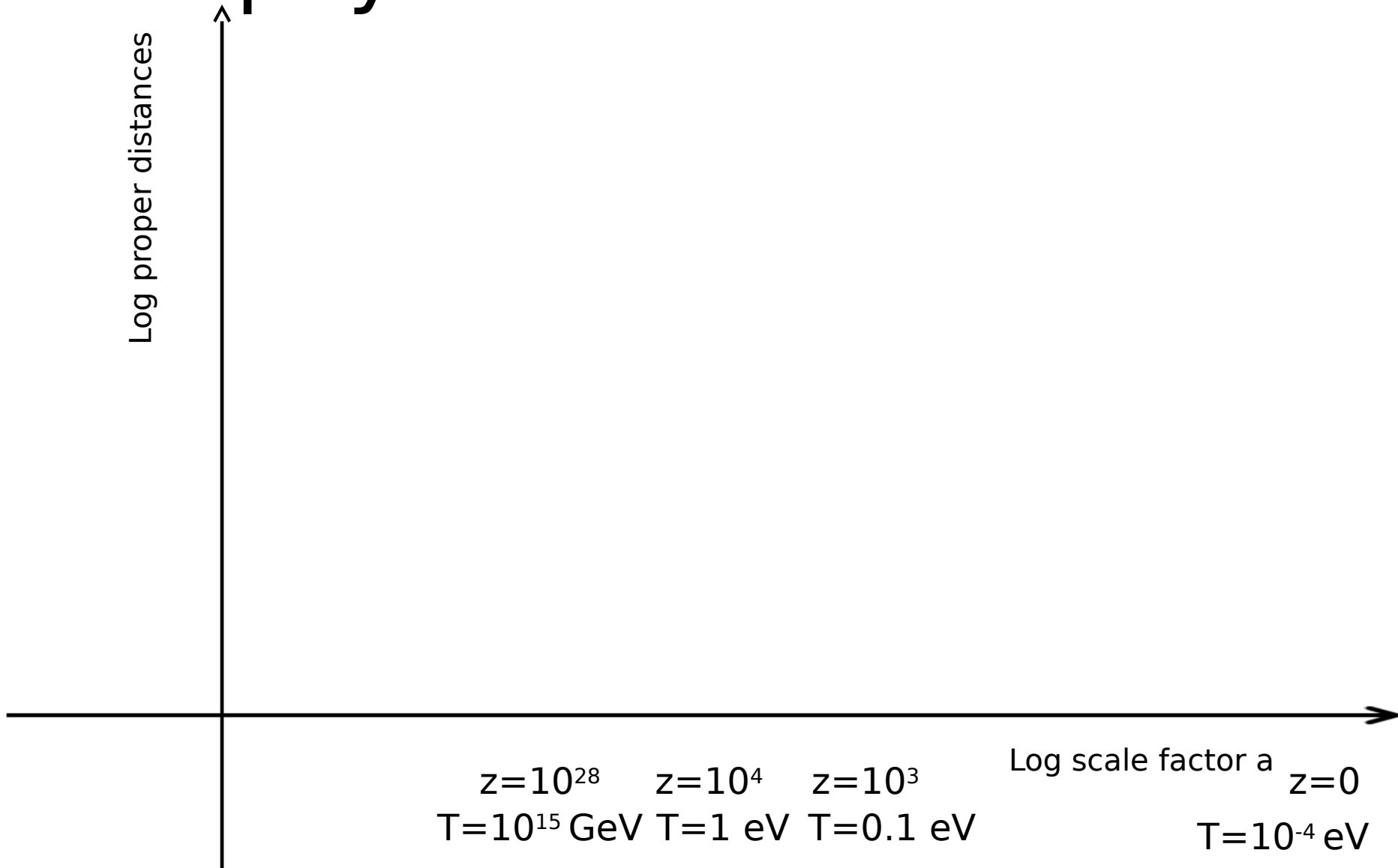
Angle  $\approx 200/l$  degrees

# Initial conditions and Gravitational Waves

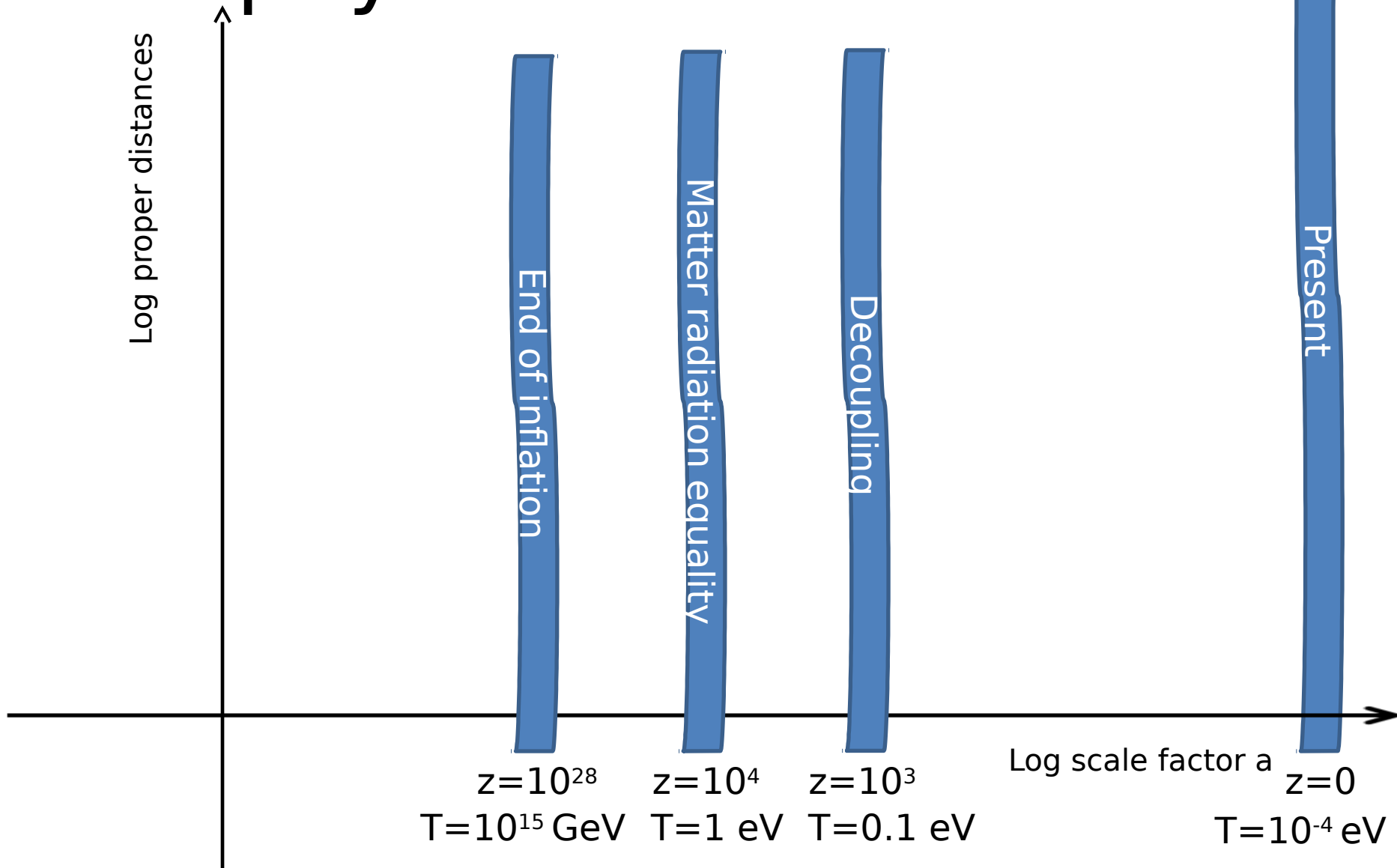
# CMB physics: initial conditions



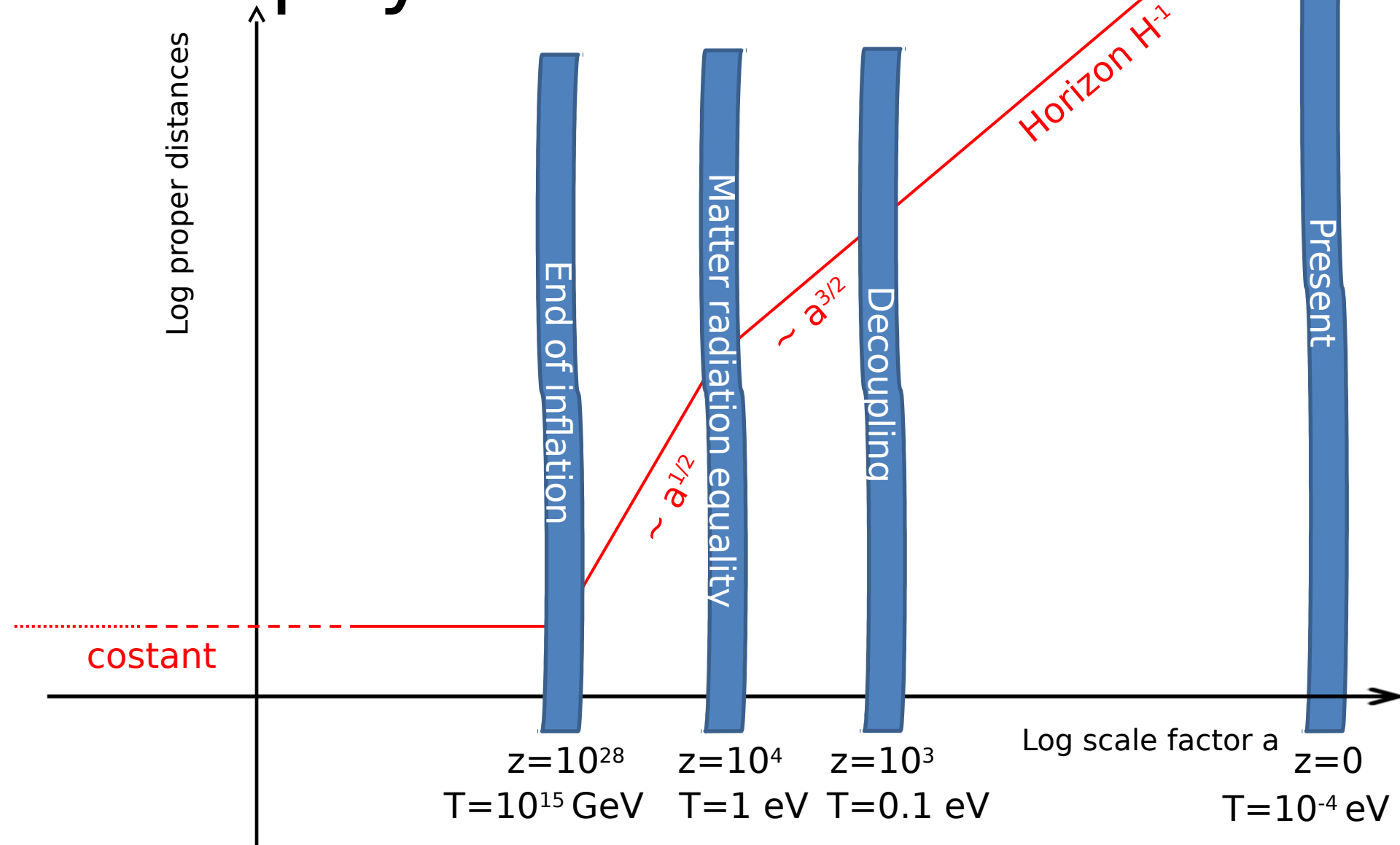
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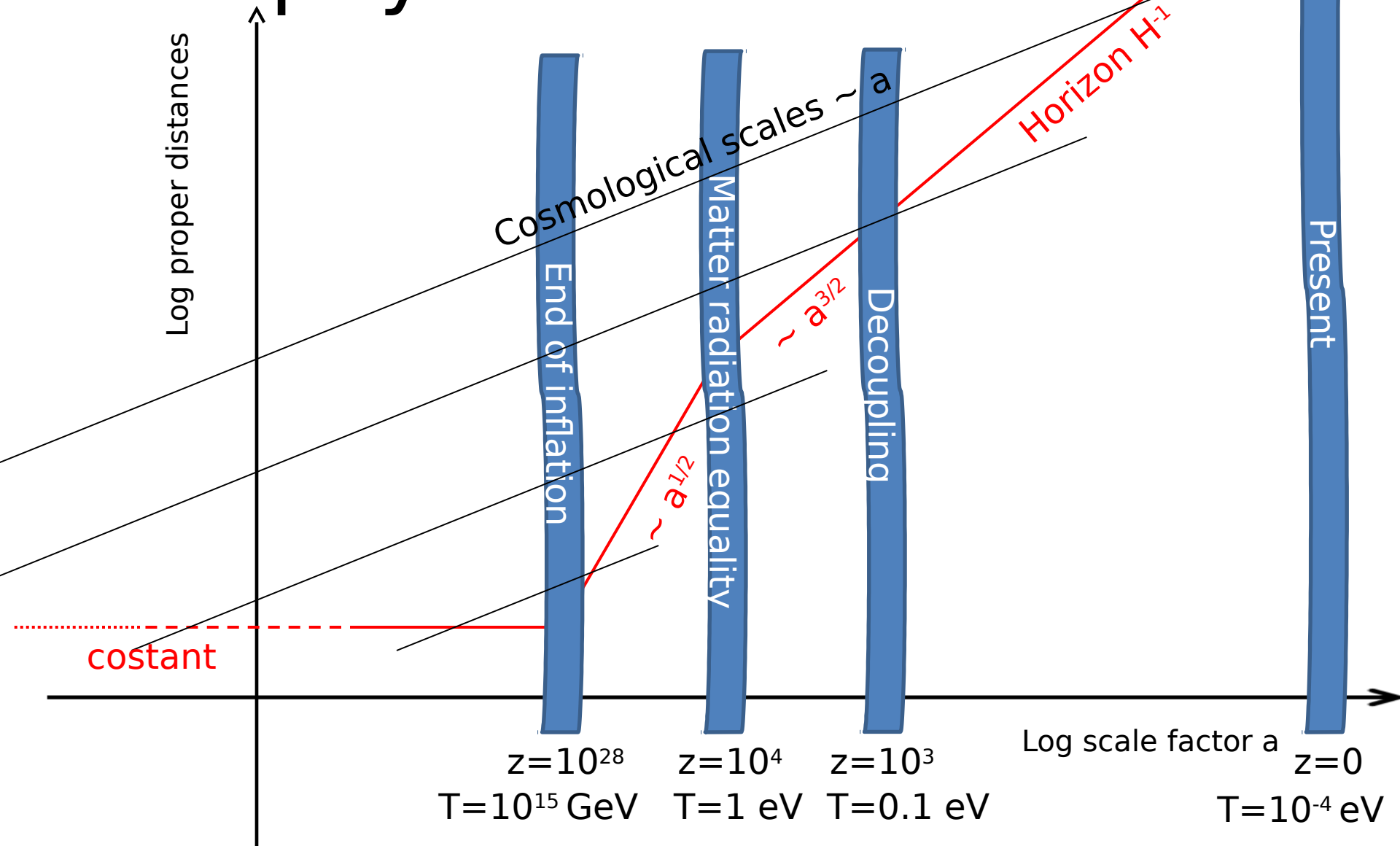
# CMB physics: initial conditions



# CMB physics: initial conditions



# CMB physics: initial conditions



# CMB physics: initial conditions

End of inflation

$\Delta(\vec{x}, t)$

Fluctuations are in the linear regime:  $\Delta \ll 1$

Fourier decomposition:  $\Delta(\vec{x}, t) \rightarrow \Delta(\vec{k}, t)$

Assumed statistics:

$$\langle \Delta(\vec{k}, t) \bar{\Delta}(\vec{k}', t) \rangle = P_{\Delta}(\vec{k}, t) \delta(\vec{k} - \vec{k}')$$

Gaussianity

# CMB physics: initial conditions

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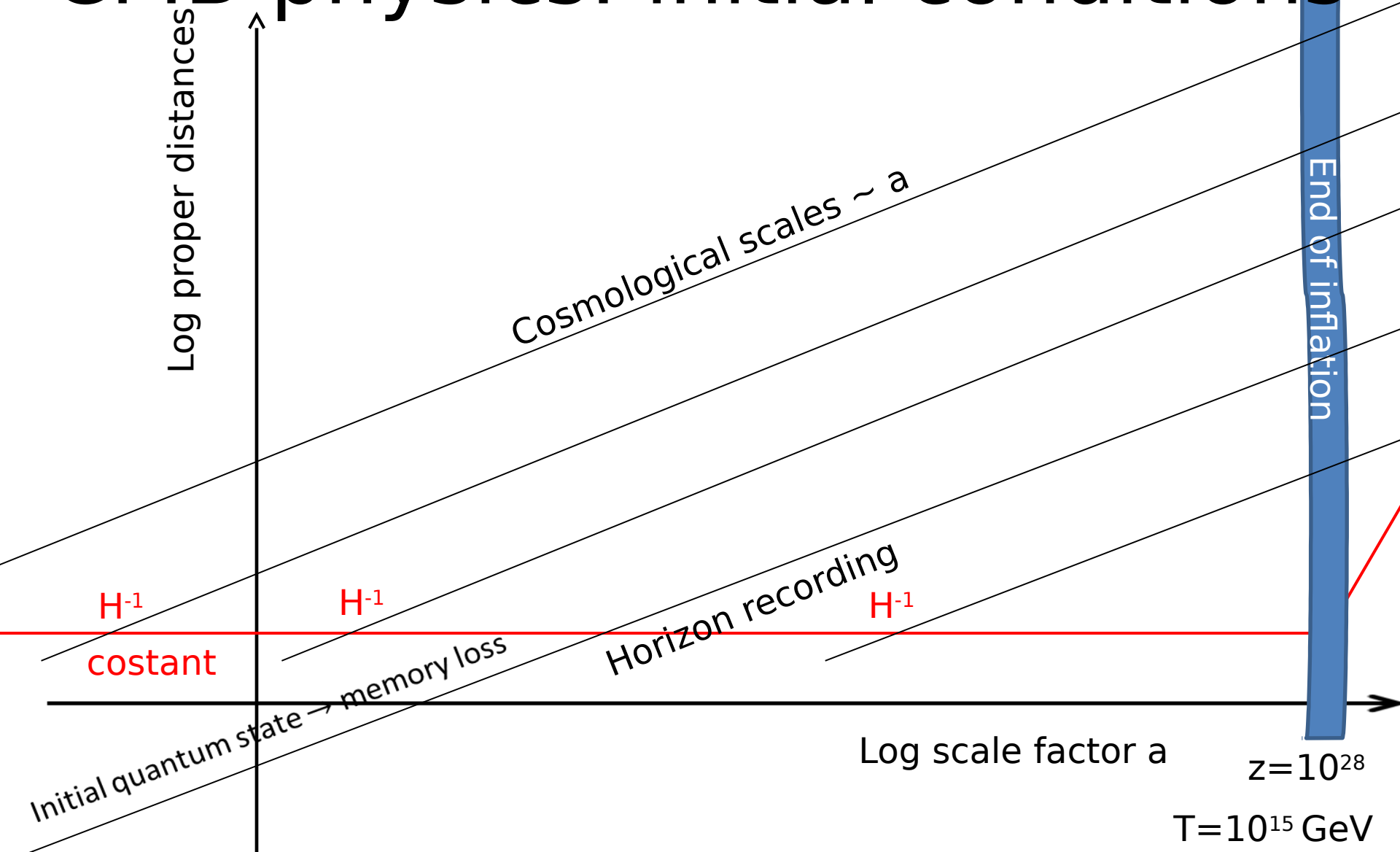
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$$\langle \Delta(\vec{k}, t) \bar{\Delta}(\vec{k}', t) \rangle = P_{\Delta}(k, t) \delta(\vec{k} - \vec{k}')$$

Gaussianity  
& Isotropy

# CMB physics: initial conditions



# CMB physics: initial conditions

End of inflation

$\Delta(\vec{x}, t)$

Harrison Zel-dovich spectrum:

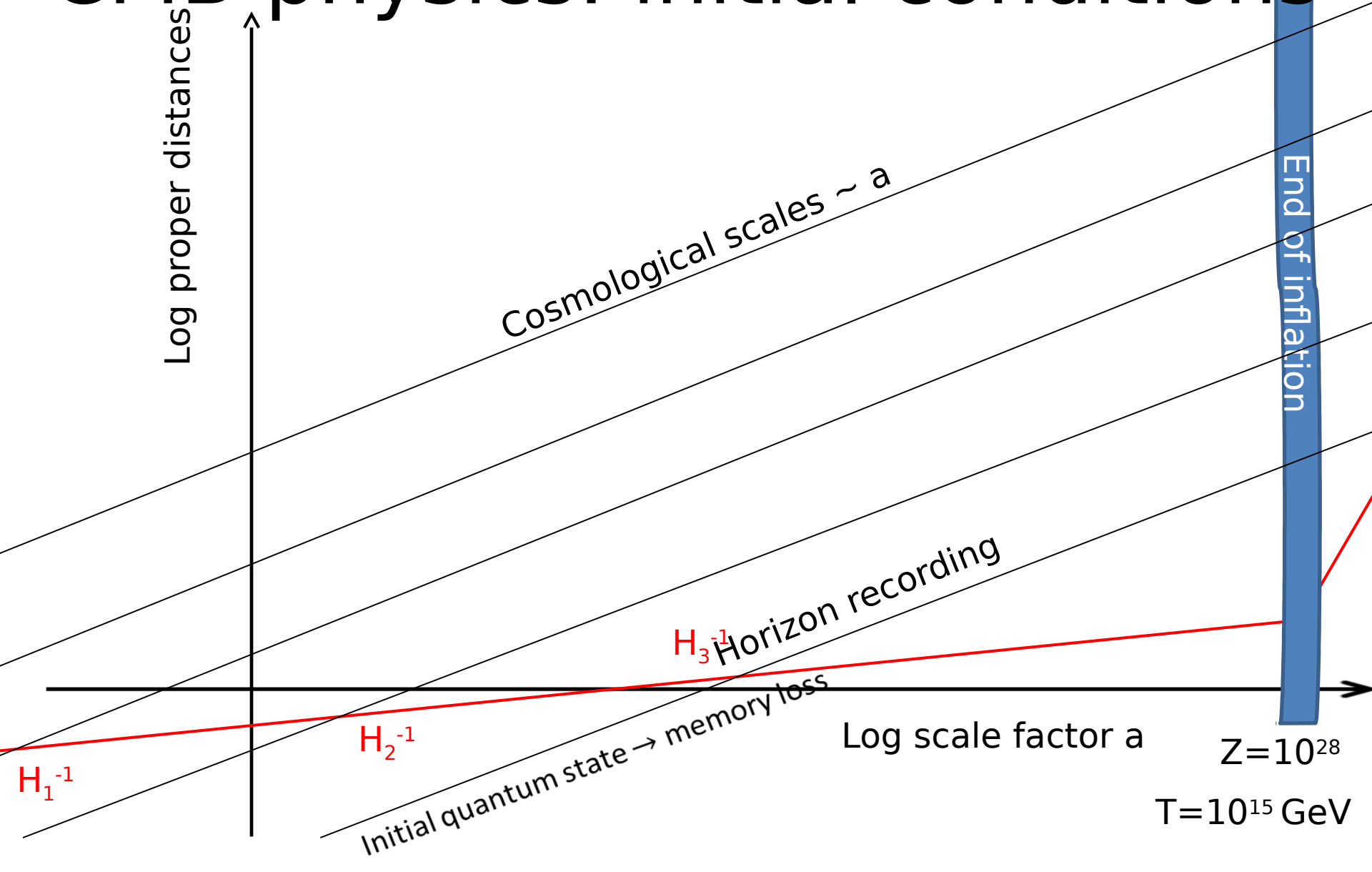
$$\langle |\Delta(\vec{x}, t)|^2 \rangle = \frac{1}{2\pi^2} \int \frac{dk}{k} k^3 P_{\Delta}(k, t)$$

Harrison Zel'dovich spectrum, also known as scale invariant:

$$k^3 P_{\Delta}(k, t) = \text{constant at horizon crossing}, \Leftrightarrow \frac{a}{k} = H^{-1}$$

at a fixed time for all scales,  $P_{\Delta}(k) = \text{normalization} \times k$

# CMB physics: initial conditions



# CMB physics: initial conditions

End of inflation

$\Delta(\vec{x}, t)$

Harrison Zel-Dovich spectrum:

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at a fixed time for all scales,  $P_{\Delta}(k) = \text{normalization} \times k$

$$P_{\Delta}(k) = \text{normalization} \times k^{n_s}$$

Inflation predicts  $n_s < 1$

# CMB physics: initial conditions

End of inflation

$\Delta(\vec{x}, t)$

Questioning Gaussianity:

$$\langle \Delta(\vec{k}, t) \bar{\Delta}(\vec{k}', t) \rangle = P_{\Delta}(k, t) \delta(\vec{k} - \vec{k}')$$

Finite correlation:

$$\langle \Delta(\vec{k}, t) \bar{\Delta}(\vec{k}', t) \rangle = P_{\Delta}(k, t) e^{-(\vec{k} - \vec{k}')^2 / 2\sigma}$$

Usually expressed in terms of second order correction:

$$\Phi = \Phi(\Phi_G) \approx \Phi_G + f_{NL} \Phi_G^2$$

# CMB physics: initial conditions

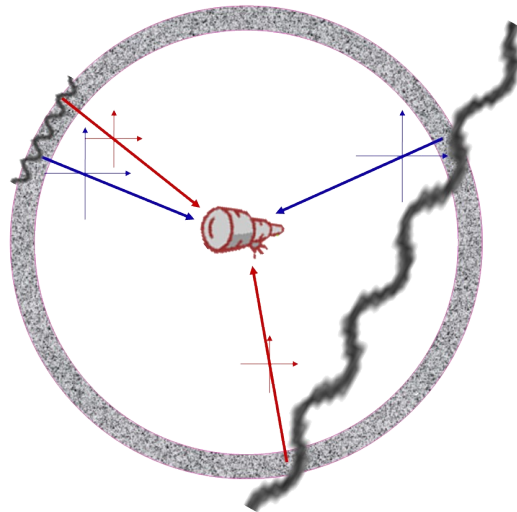
End of inflation

Non-scalar metric modes:

Vectors decay

Scalars and tensors keep constant

For total intensity, tensors they contribute on all scales:



# CMB physics: initial conditions

End of inflation

Non-scalar metric modes:

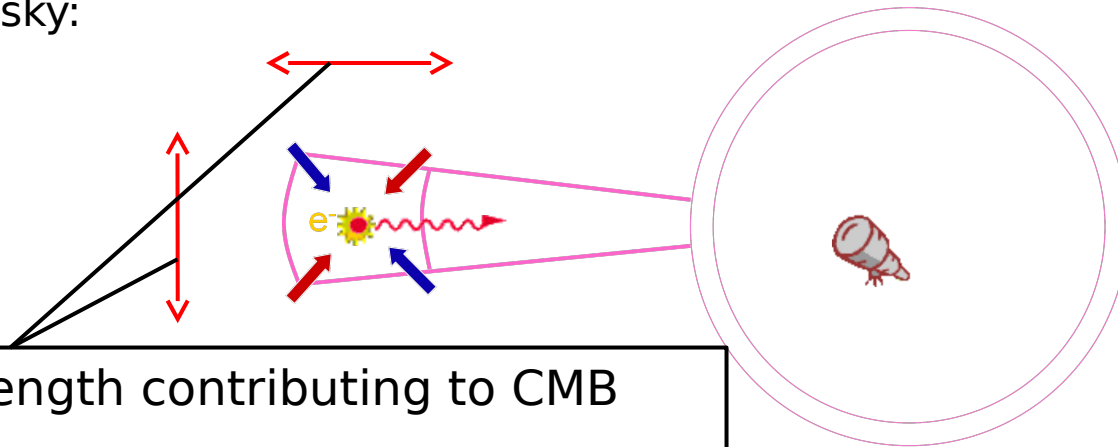
Vectors decay

Scalars and tensors keep constant

Tensors are not sub-stained by radiation pressure when sub-horizon, implying rapid diffusion damping

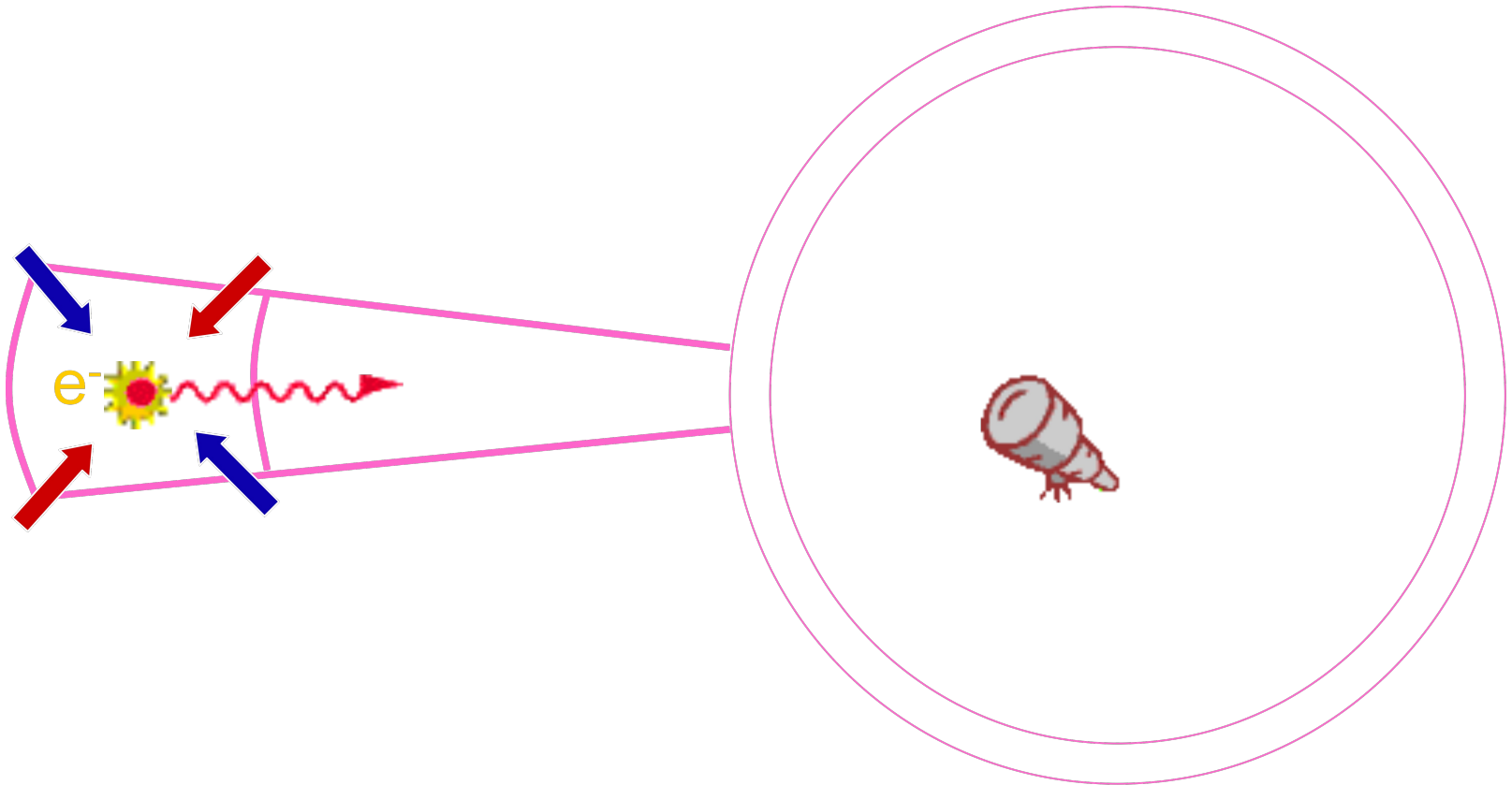
Super-horizon tensor perturbation cannot be recorded in polarisation

Polarisation CMB anisotropies from tensors peak on the angular scale corresponding to the width of the last scattering region, about 1 degree in the sky:

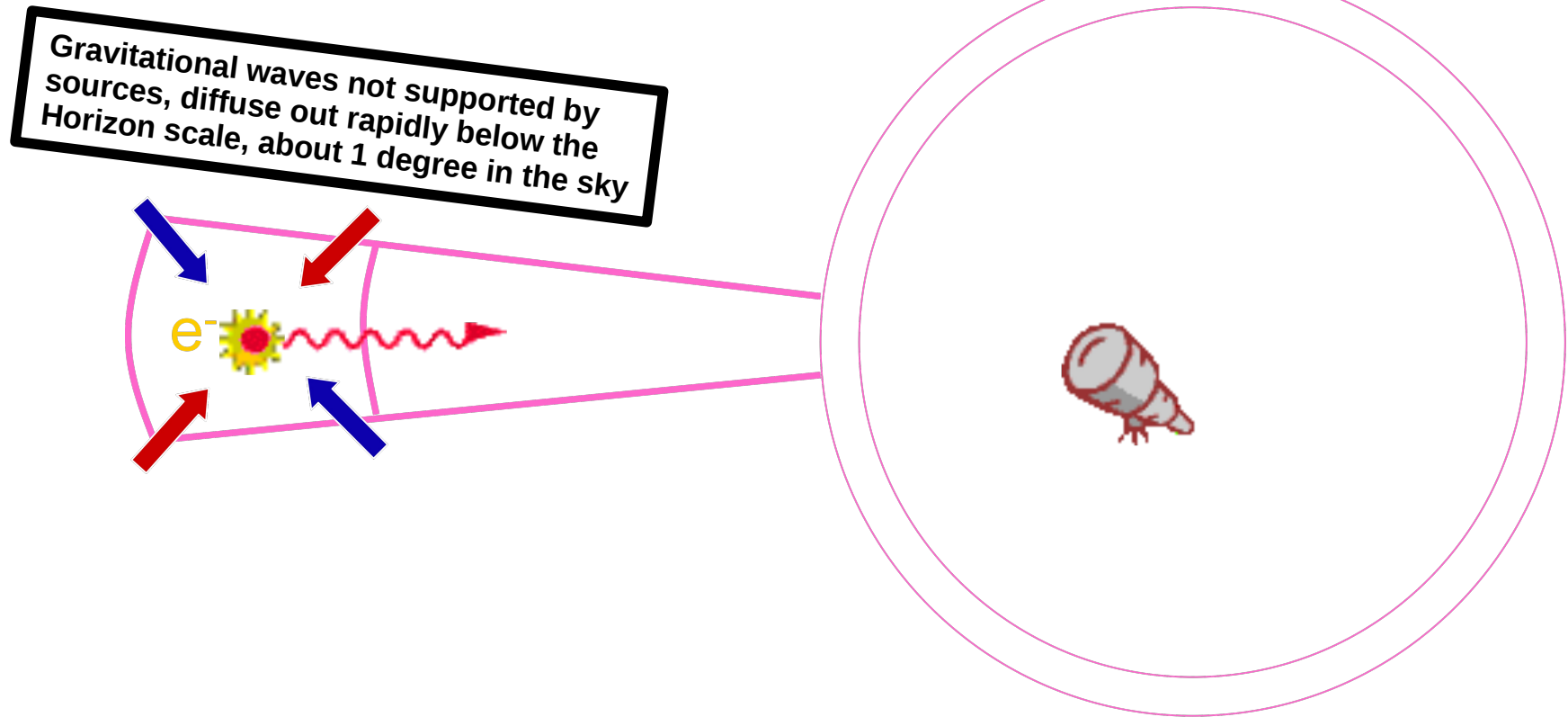


GWs wavelength contributing to CMB polarisation

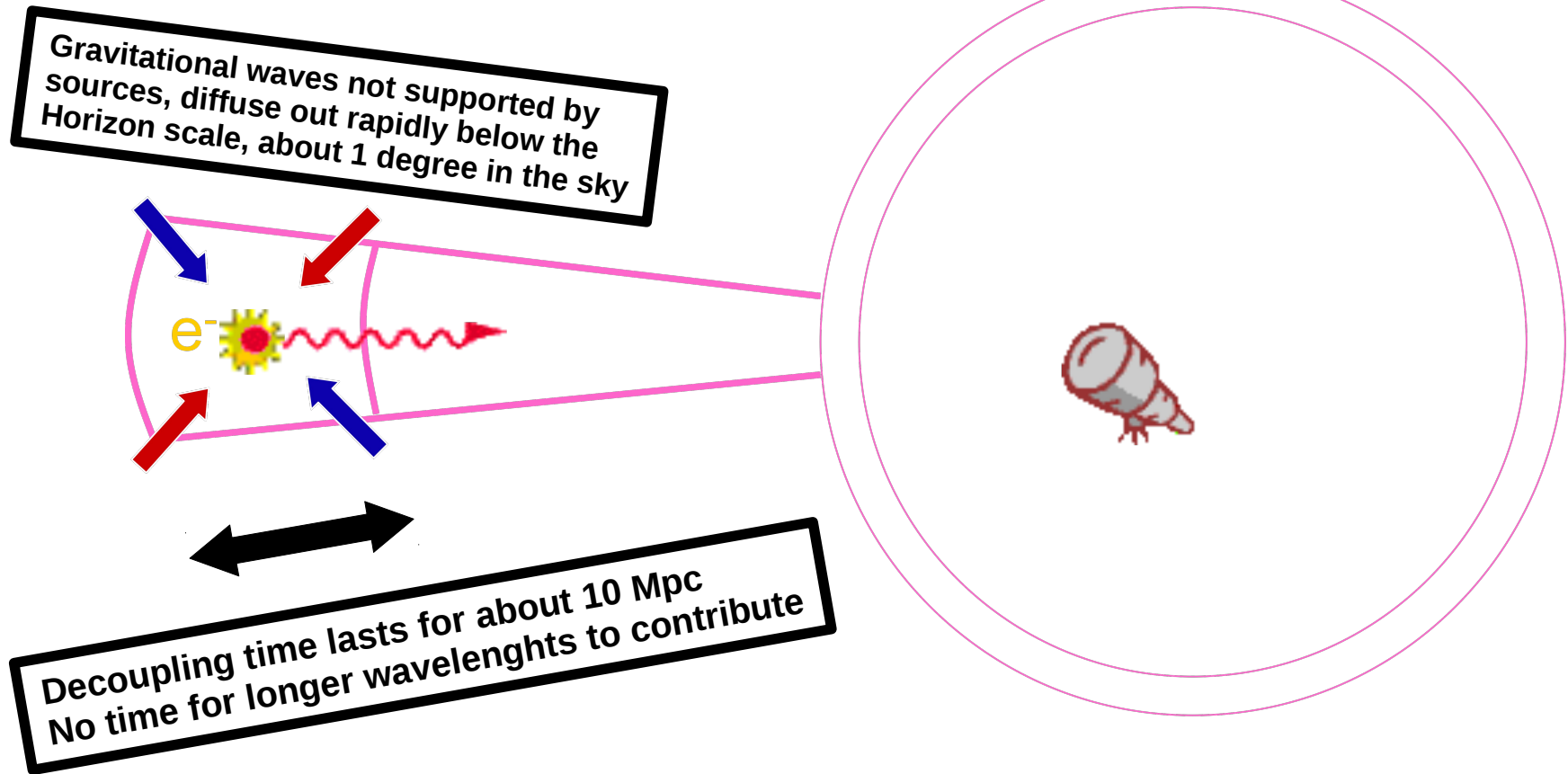
# The B-modes record GWs monochromatically in angle



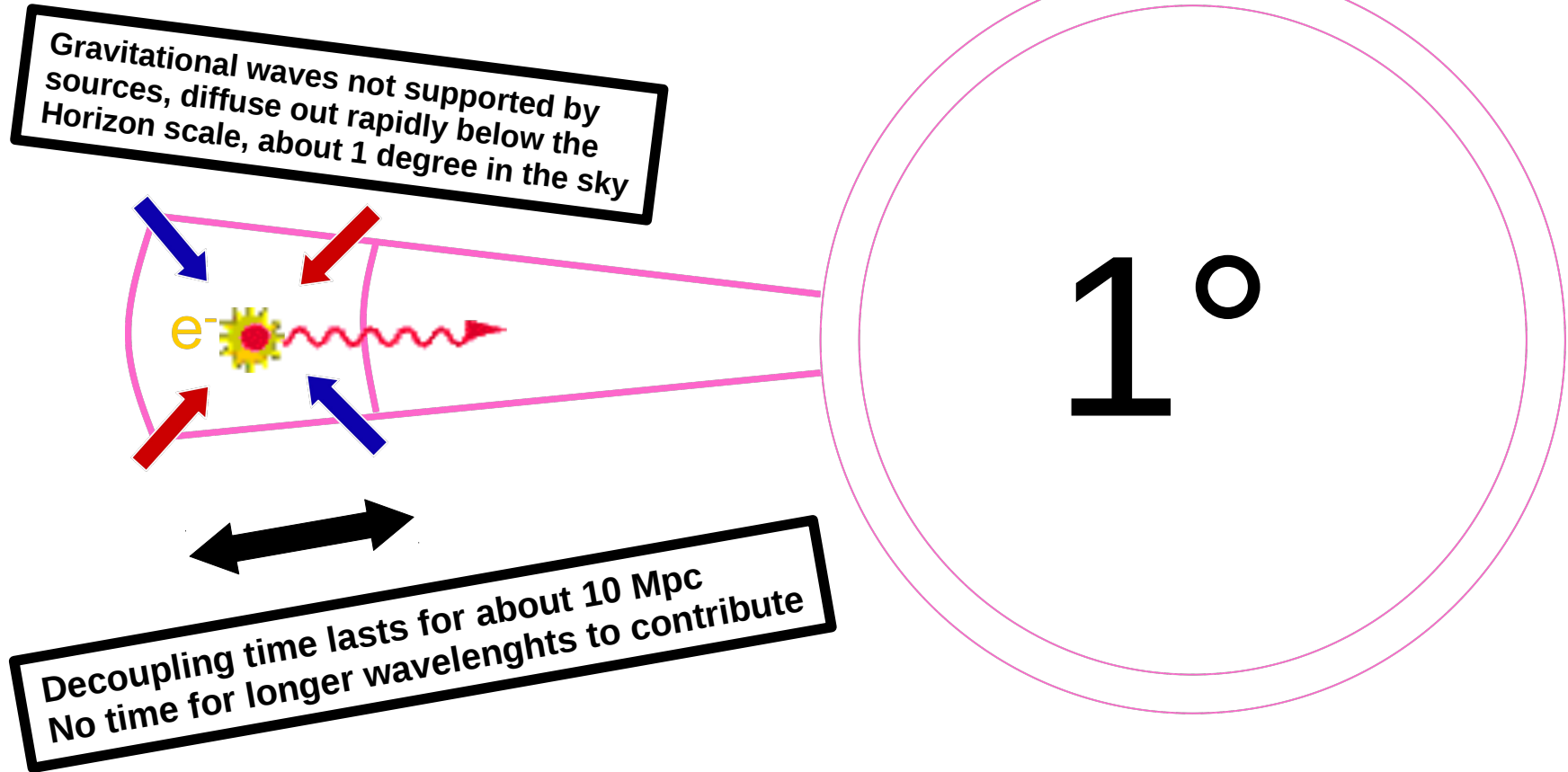
# The B-modes record GWs mono-chromatically in angle



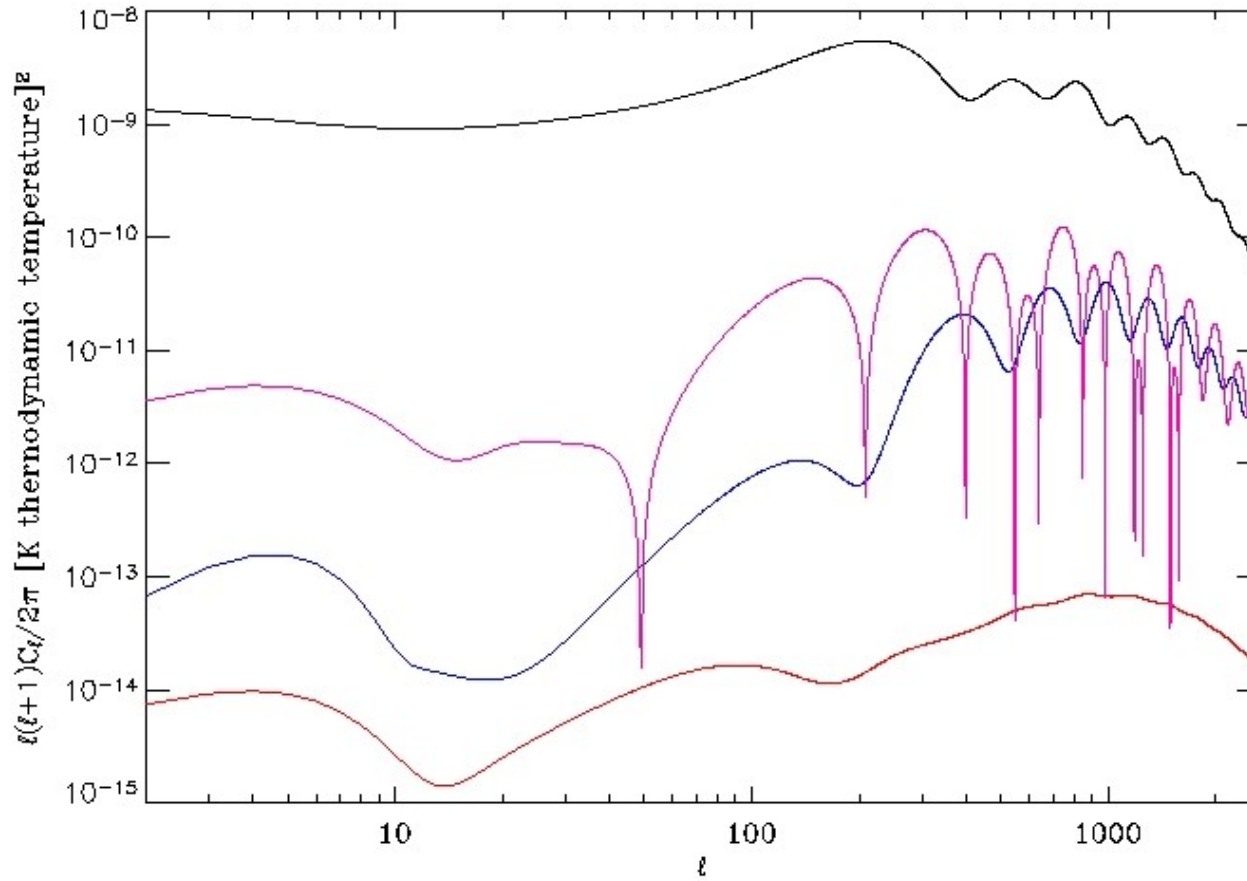
# The B-modes record GWs mono-chromatically in angle



# The B-modes record GWs mono-chromatically in angle



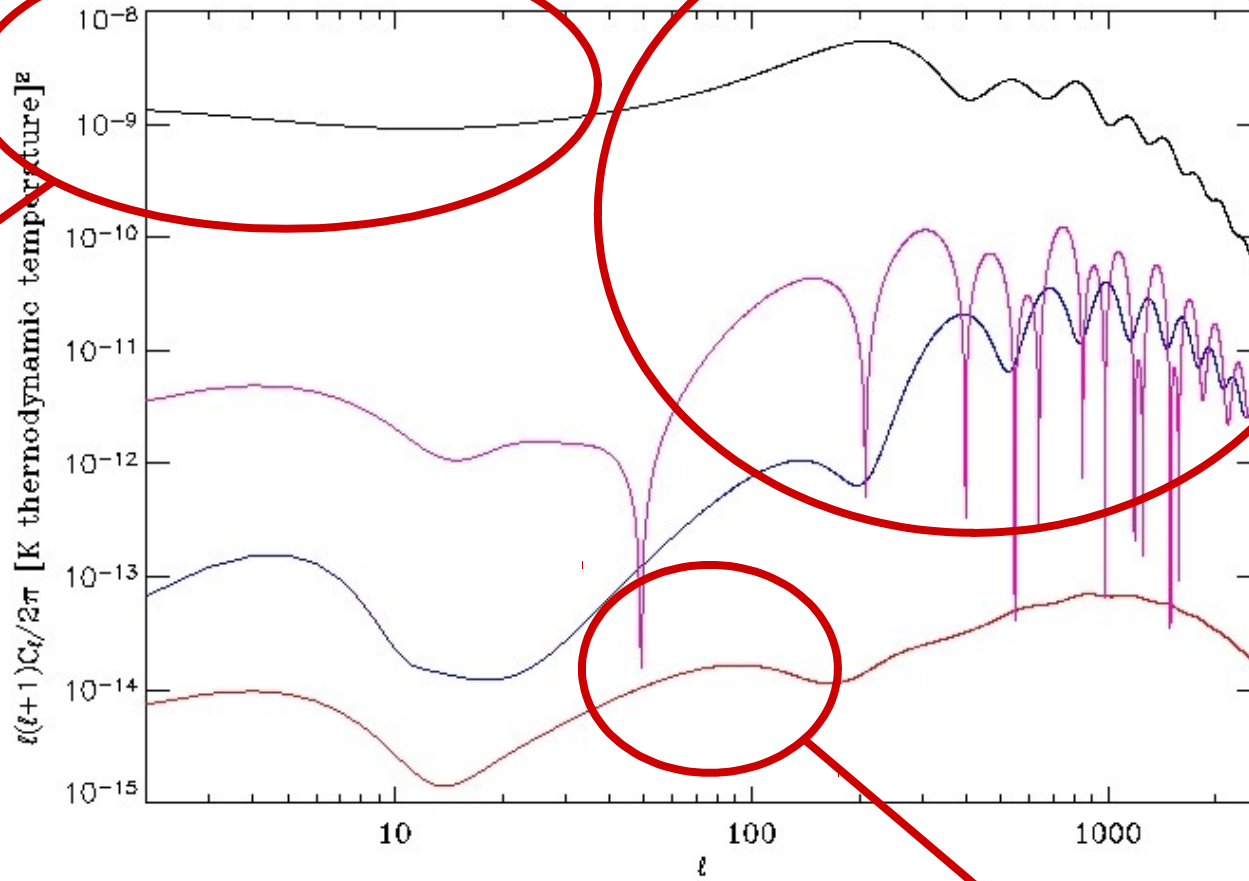
# CMB angular power spectrum



Angle  $\approx 200/\ell$  degrees

# CMB angular power spectrum

Acoustic oscillations

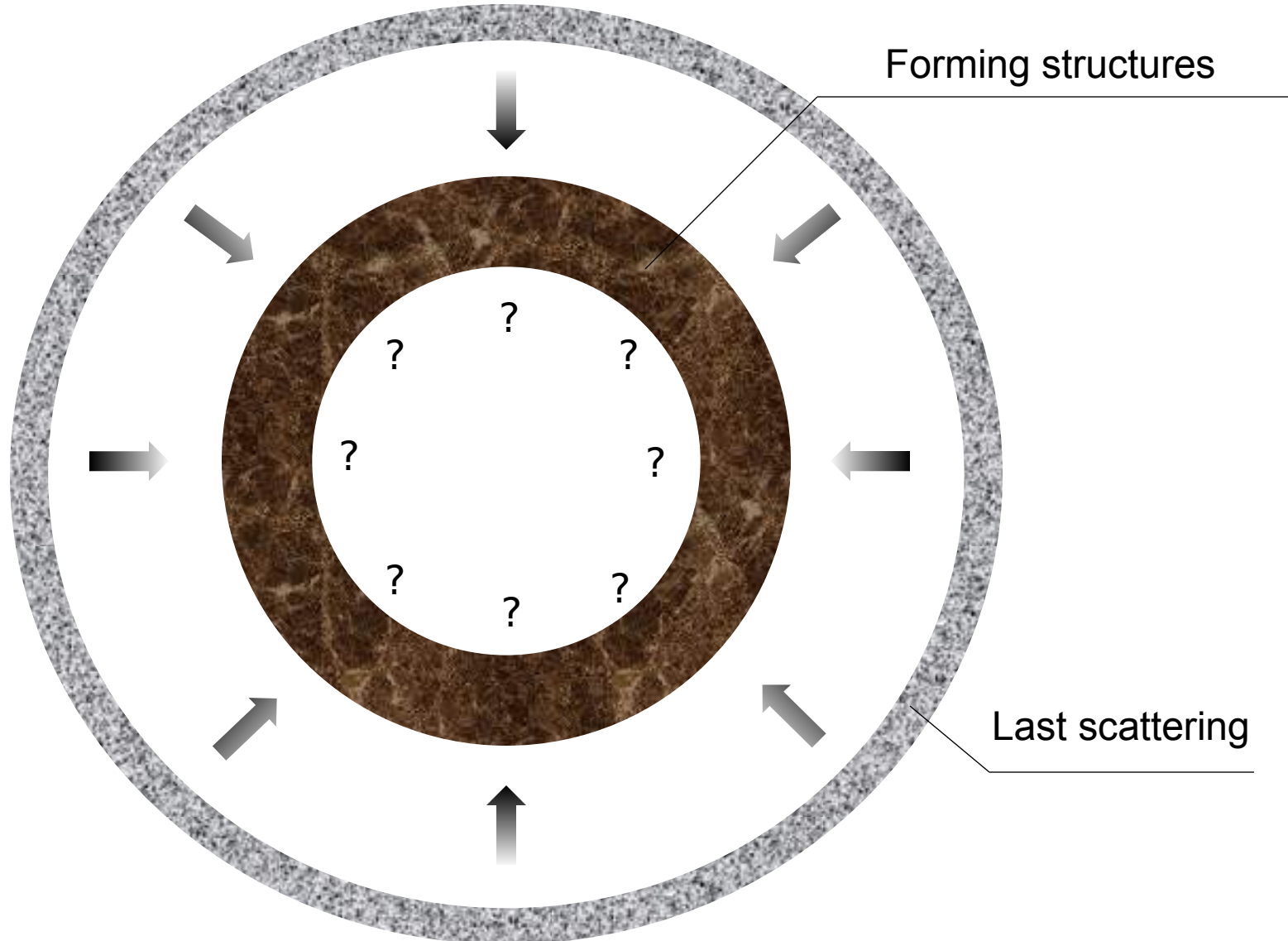


Angle  $\approx 200/\ell$  degrees

Gravitational waves

# Large Scale Structure and Dark Energy

# LSS effects on CMB



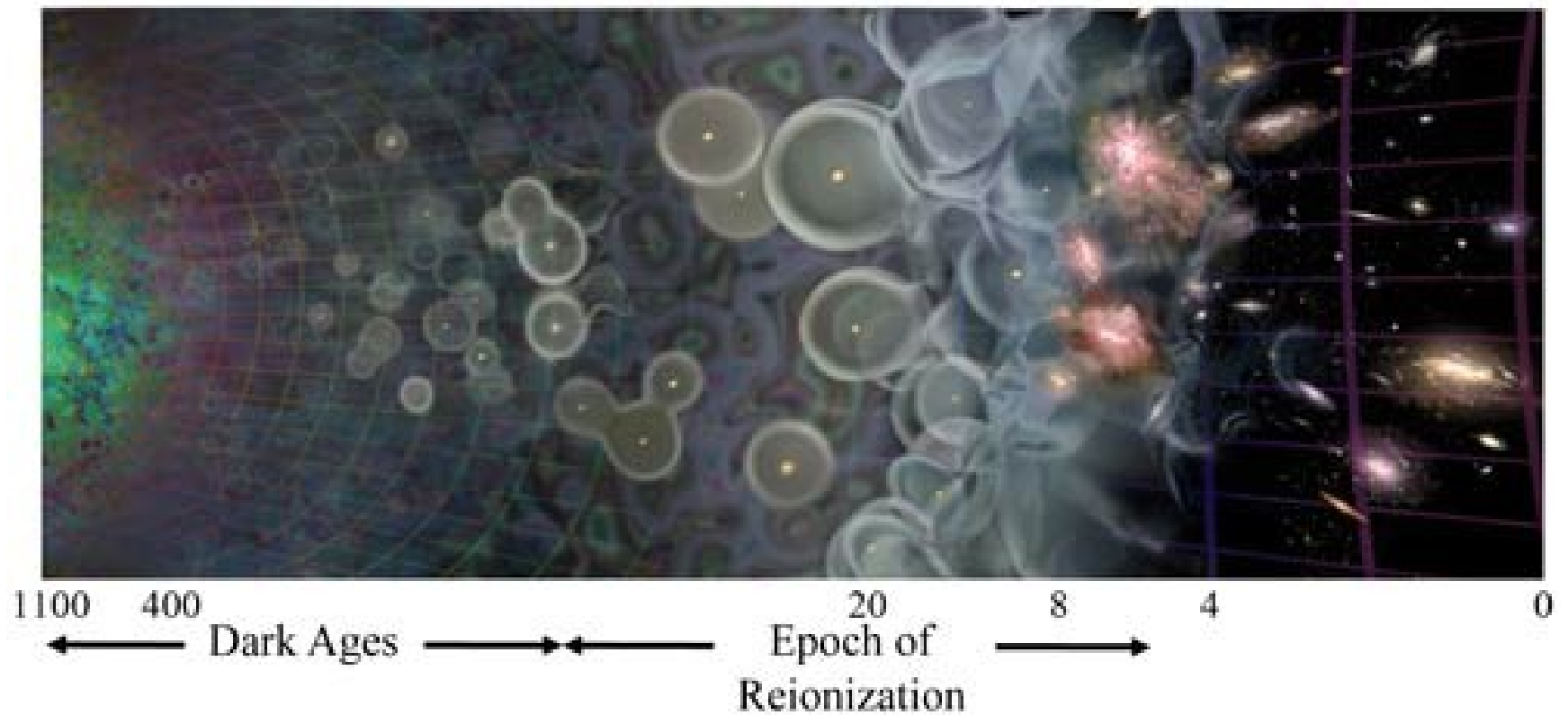
# Categories for LSS effects on CMB

- Re-scattering
- Gravitation

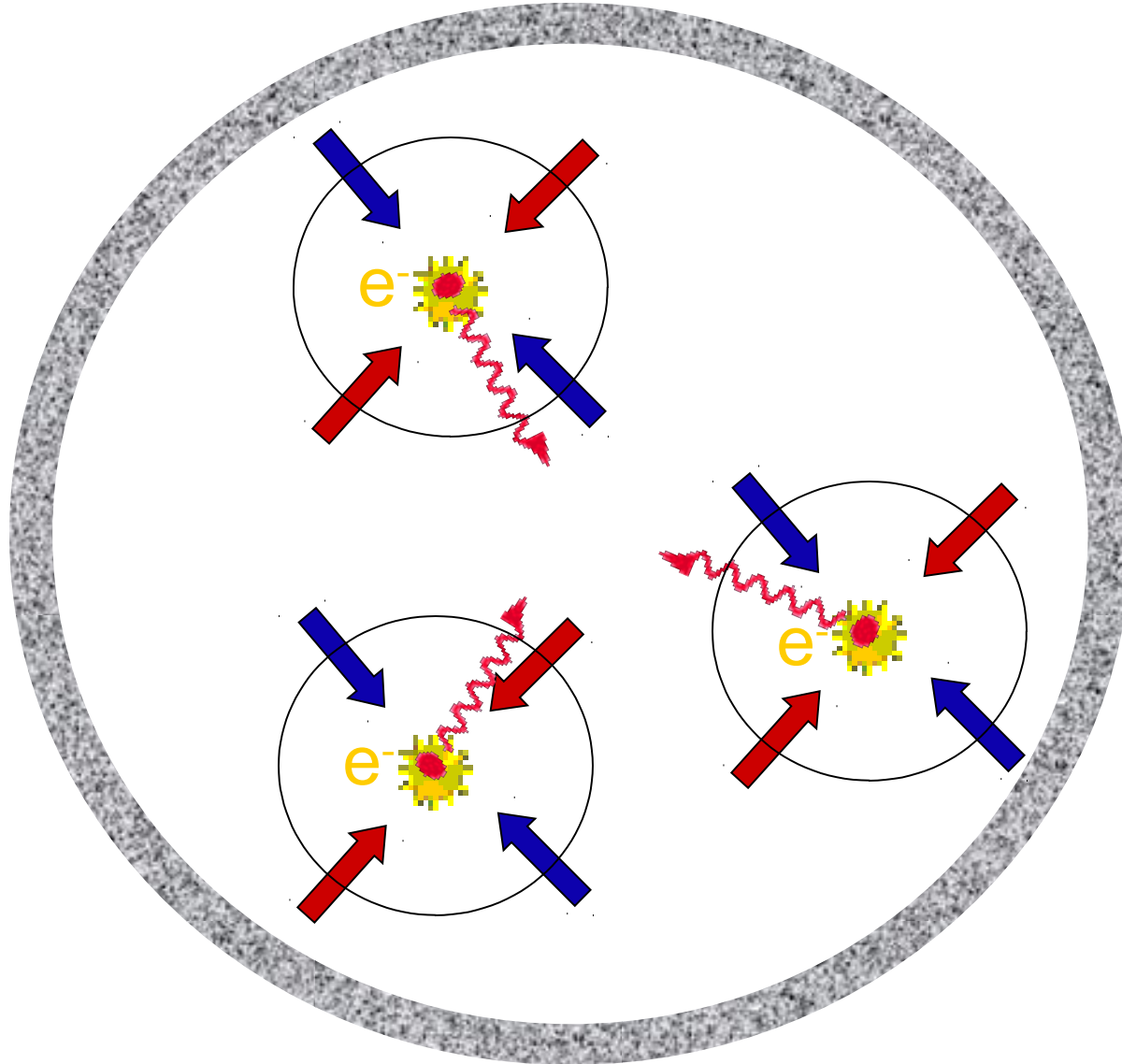
# Categories for LSS effects on CMB

- Re-scattering
  - Re-ionization
- Gravitation
  - Time evolution of the metric tensor
  - Deflection

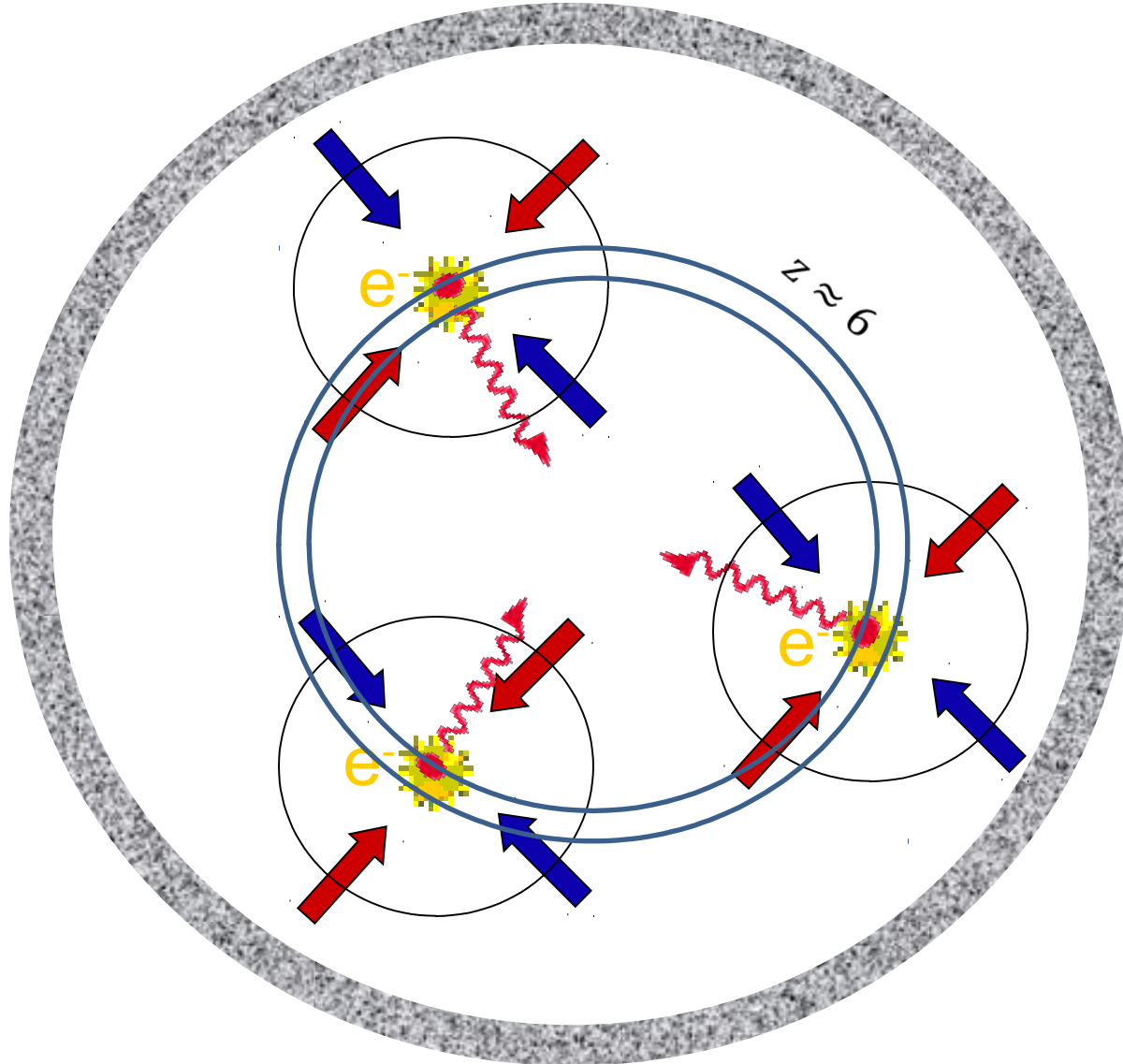
# Reionization



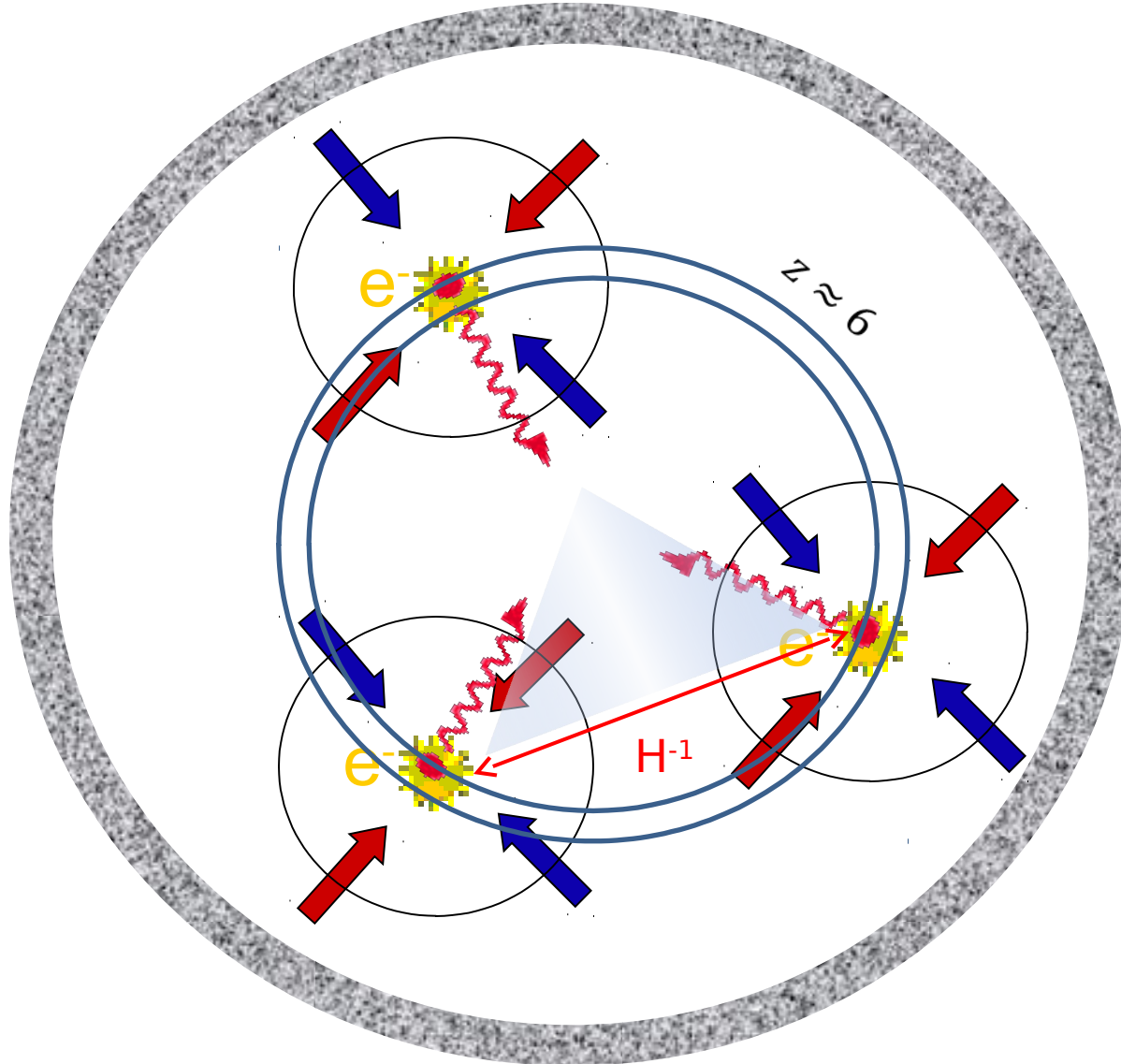
# CMB anisotropy: reionization



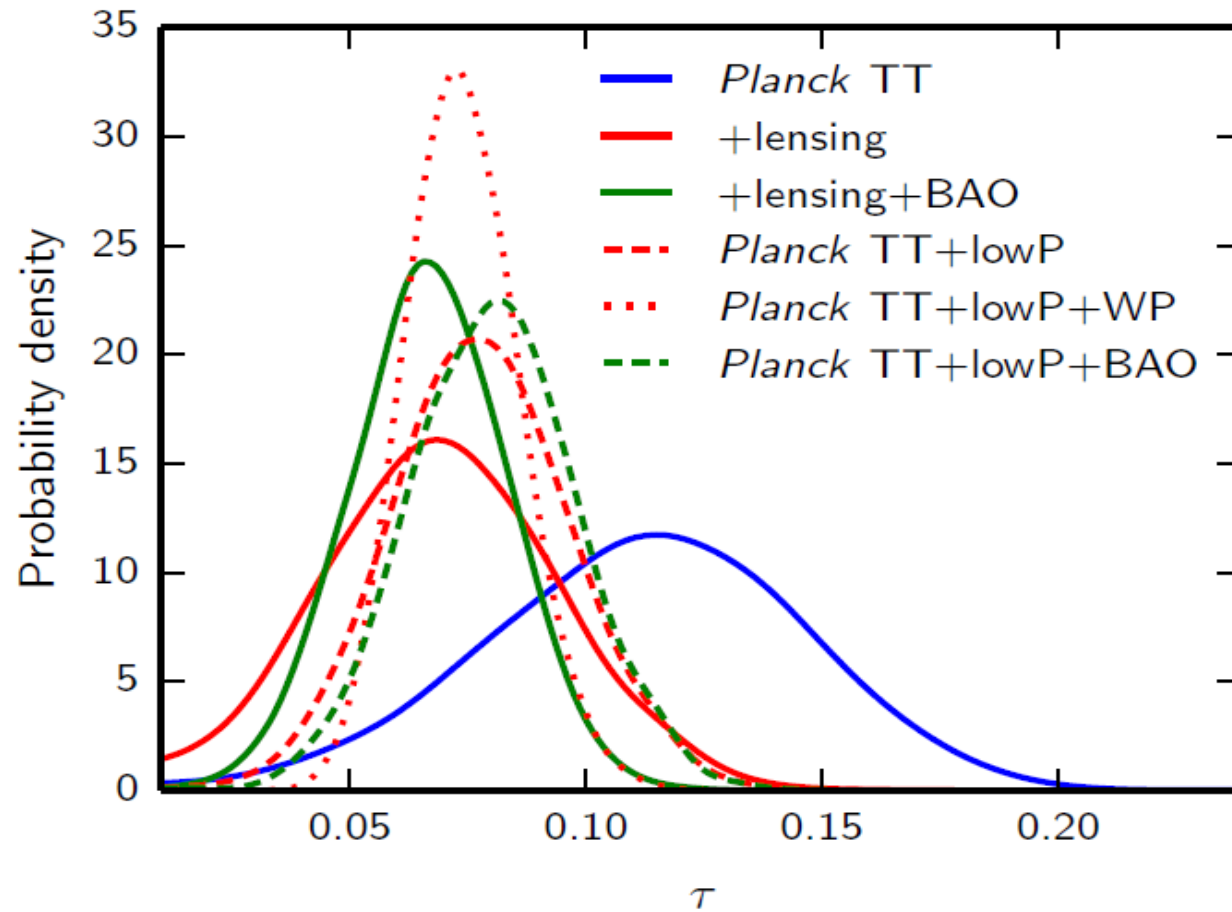
# CMB anisotropy: reionization



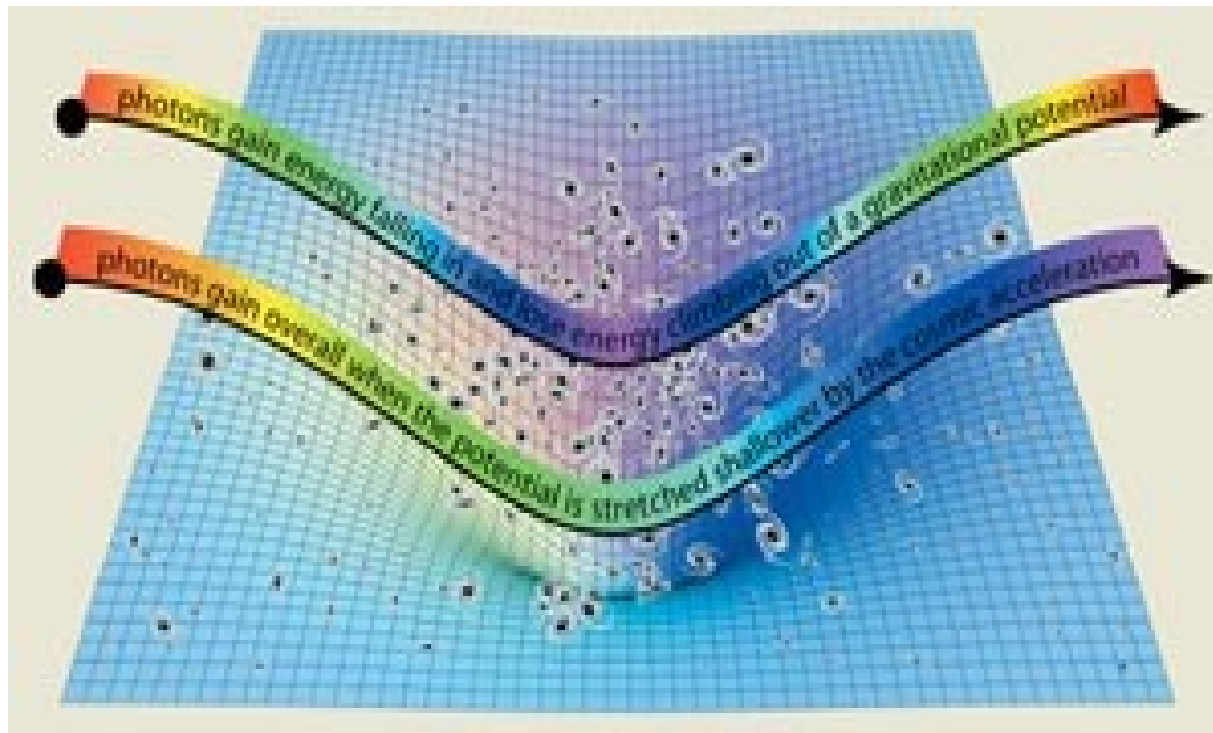
# CMB anisotropy: reionization



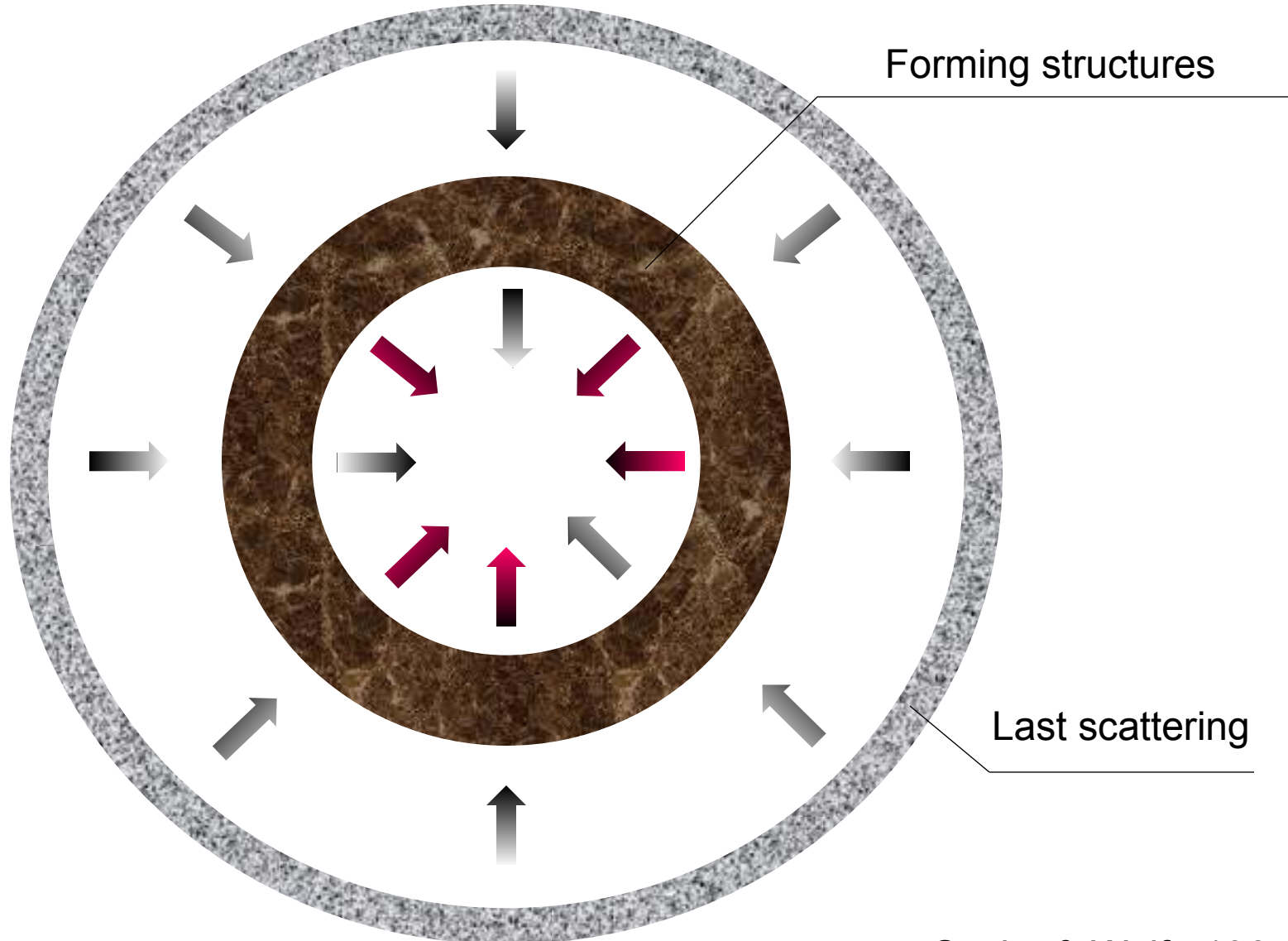
# Planck 2015: reionization epoch



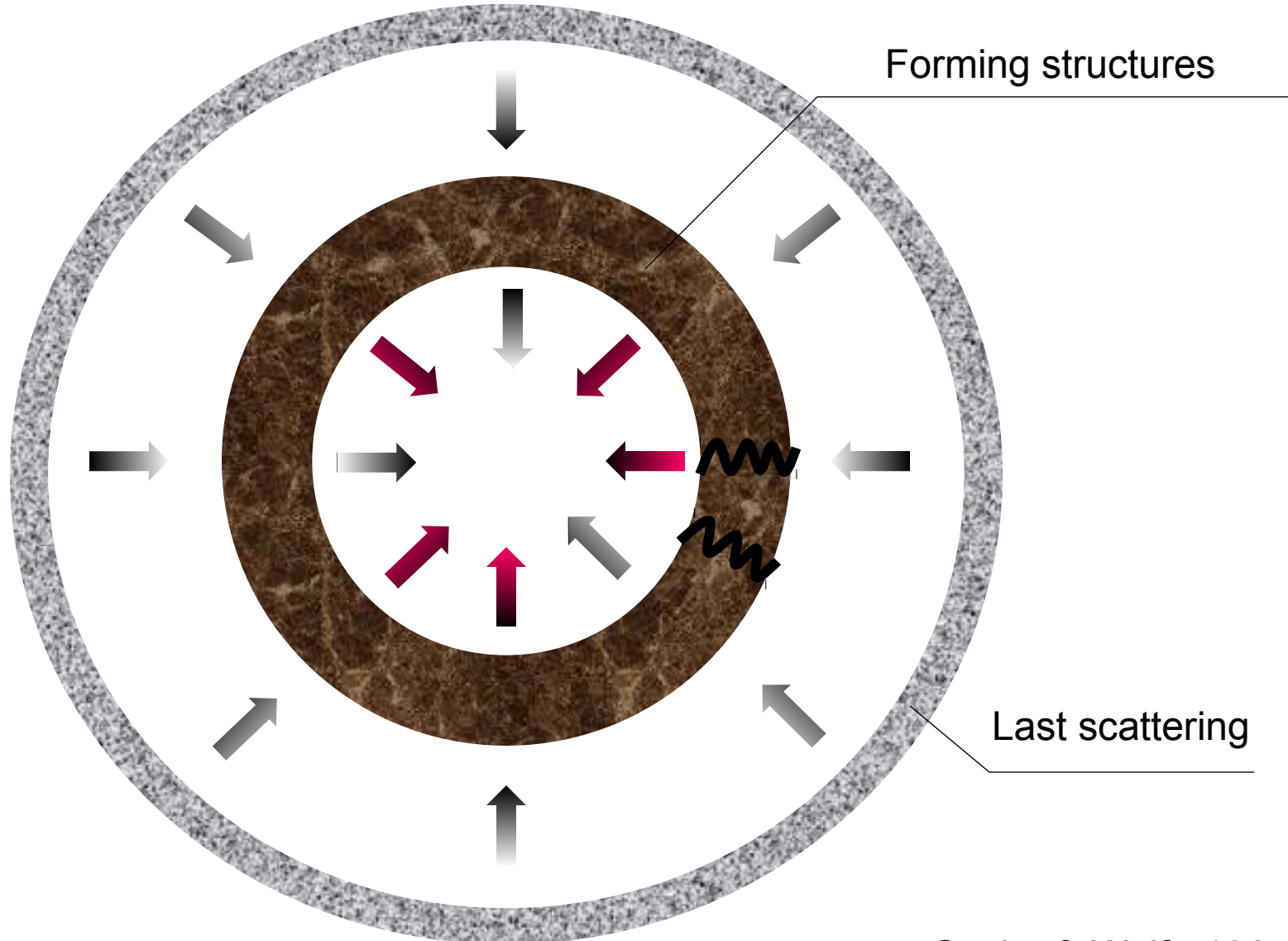
# Integrated Sachs-Wolfe



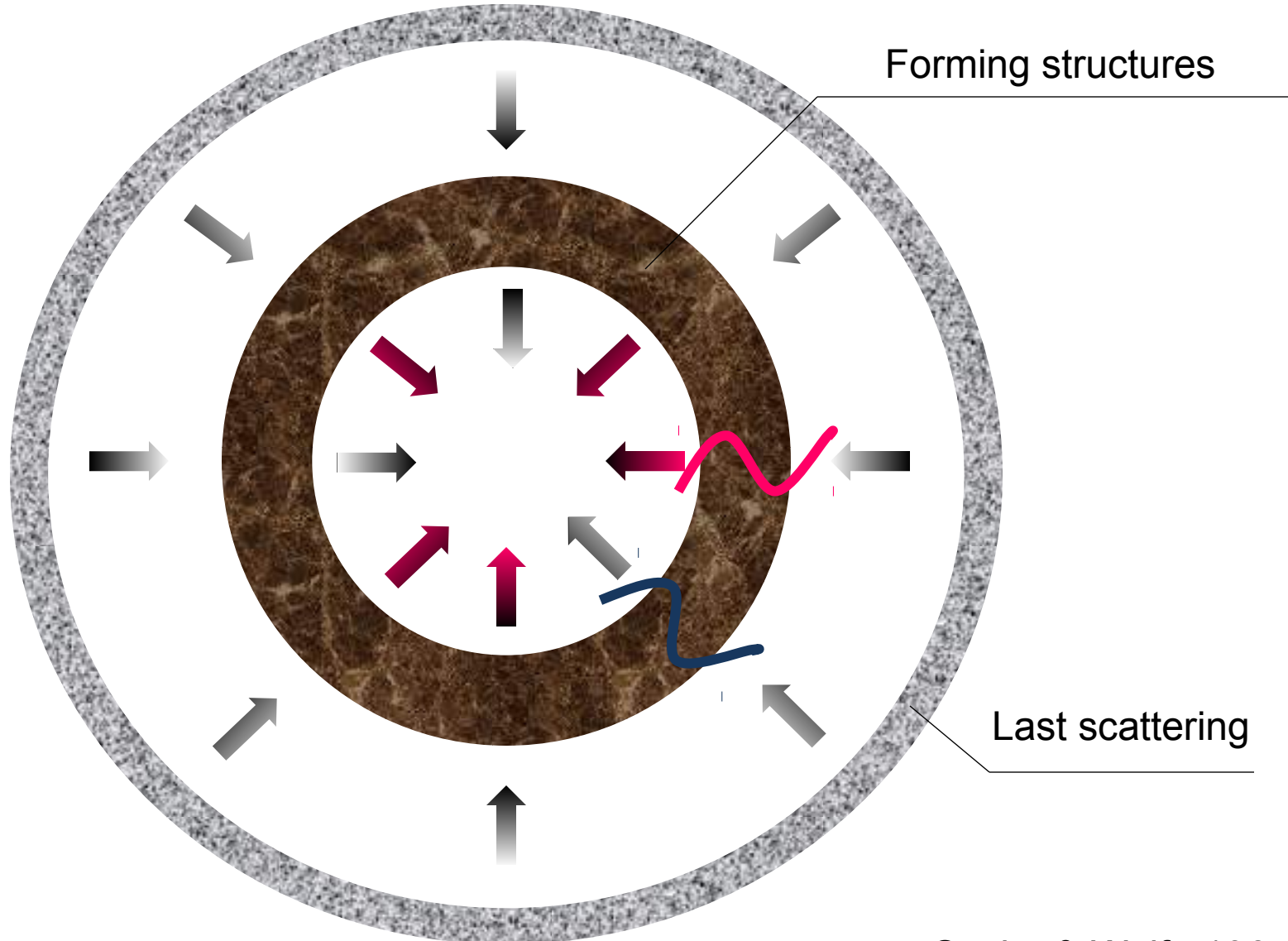
# Integrated Sachs-Wolfe



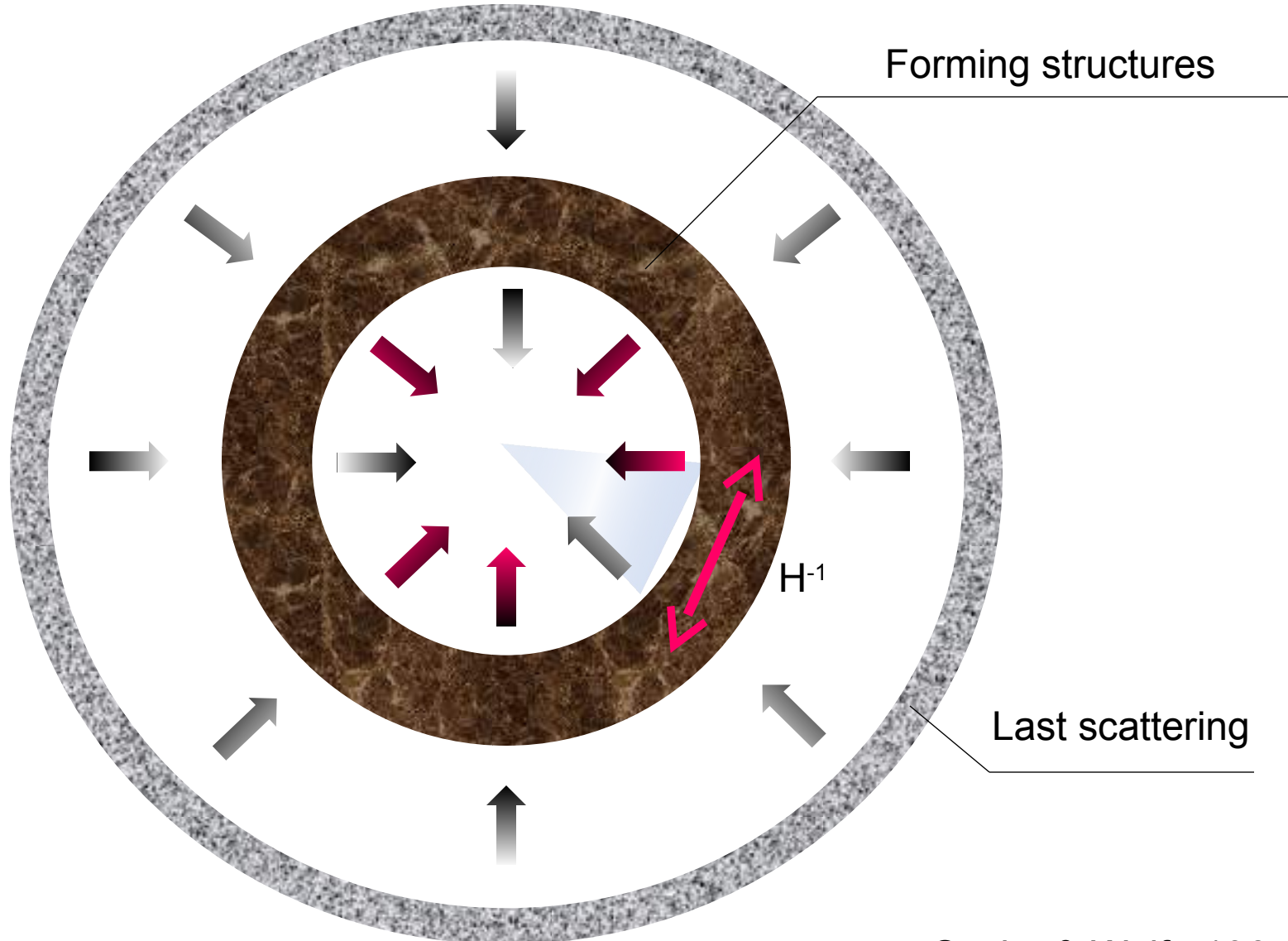
# Integrated Sachs-Wolfe



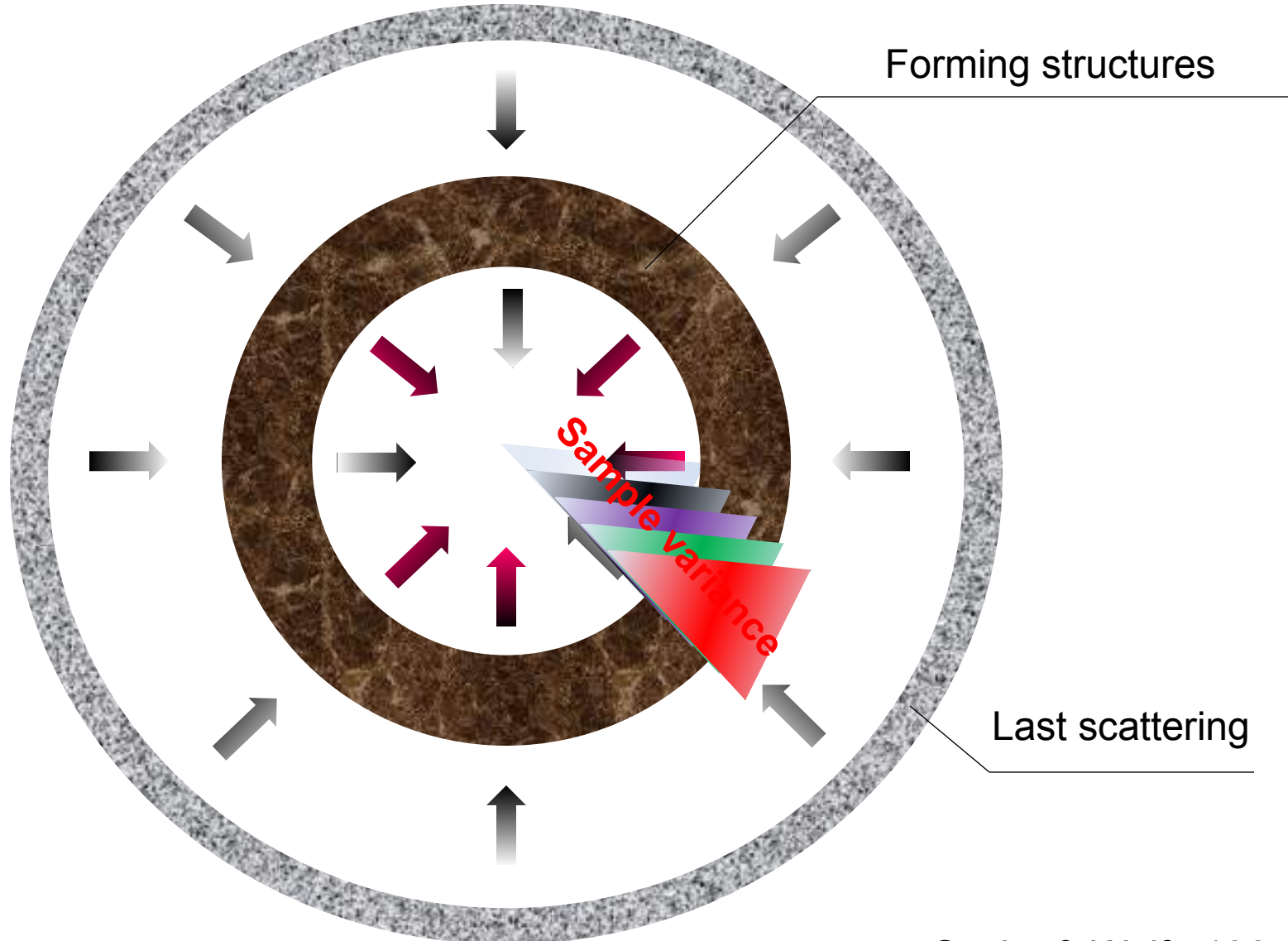
# Integrated Sachs-Wolfe



# Integrated Sachs-Wolfe



# Integrated Sachs-Wolfe



# Cross-correlating with LSS

$$\frac{\Delta T^{\text{ISW}}}{T}(\hat{n}) = -2 \int \dot{\Phi}[\eta, \hat{n}(\eta_0 - \eta)] d\eta \quad \delta_g(\hat{n}_1) = \int f(z) \delta_m(\hat{n}_1, z) dz = \int b_g(z) \frac{dN}{dz}(z) \delta_m(\hat{n}_1, z) dz$$

$$C^{\kappa\text{T}}(\theta) \equiv \left\langle \frac{\Delta T}{T}(\hat{n}_1) \delta_g(\hat{n}_2) \right\rangle = \sum_{\ell=2}^{\infty} \frac{2\ell+1}{4\pi} C_{\ell}^{\kappa\text{T}} P_{\ell}[\cos(\theta)] \exp[-0.5(\theta_s \ell)^2]$$

$$C_{\ell}^{\kappa\text{T}} = \frac{2}{\pi} \int k^2 dk P(k) I_{\ell}^{\text{ISW}}(k) I_{\ell}^g(k)$$

$$I_{\ell}^{\text{ISW}}(k) = -2 \int \frac{d\Phi(k)}{dz} j_{\ell}[k\chi(z)] dz \quad I_{\ell}^g(k) = \int b_g(z) \frac{dN}{dz}(z) \delta_m(k, z) j_{\ell}[k\chi(z)] dz$$

# Cross-correlating with LSS

$$\frac{\Delta T^{\text{ISW}}}{T}(\hat{n}) = -2 \int \dot{\Phi}[\eta, \hat{n}(\eta_0 - \eta)] d\eta \quad \delta_g(\hat{n}_1) = \int f(z) \delta_m(\hat{n}_1, z) dz = \int b_g(z) \frac{dN}{dz}(z) \delta_m(\hat{n}_1, z) dz$$

$$C^{\text{gT}}(\theta) \equiv \left\langle \frac{\Delta T}{T}(\hat{n}_1) \delta_g(\hat{n}_2) \right\rangle = \sum_{\ell=2}^{\infty} \frac{2\ell+1}{4\pi} C_{\ell}^{\text{gT}} P_{\ell}[\cos(\theta)] \exp[-0.5(\theta_s \ell)^2]$$

$$C_{\ell}^{\text{gT}} = \frac{2}{\pi} \int k^2 dk P(k) I_{\ell}^{\text{ISW}}(k) I_{\ell}^{\text{g}}(k)$$

$$I_{\ell}^{\text{ISW}}(k) = -2 \int \frac{d\Phi(k)}{dz} j_{\ell}[k\chi(z)] dz \quad I_{\ell}^{\text{g}}(k) = \int b_g(z) \frac{dN}{dz}(z) \delta_m(k, z) j_{\ell}[k\chi(z)] dz$$

Perturbation **statistics and dynamics** → early Universe

# Cross-correlating with LSS

$$\frac{\Delta T^{\text{ISW}}}{T}(\hat{n}) = -2 \int \dot{\Phi}[\eta, \hat{n}(\eta_0 - \eta)] d\eta \quad \delta_g(\hat{n}_1) = \int f(z) \delta_m(\hat{n}_1, z) dz = \int b_g(z) \frac{dN}{dz}(z) \delta_m(\hat{n}_1, z) dz$$

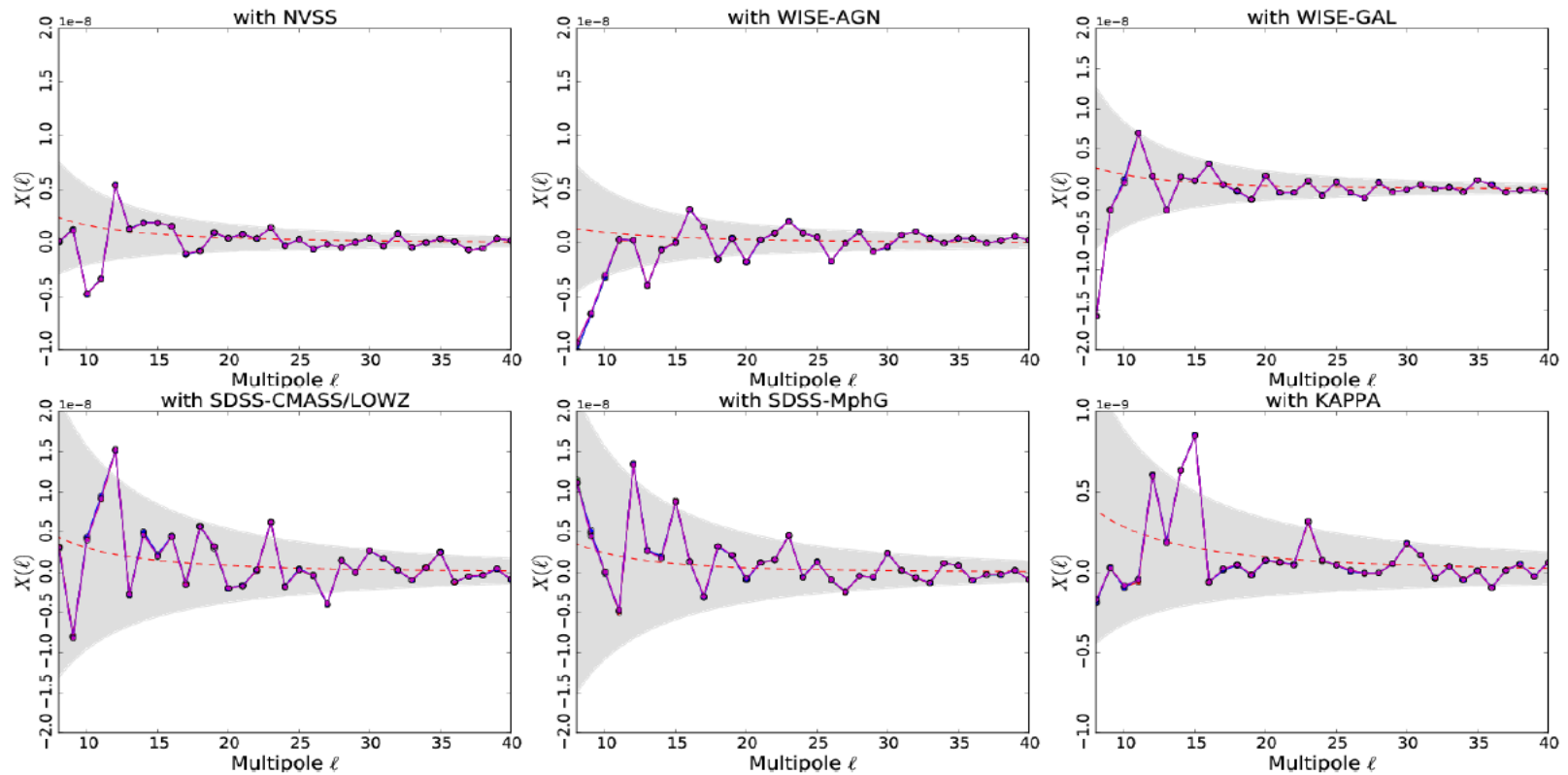
$$C^{\text{gT}}(\theta) \equiv \left\langle \frac{\Delta T}{T}(\hat{n}_1) \delta_g(\hat{n}_2) \right\rangle = \sum_{\ell=2}^{\infty} \frac{2\ell+1}{4\pi} C_{\ell}^{\text{gT}} P_{\ell}[\cos(\theta)] \exp[-0.5(\theta_s \ell)^2]$$

$$C_{\ell}^{\text{gT}} = \frac{2}{\pi} \int k^2 dk P(k) I_{\ell}^{\text{ISW}}(k) I_{\ell}^{\text{g}}(k)$$

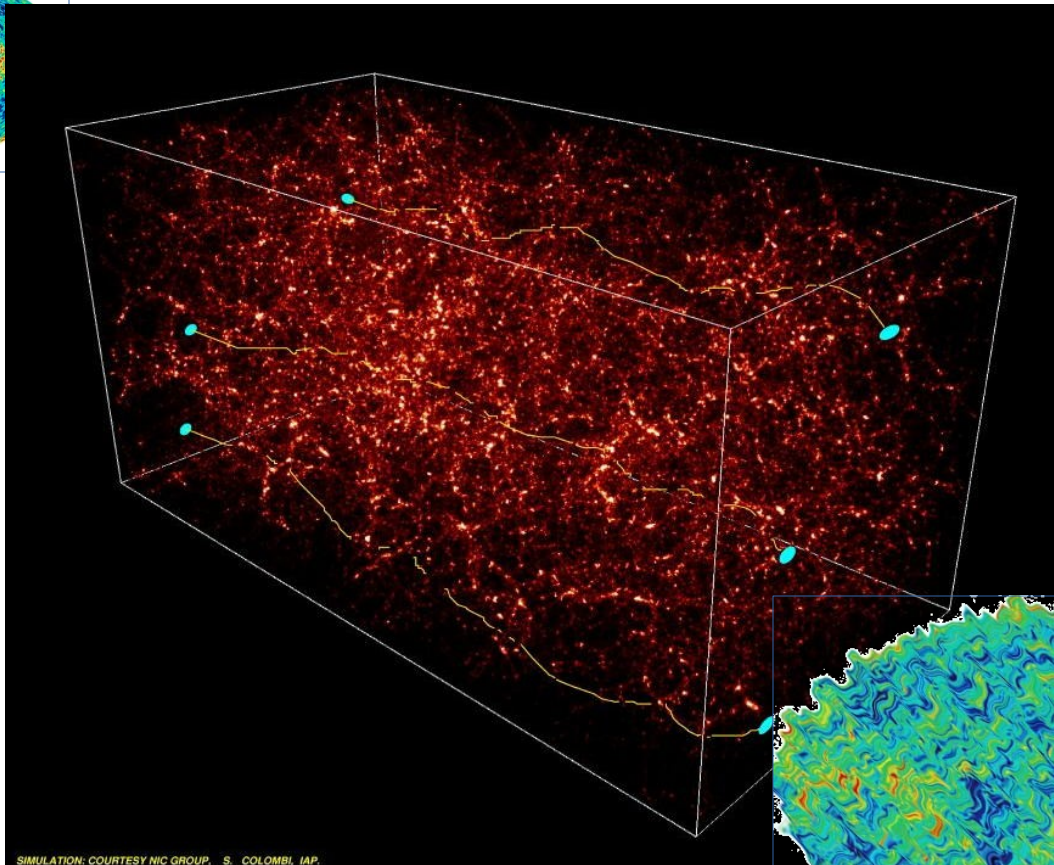
$$I_{\ell}^{\text{ISW}}(k) = -2 \int \frac{d\Phi(k)}{dz} j_{\ell}[k\chi(z)] dz \quad I_{\ell}^{\text{g}}(k) = \int b_g(z) \frac{dN}{dz}(z) \delta_m(k, z) j_{\ell}[k\chi(z)] dz$$

Geometry and **perturbation dynamics** → dark energy effects

# Planck 2015: non-zero ISW-LSS Cross-Spectra at $4\sigma$



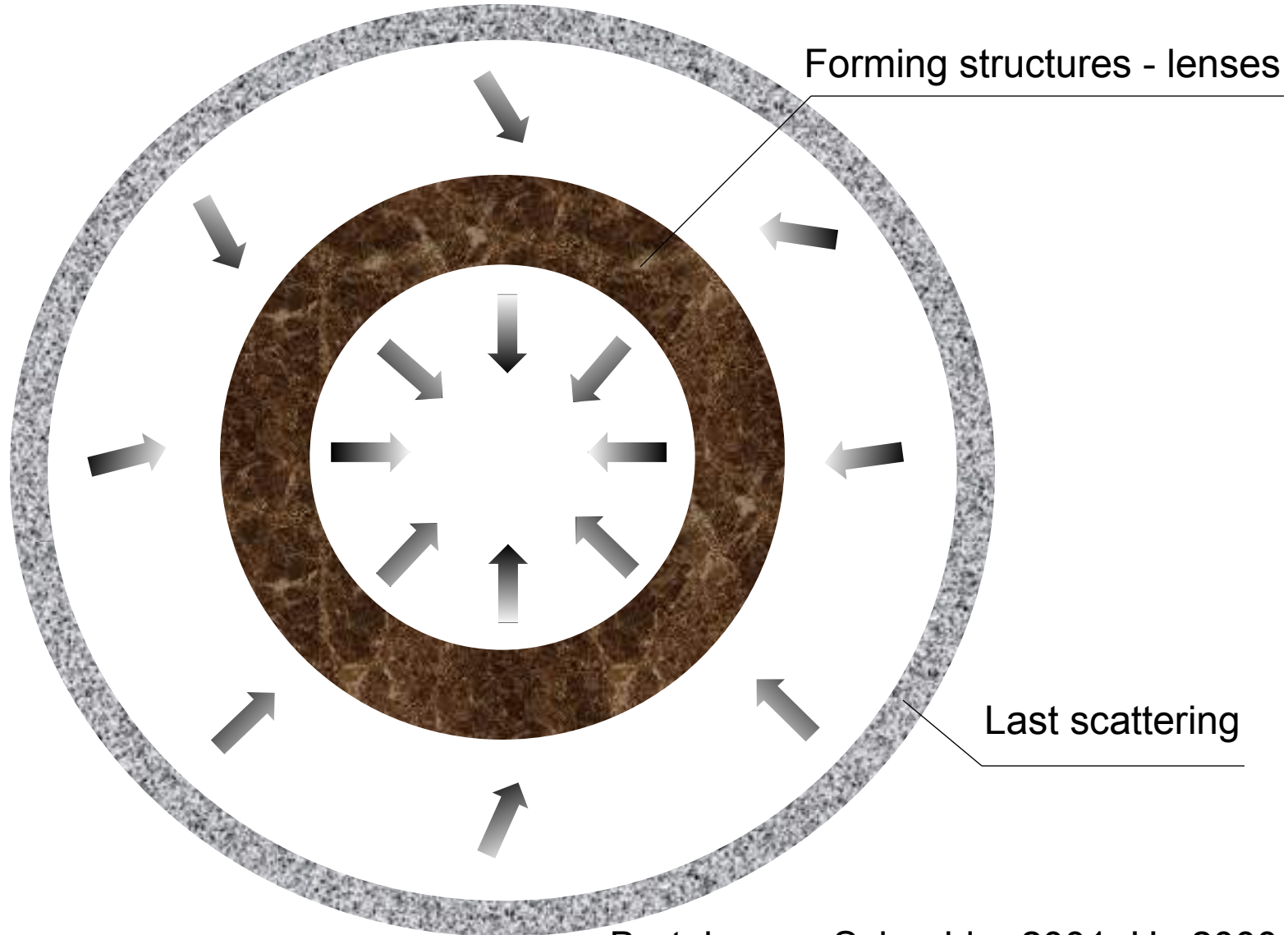
# CMB lensing



SIMULATION: COURTESY NIC GROUP, S. COLOMBI, IAP.

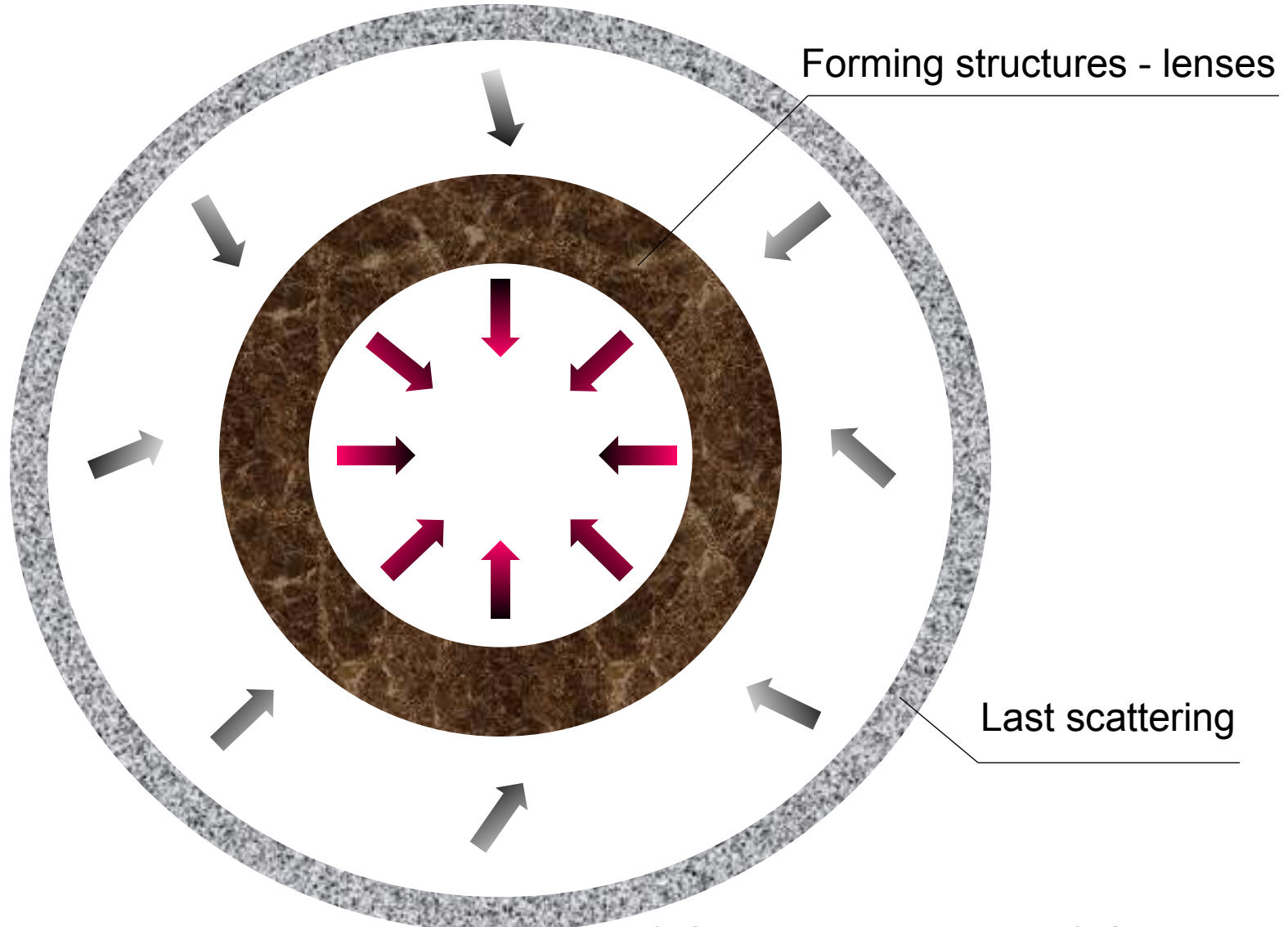
# CMB lensing

T



# CMB lensing

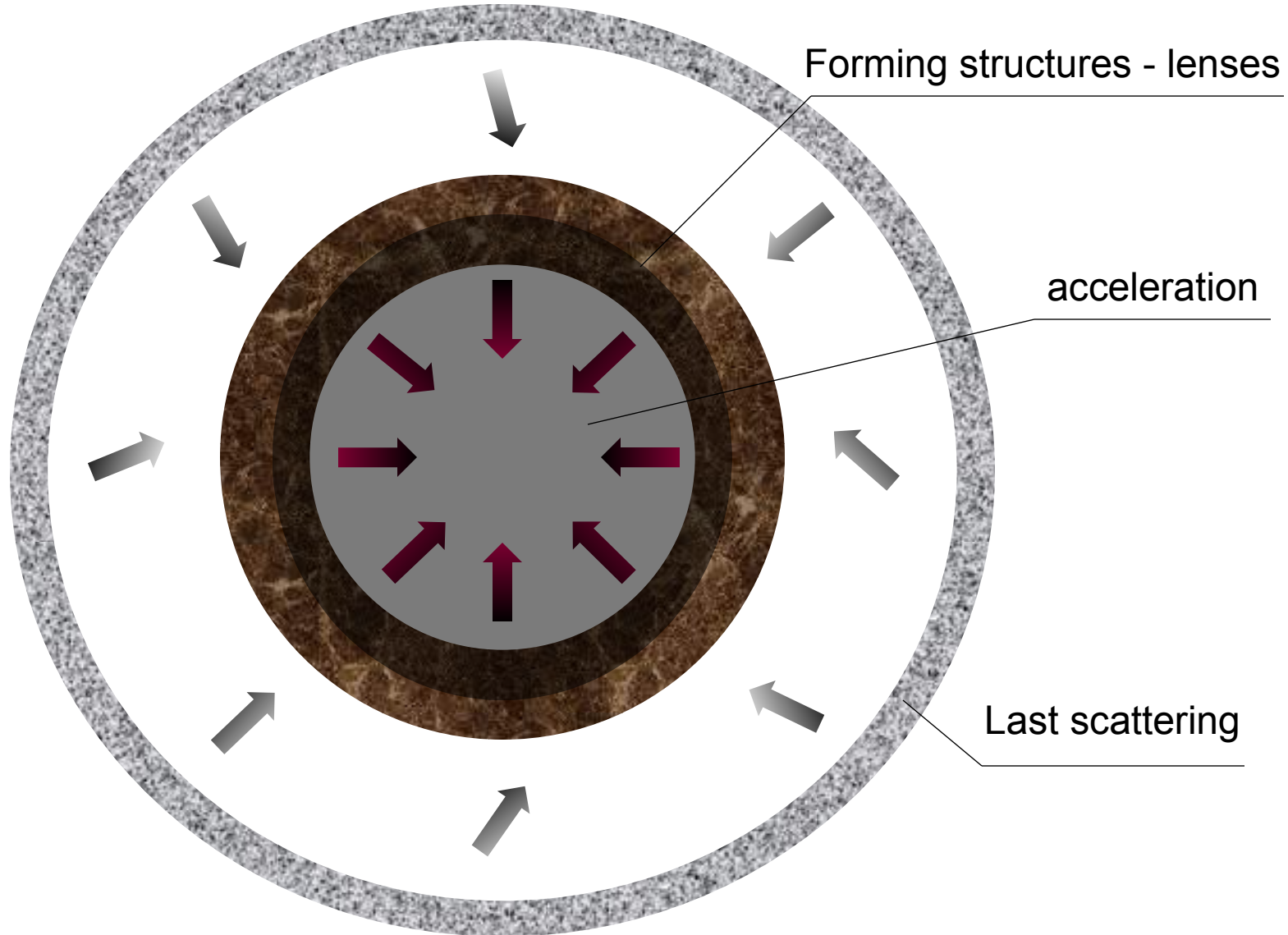
E  
B



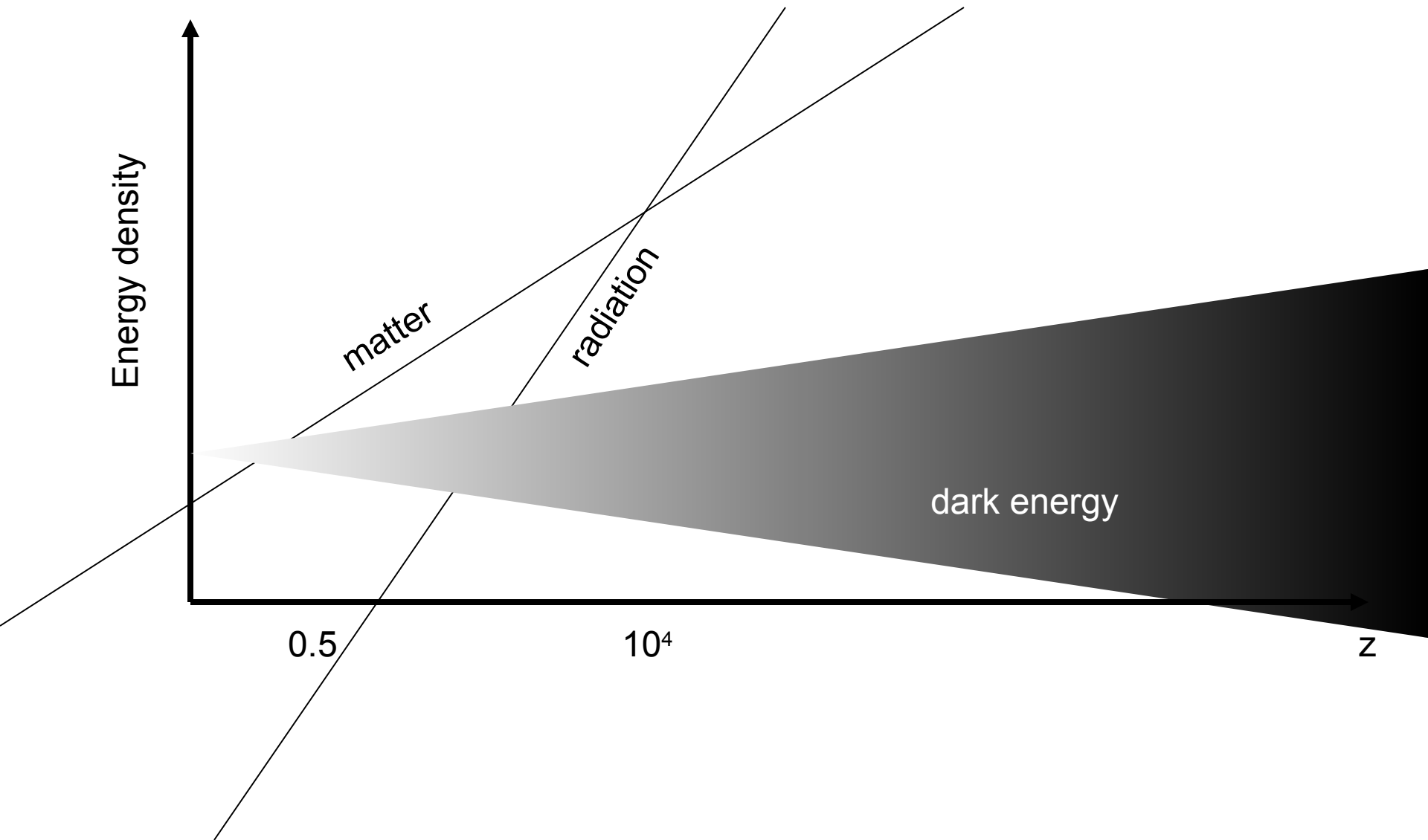
Kamionkowski, Kosowsky & Stebbins, Zaldarriaga & Seljak 1998

# CMB lensing

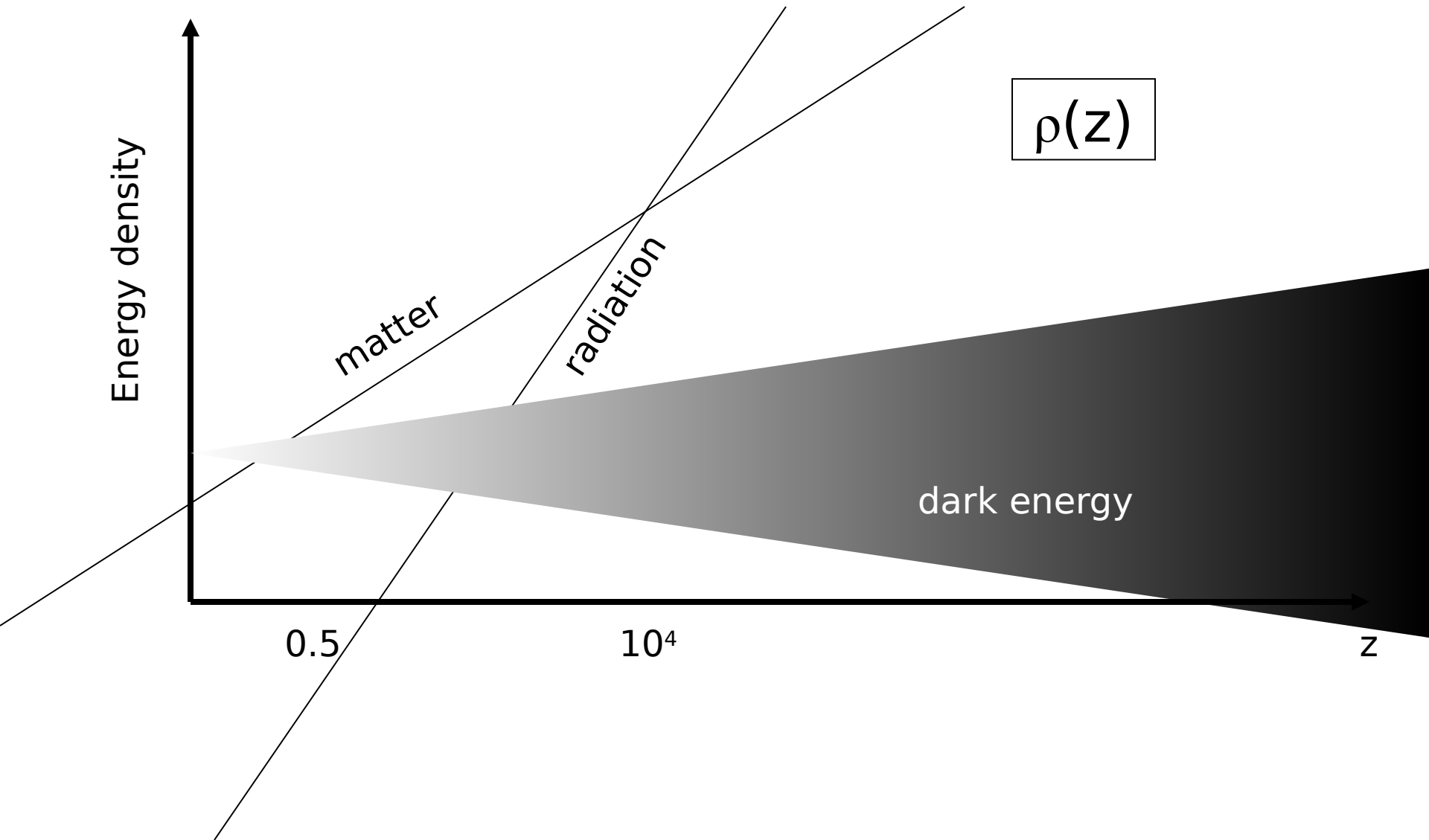
$T$   
 $E$   
 $B$



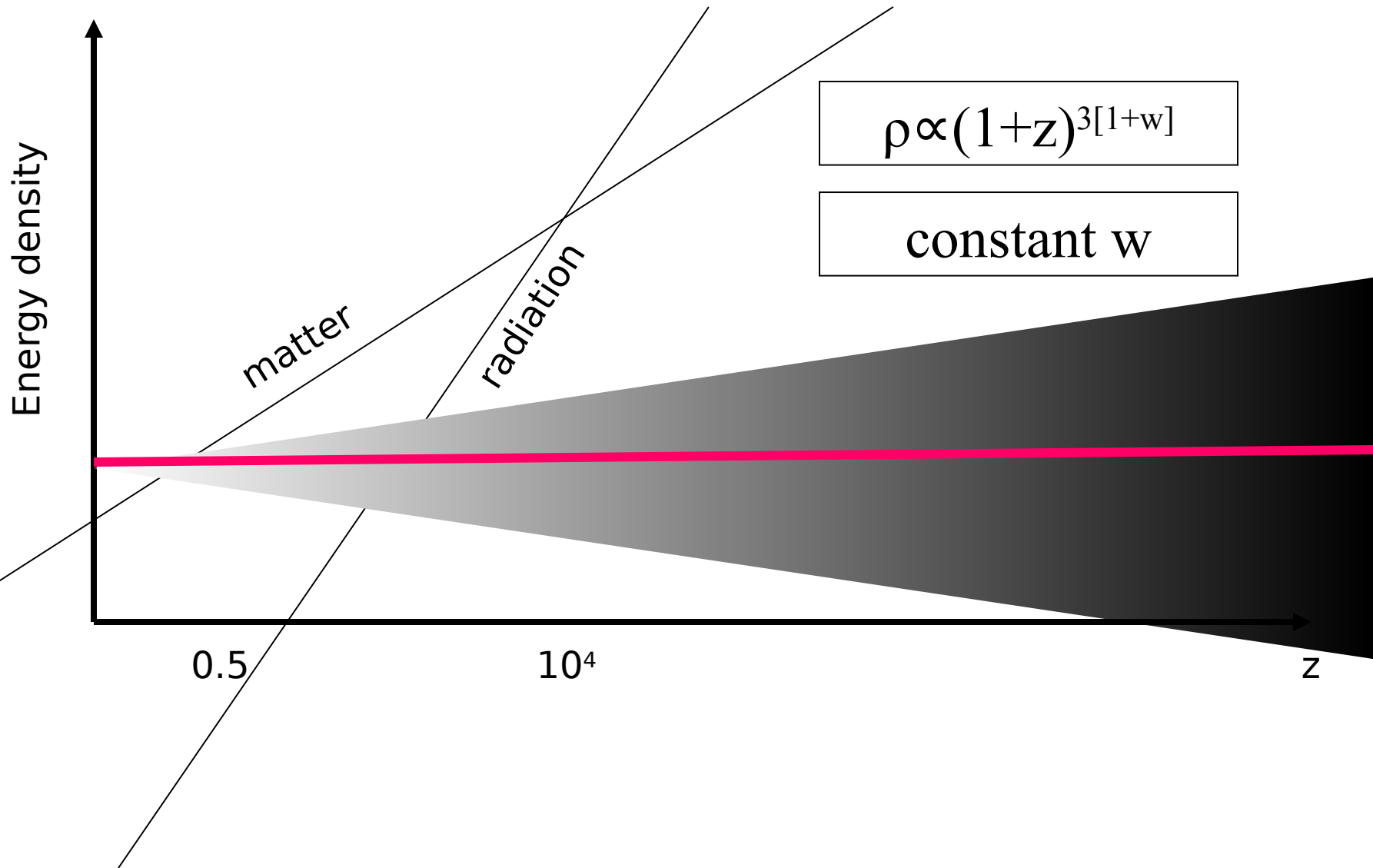
# The promise of CMB lensing



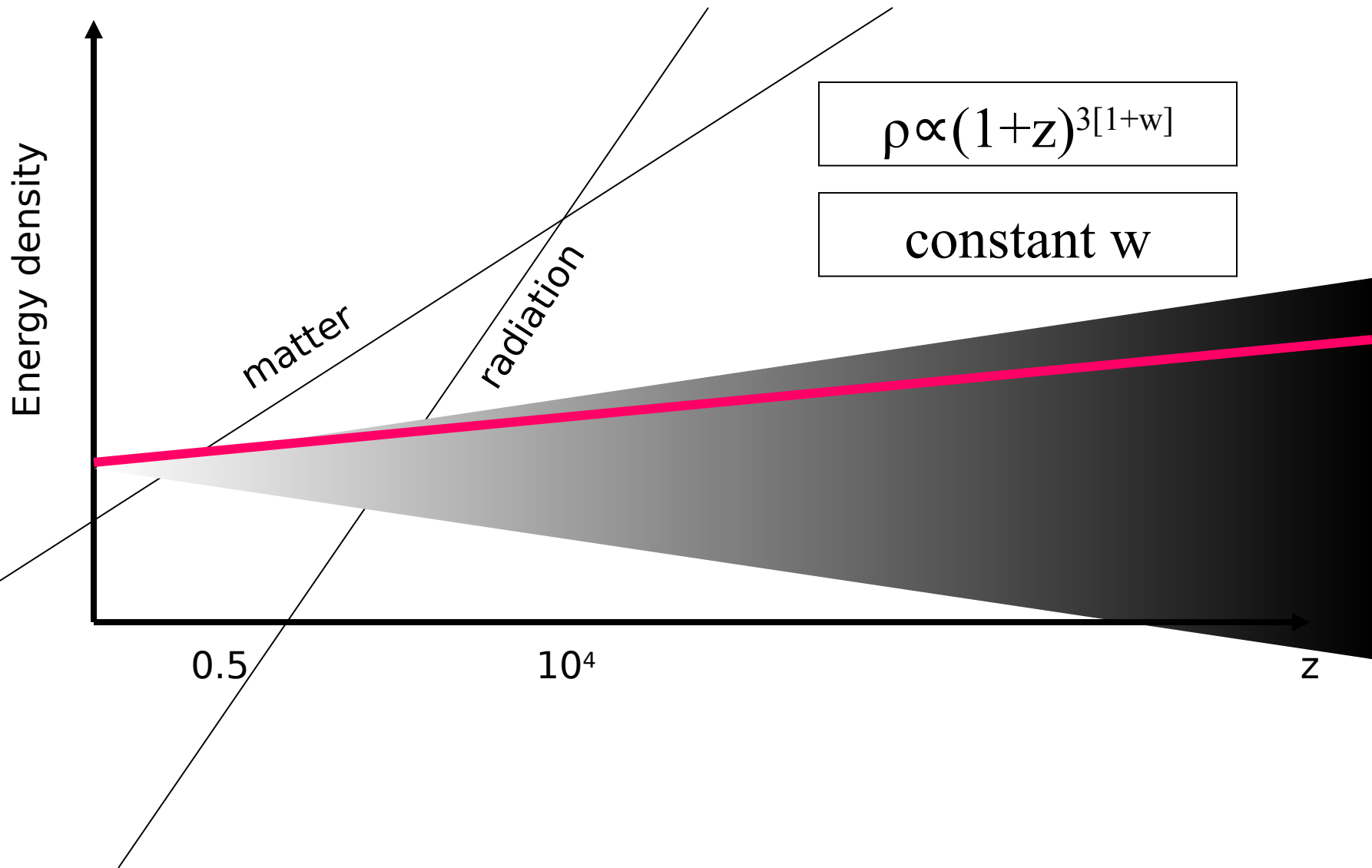
# Dark energy density



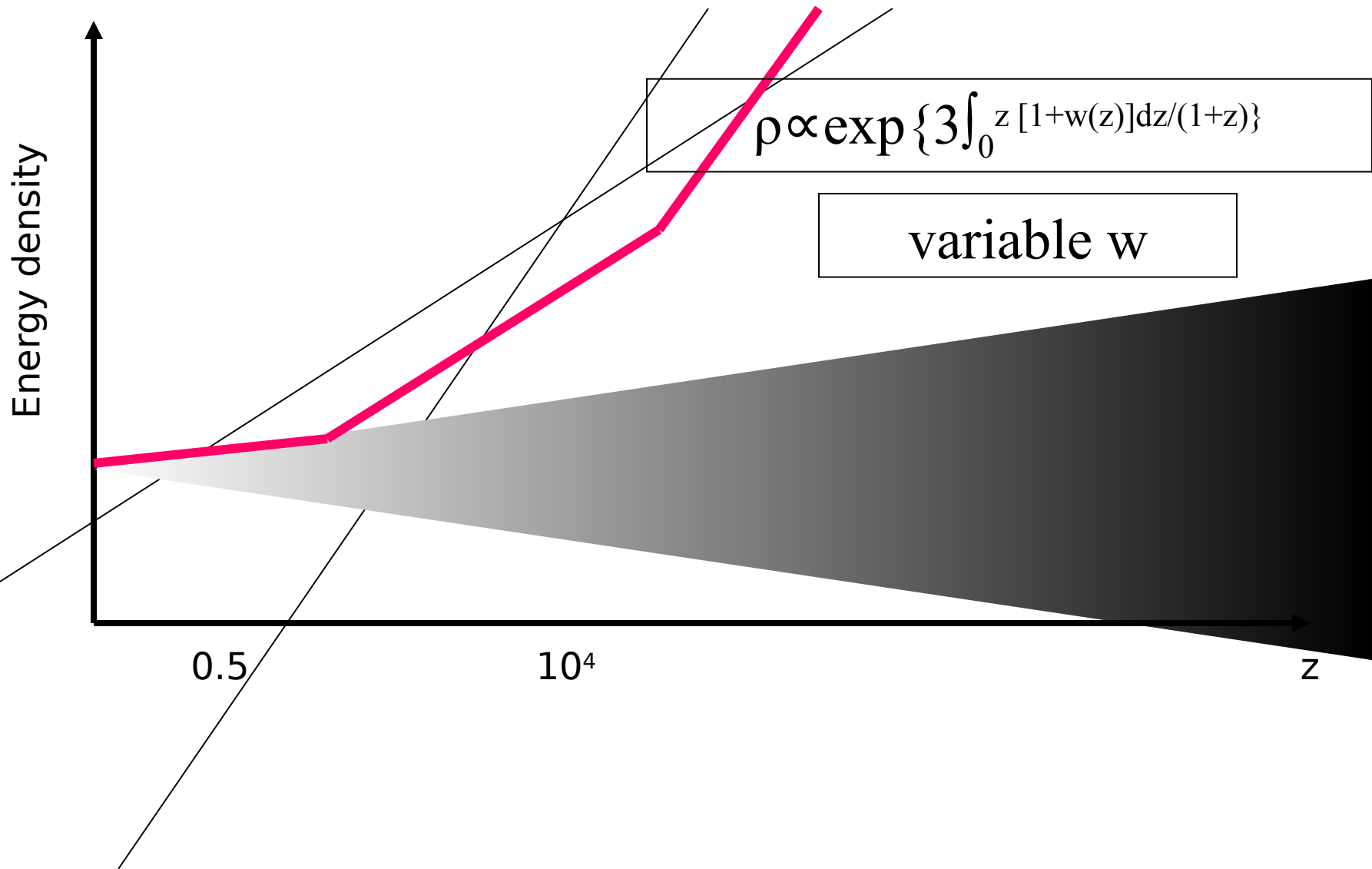
# Constant Dark Energy ( $\Lambda$ )



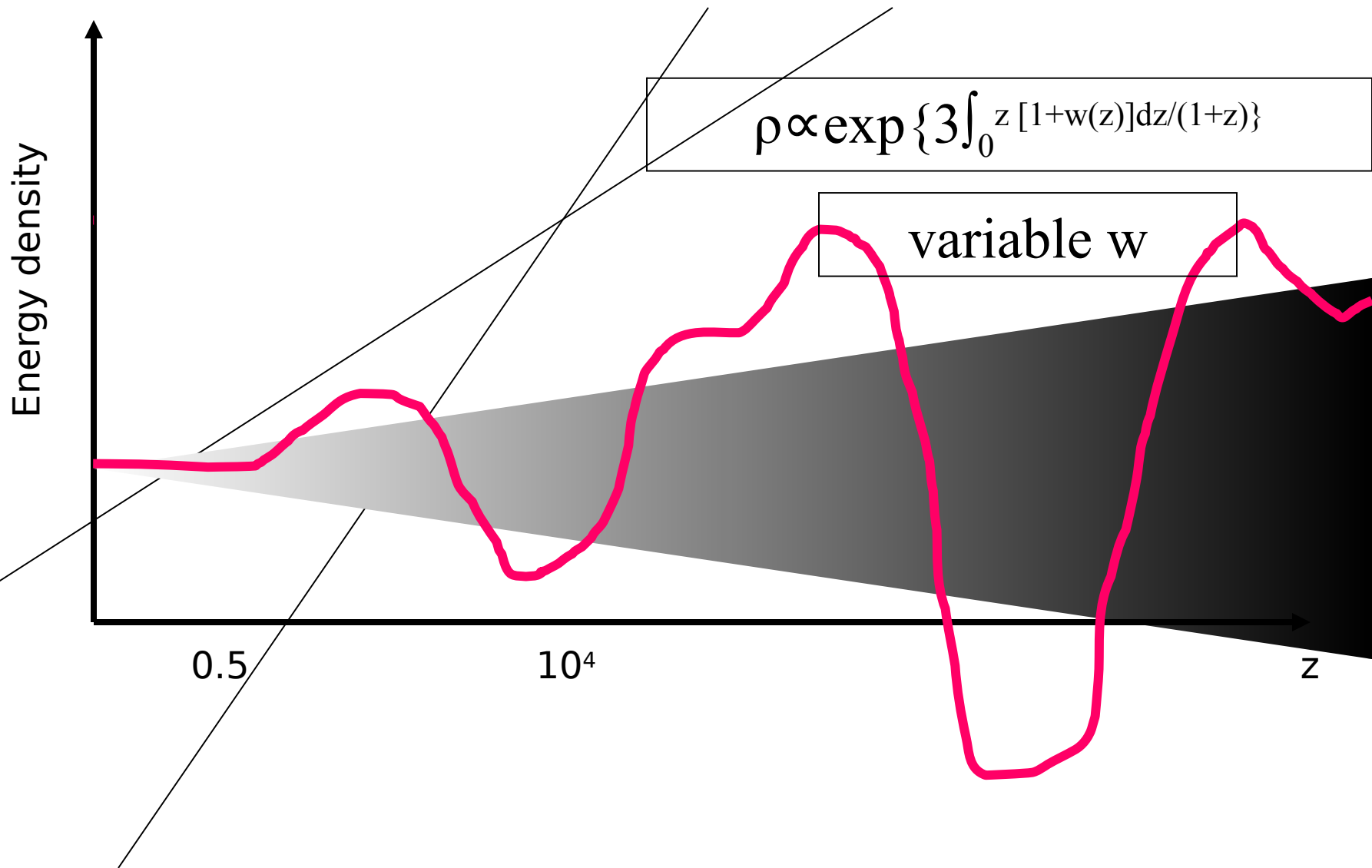
# Scaling Dark Energy



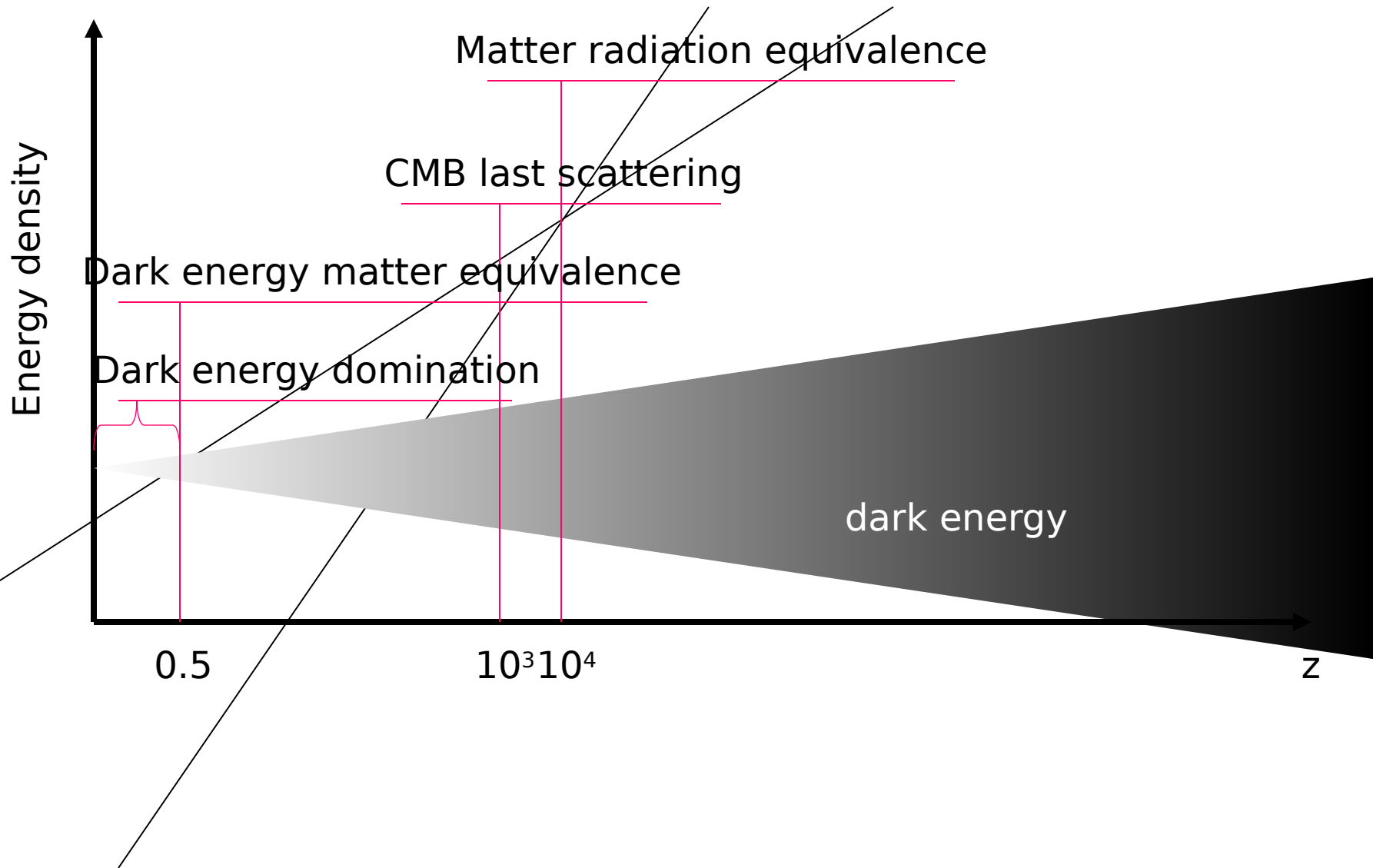
# Early Dark Energy



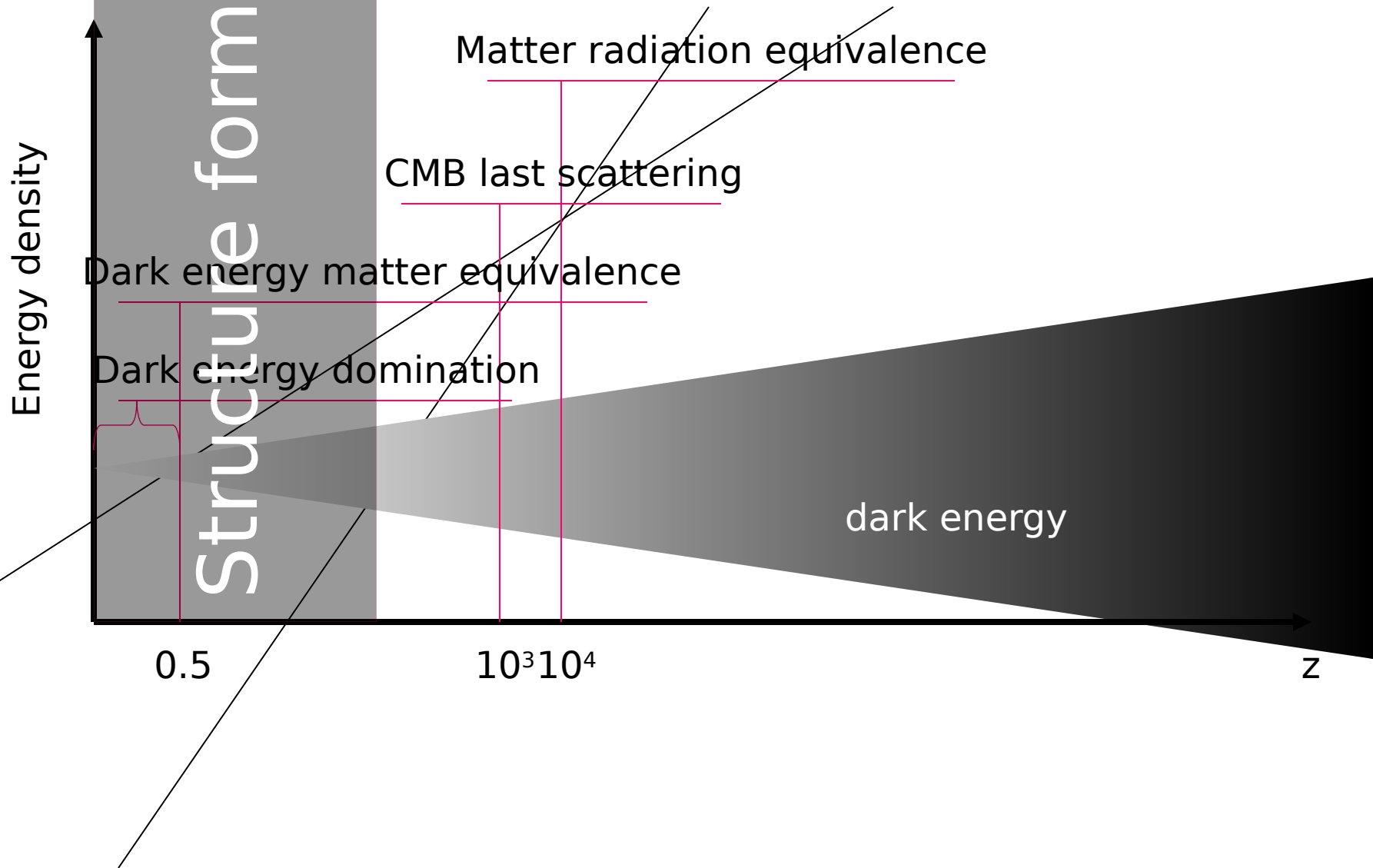
# Messy Dark Energy



# The promise of CMB lensing

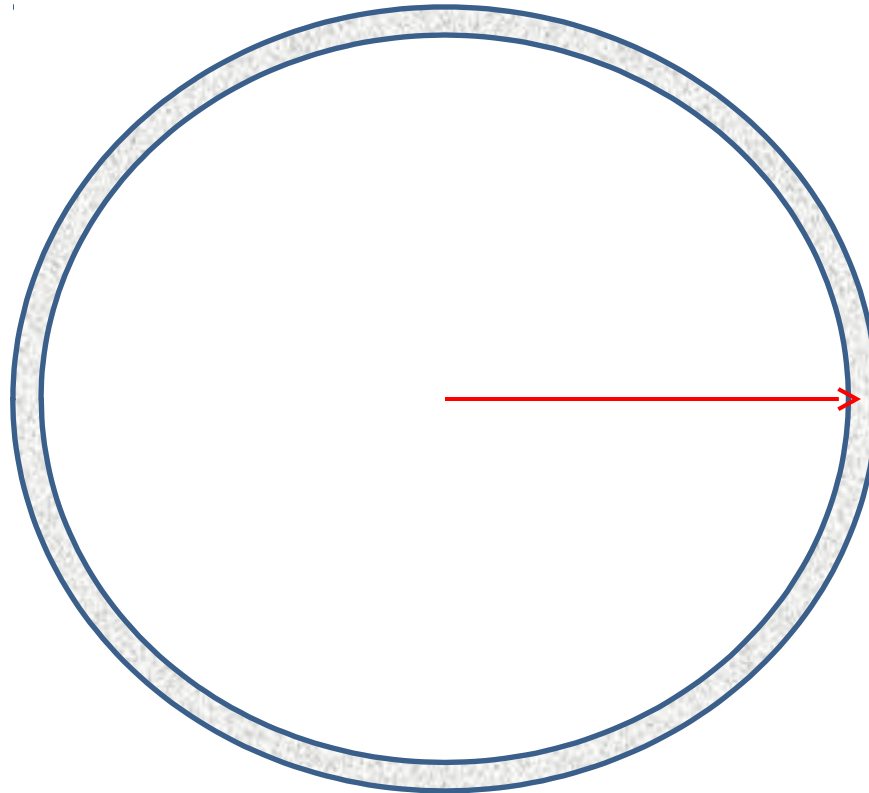


# The promise of CMB lensing



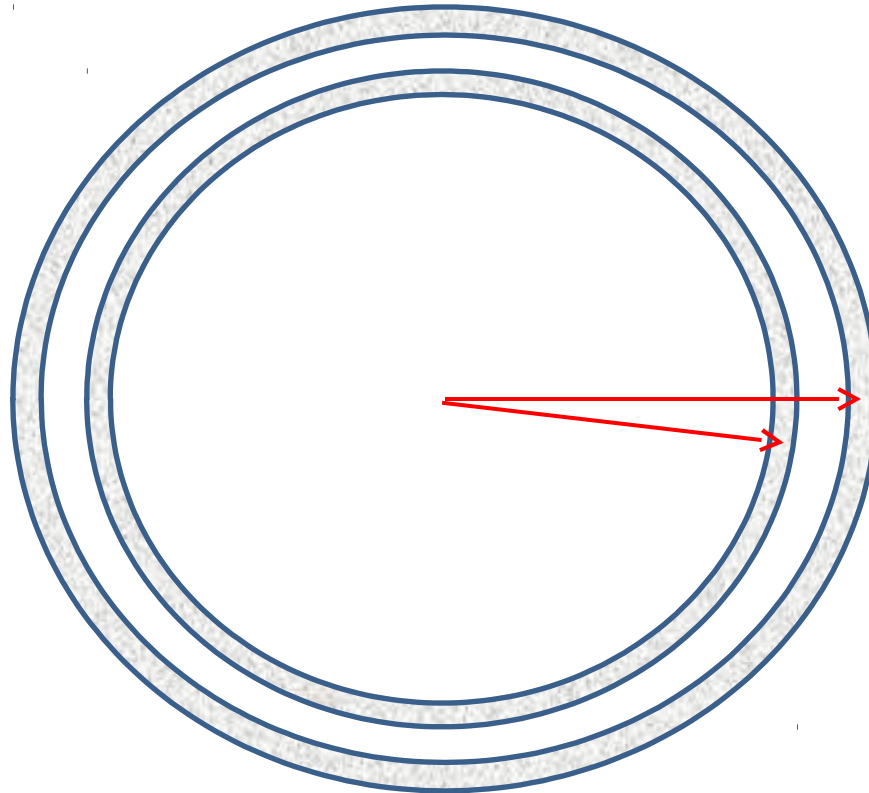
# Constraining the expansion with distances

$$D = H_0^{-1} \int_0^z \frac{dz}{[\sum_i \Omega_i (1+z)^{3(1+w_i)}]^{1/2}}$$



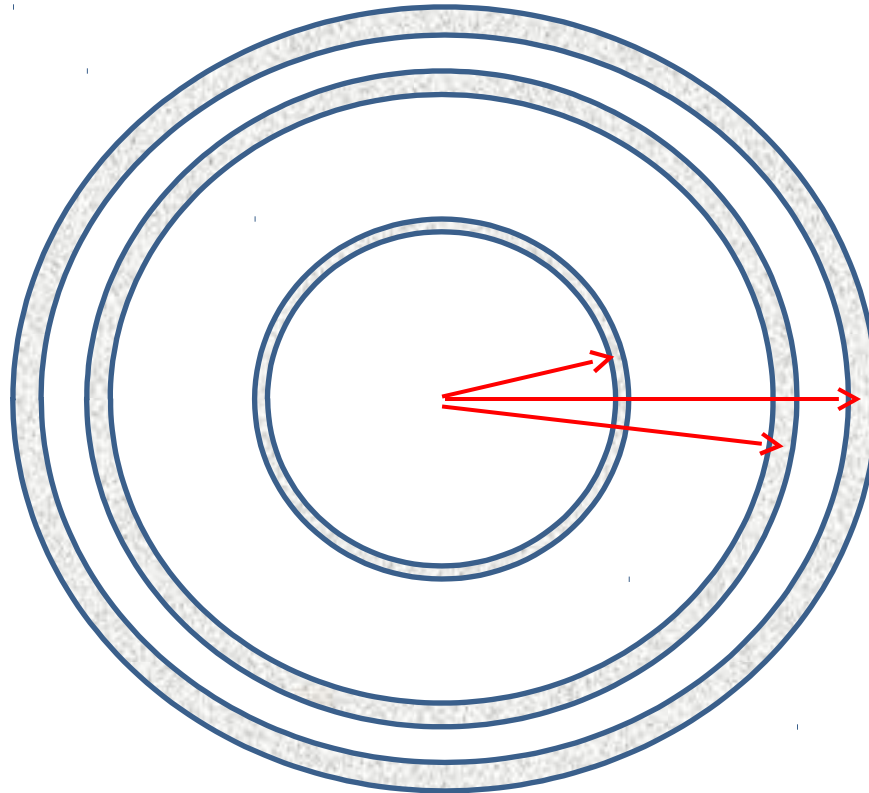
# Constraining the expansion with distances

$$D = H_0^{-1} \int_0^z \frac{dz}{[\sum_i \Omega_i (1+z)^{3(1+w_i)}]^{1/2}}$$

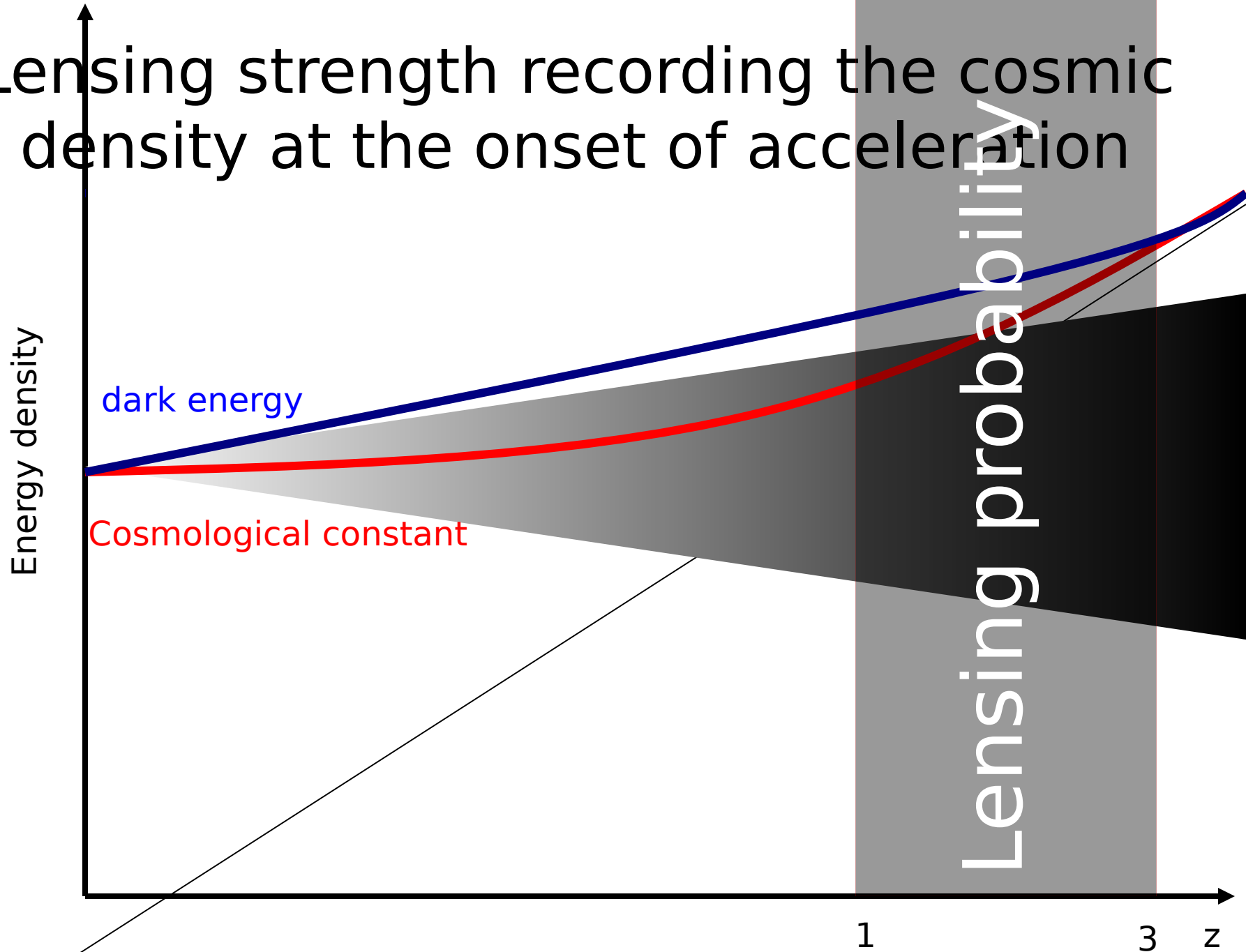


# Constraining the expansion with distances

$$D = H_0^{-1} \int_0^z \frac{dz}{[\sum_i \Omega_i (1+z)^{3(1+w_i)}]^{1/2}}$$



# Lensing strength recording the cosmic density at the onset of acceleration



# CMB lensing

$$\phi(\hat{n}) = \int_0^{\chi_\infty} d\chi D(\chi) [\Psi(\hat{n}, \chi) - \Phi(\hat{n}, \chi)] \int_\chi^{\chi_\infty} d\chi' \frac{D(\chi' - \chi)}{D(\chi')} g(\chi')$$

$$C_l^{\phi\phi} = \frac{32}{\pi} \int_0^{\chi_\infty} d\chi g'(\chi) \int_0^{\chi_\infty} d\chi' g'(\chi') \int k^2 dk u_l[k D(\chi)] u_l[k D(\chi')] \cdot \left[ \frac{1}{4} \langle \Psi(k, \chi) \Psi(k, \chi') \rangle + \frac{1}{4} \langle \Phi(k, \chi) \Phi(k, \chi') \rangle - \frac{1}{2} \langle \Psi(k, \chi) \Phi(k, \chi') \rangle \right]$$

$$C_l^{\Theta\phi} = \frac{2}{\pi} \int_0^{\chi_\infty} d\chi g'(\chi) \int_0^{\chi_\infty} d\chi' \int k^2 dk u_l(k D(\chi)) u_l(k D(\chi')) \cdot \left[ \langle \Psi(k, \chi) \dot{\Psi}(k, \chi') \rangle + \langle \Phi(k, \chi) \dot{\Phi}(k, \chi') \rangle - \langle \Psi(k, \chi) \dot{\Phi}(k, \chi') \rangle - \langle \dot{\Psi}(k, \chi') \Phi(k, \chi) \rangle \right]$$

# CMB lensing

$$\phi(\hat{n}) = \int_0^{\chi_\infty} d\chi \left[ D(\chi) [\Psi(\hat{n}, \chi) - \Phi(\hat{n}, \chi)] \int_\chi^{\chi_\infty} d\chi' \frac{D(\chi' - \chi)}{D(\chi')} \right] \gamma(\chi')$$

$$C_l^{\phi\phi} = \frac{32}{\pi} \int_0^{\chi_\infty} d\chi g'(\chi) \int_0^{\chi_\infty} d\chi' g'(\chi') \int k^2 dl \left[ u_l[k D(\chi)] u_l[k D(\chi')] \right] \cdot \left[ \frac{1}{4} \langle \Psi(k, \chi) \Psi(k, \chi') \rangle + \frac{1}{4} \langle \Phi(k, \chi) \Phi(k, \chi') \rangle - \frac{1}{2} \langle \Psi(k, \chi) \Phi(k, \chi') \rangle \right]$$

$$C_l^{\Theta\phi} = \frac{2}{\pi} \int_0^{\chi_\infty} d\chi g'(\chi) \int_0^{\chi_\infty} d\chi' \int k^2 dl \left[ u_l(k D(\chi)) u_l(k D(\chi')) \right] \cdot \left[ \langle \Psi(k, \chi) \dot{\Psi}(k, \chi') \rangle + \langle \Phi(k, \chi) \dot{\Phi}(k, \chi') \rangle - \langle \Psi(k, \chi) \dot{\Phi}(k, \chi') \rangle - \langle \dot{\Psi}(k, \chi') \Phi(k, \chi) \rangle \right]$$

# CMB lensing

$$\phi(\hat{n}) = \int_0^{\chi_\infty} d\chi D(\chi) [\Psi(\hat{n}, \chi) - \Phi(\hat{n}, \chi)] \int_\chi^{\chi_\infty} d\chi' \frac{D(\chi' - \chi)}{D(\chi')} g(\chi')$$

$$C_l^{\phi\phi} = \frac{32}{\pi} \int_0^{\chi_\infty} d\chi g'(\chi) \int_0^{\chi_\infty} d\chi' g'(\chi') \int k^2 dk u_l[k D(\chi)] u_l[k D(\chi')] \cdot$$

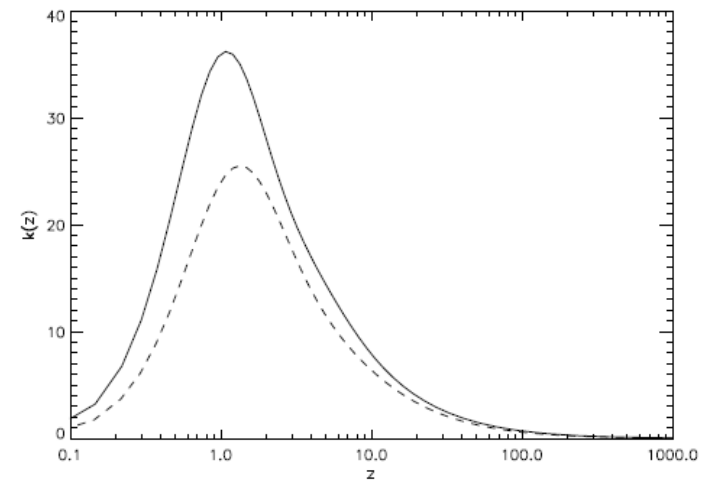
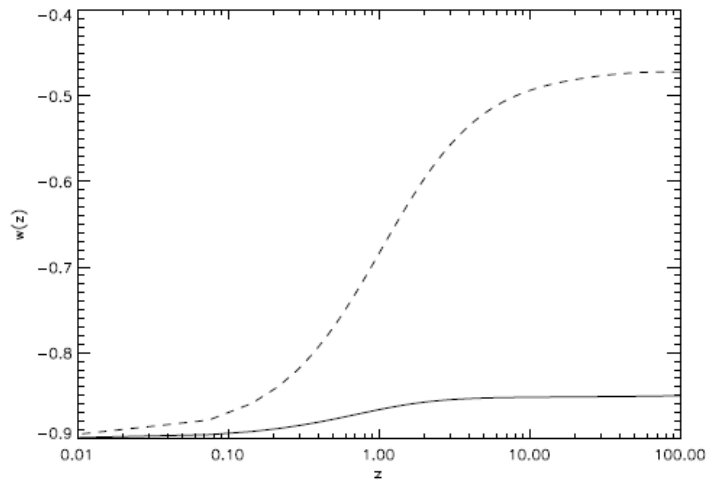
$$\cdot \left[ \frac{1}{4} \langle \Psi(k, \chi) \Psi(k, \chi') \rangle + \frac{1}{4} \langle \Phi(k, \chi) \Phi(k, \chi') \rangle - \frac{1}{2} \langle \Psi(k, \chi) \Phi(k, \chi') \rangle \right]$$

$$C_l^{\Theta\phi} = \frac{2}{\pi} \int_0^{\chi_\infty} d\chi g'(\chi) \int_0^{\chi_\infty} d\chi' \int k^2 dk u_l(k D(\chi)) u_l(k D(\chi')) \cdot$$

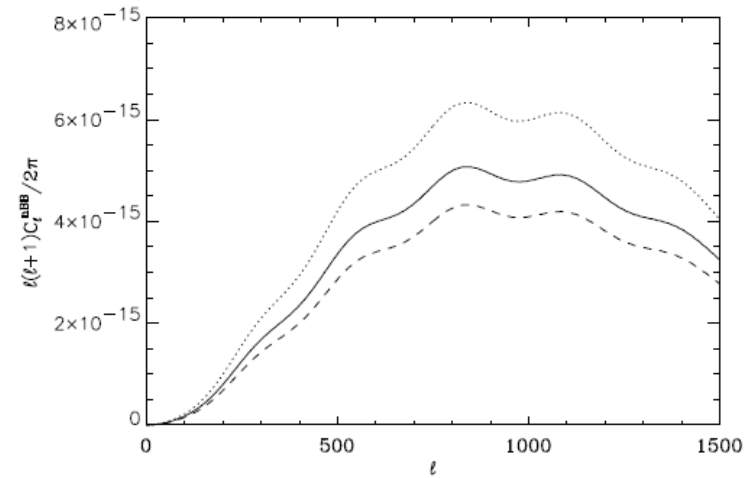
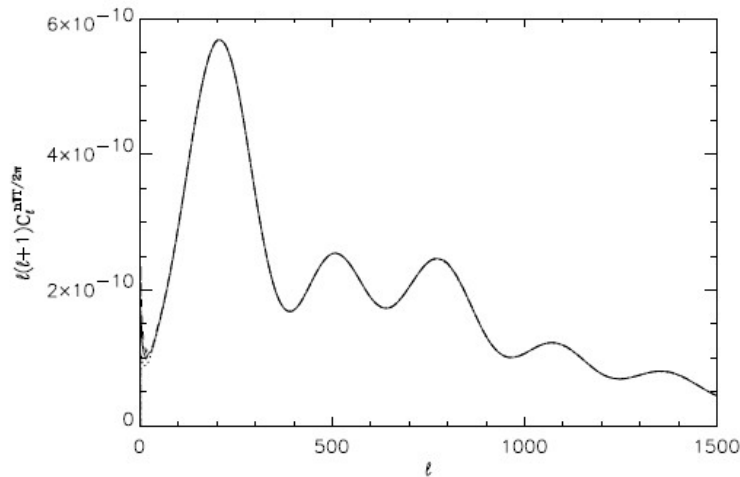
$$\cdot \left[ \langle \Psi(k, \chi) \dot{\Psi}(k, \chi') \rangle + \langle \Phi(k, \chi) \dot{\Phi}(k, \chi') \rangle - \langle \Psi(k, \chi) \dot{\Phi}(k, \chi') \rangle - \langle \dot{\Psi}(k, \chi') \Phi(k, \chi) \rangle \right]$$

**Perturbation dynamics**

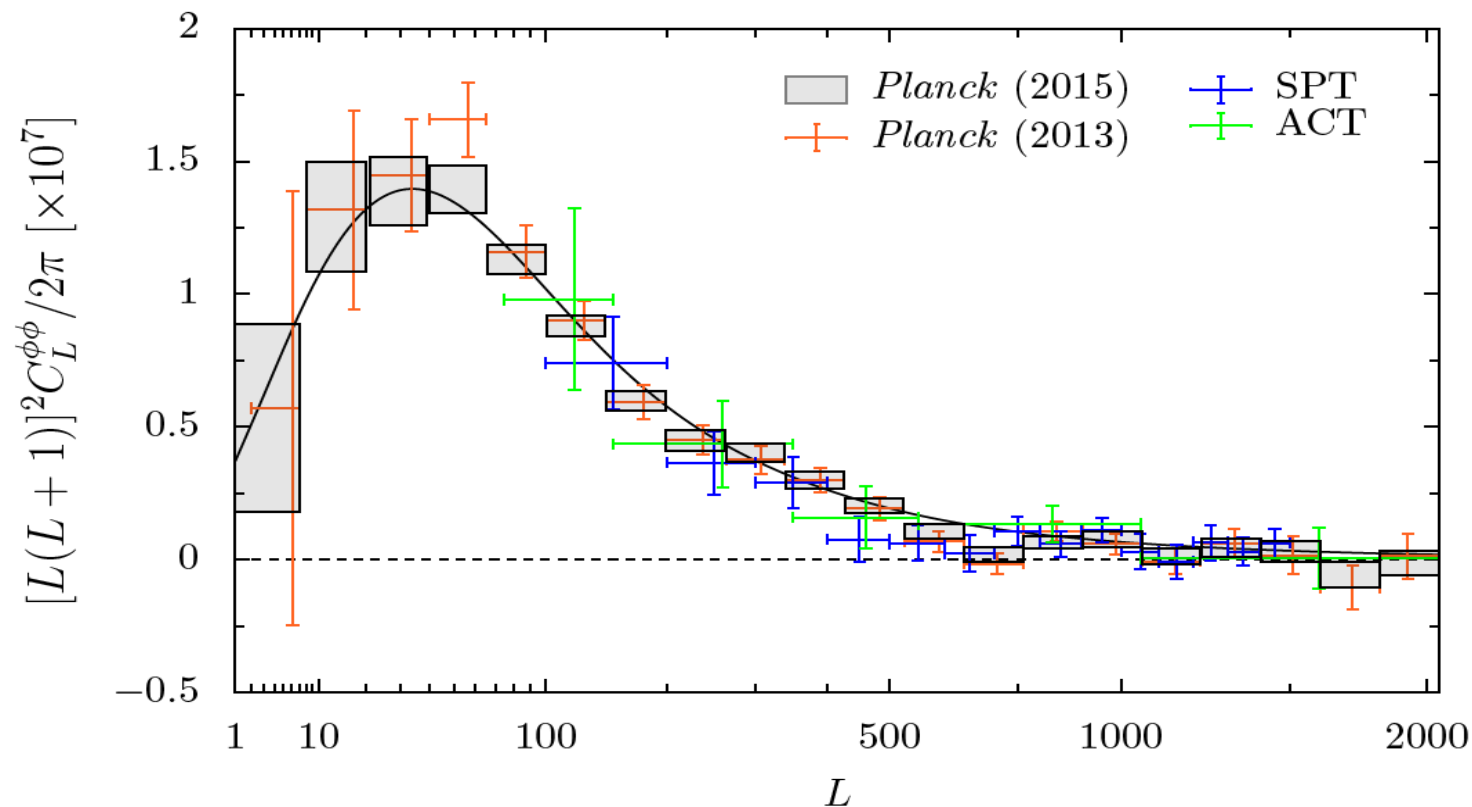
# Early dark energy and CMB lensing



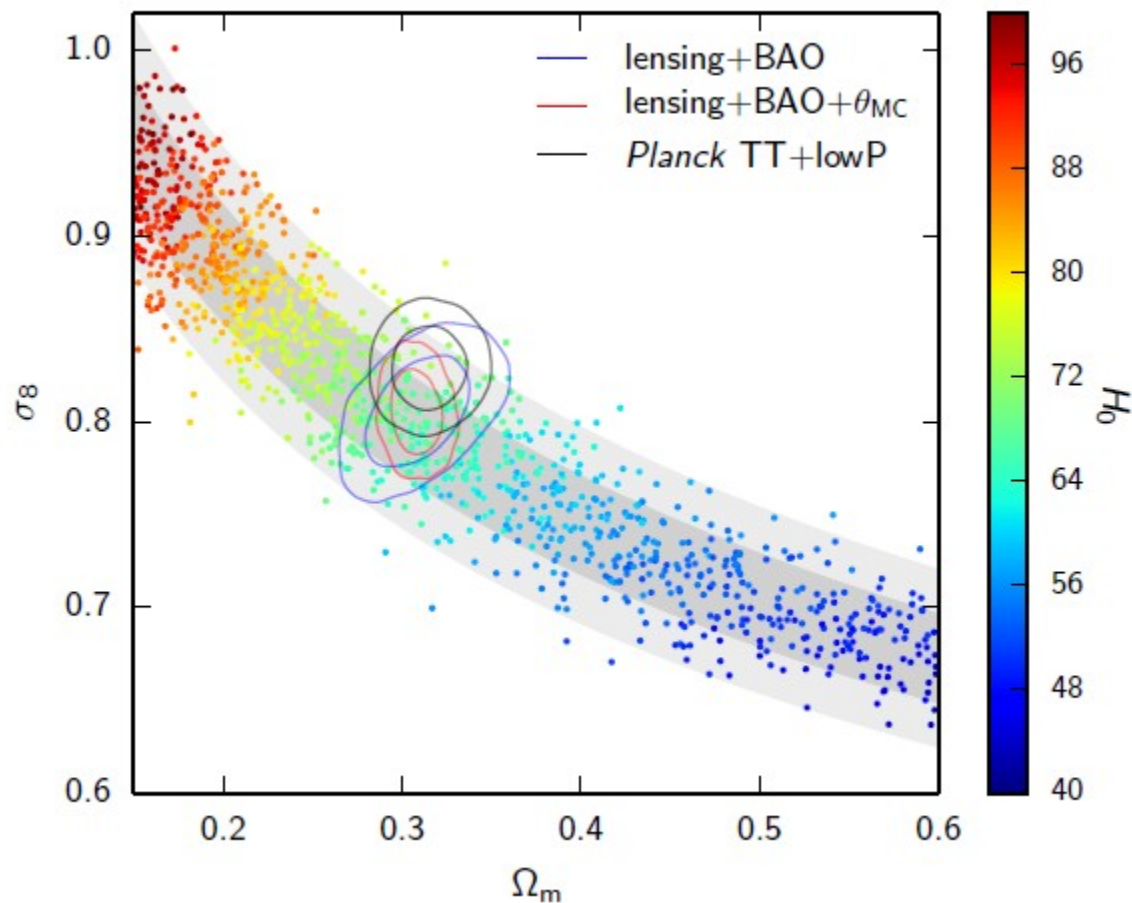
# CMBreaking projection degeneracy



# Planck lensing 2015



# Planck lensing 2015



# Lensing-LSS Cross-Correlations

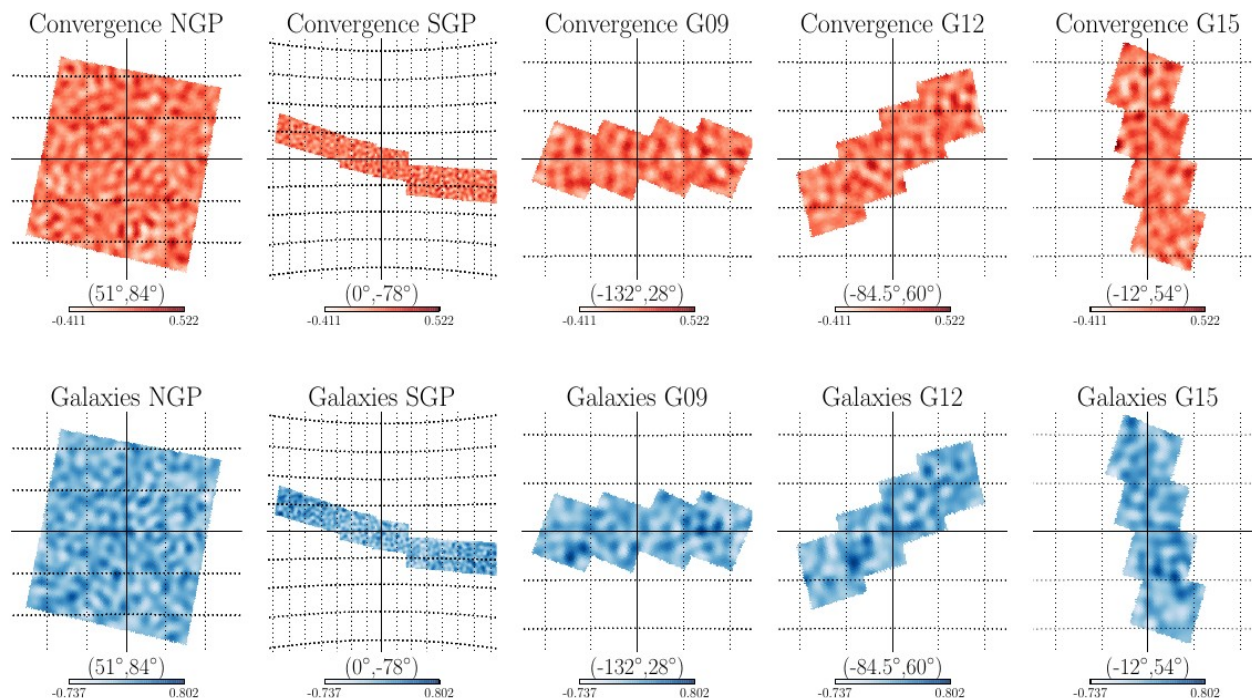


Fig. 5.— Convergence maps (upper row) and galaxy overdensity maps (lower row) in the H-ATLAS fields: noise-dominated multipoles  $\ell > 400$  have been filtered out. Galactic longitude and latitude ( $l, b$ ) of patches centers are provided in brackets. The grid overlay has spacing of  $3^\circ$  in each box.

# Lensing-LSS Cross-Correlations

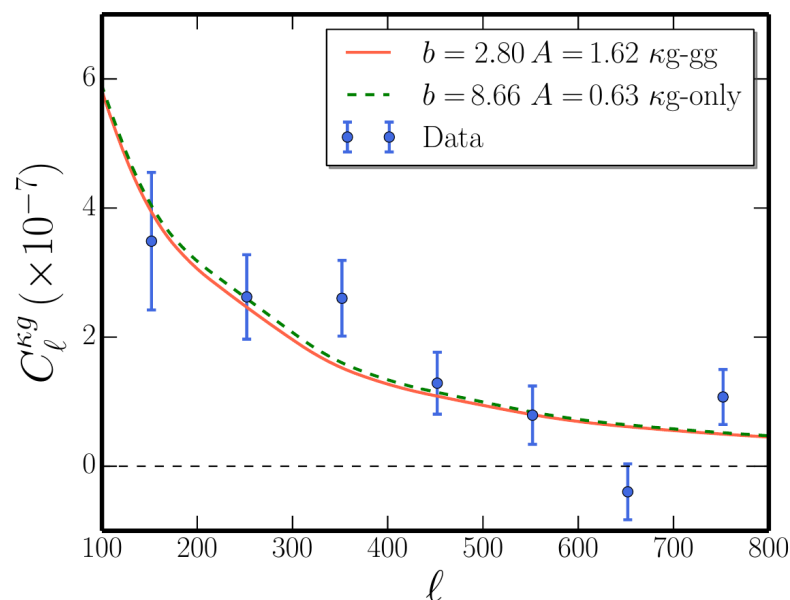
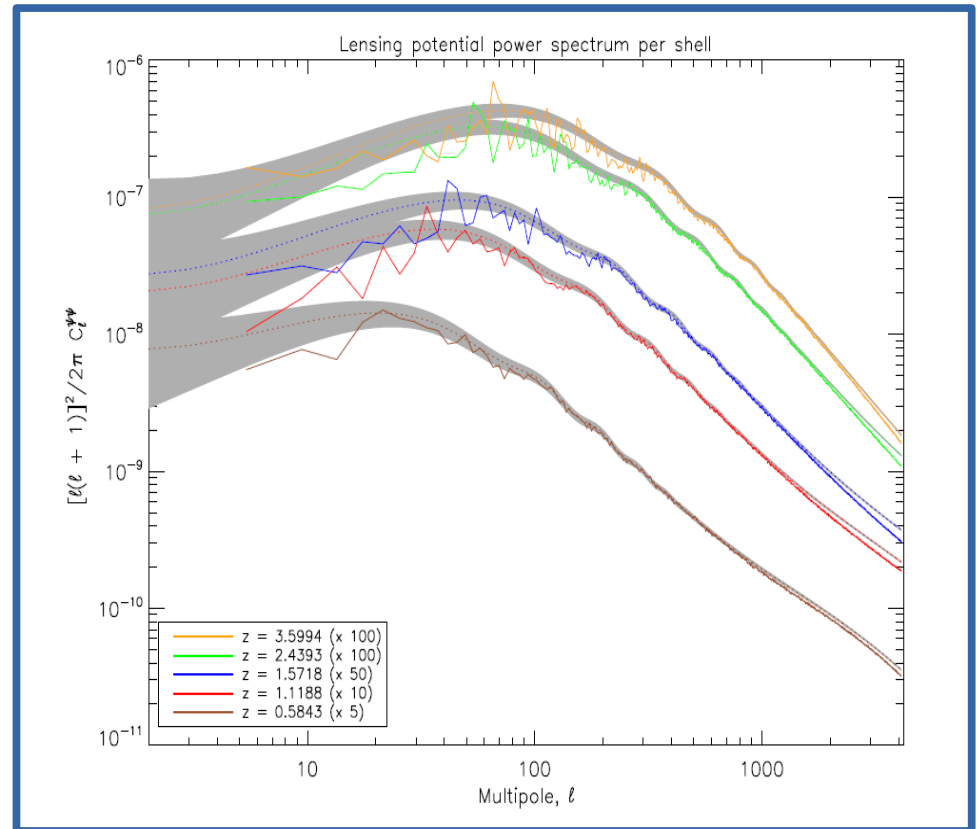
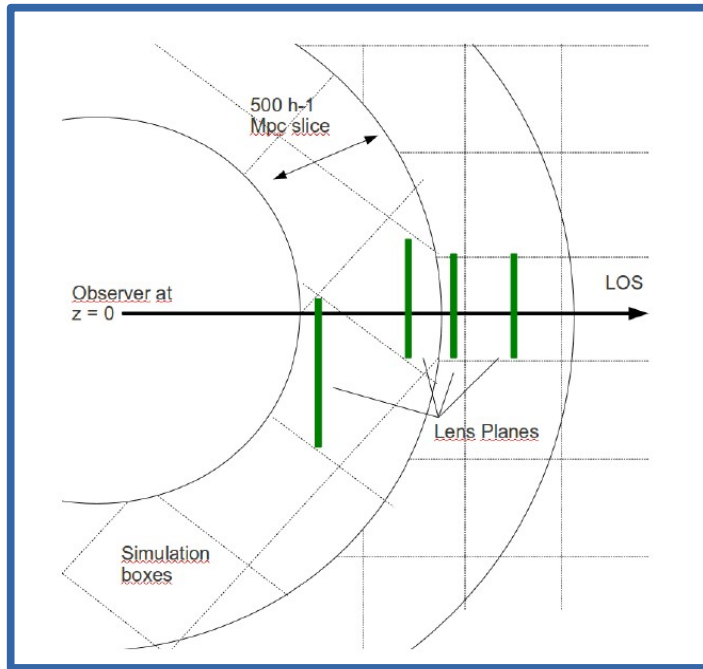


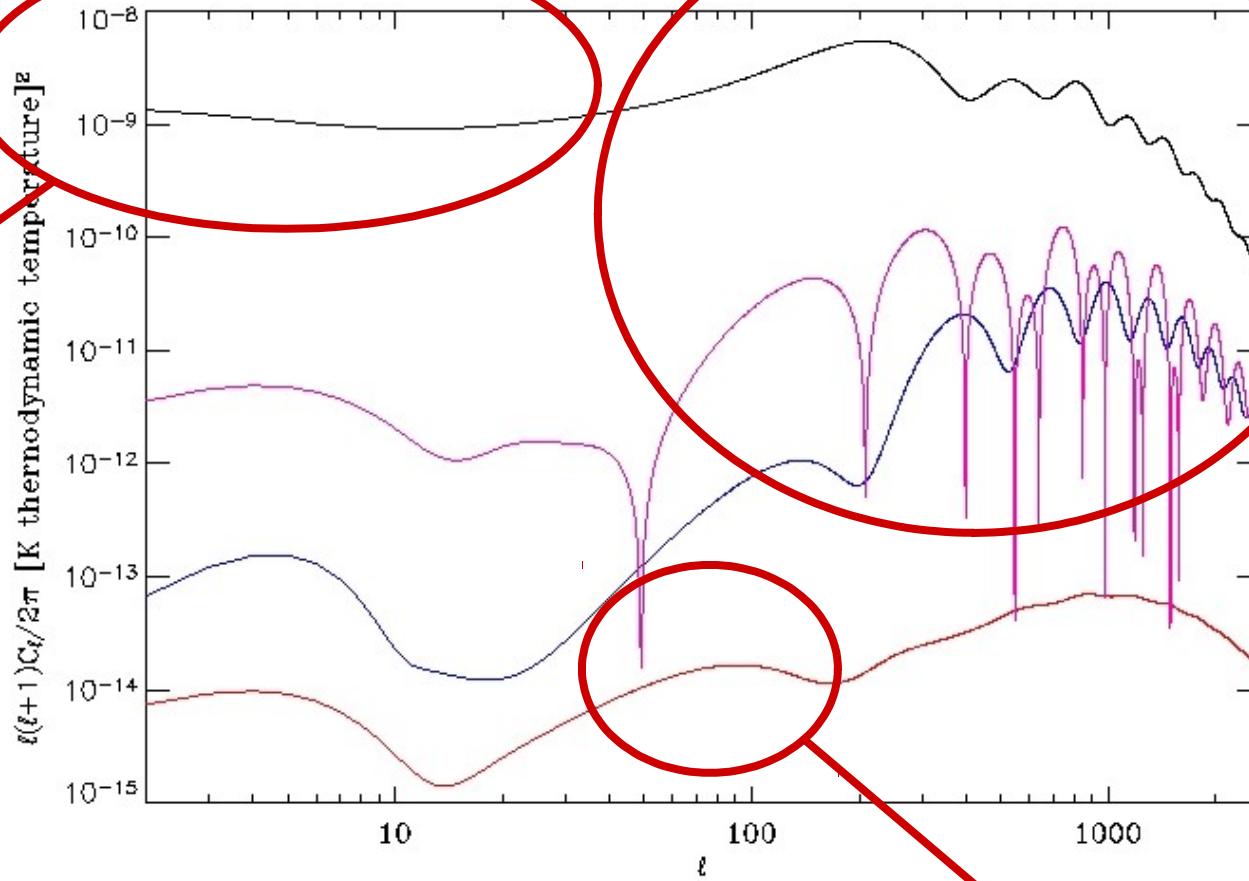
Fig. 9.— The CMB convergence - galaxy density cross-spectrum as measured from *Planck* and *Herschel* data. The data points are shown in blue, with error bars computed using the full covariance matrix obtained from Monte Carlo realizations of convergence maps. The theoretical spectra calculated with the bias values inferred from the likelihood analysis (as described in text) using the cross-correlation data only (solid red line) and the cross-correlation together with the galaxy auto-correlation data (dot-dashed green line) are also shown; we fix  $\alpha = 3$  in this analysis. The null (no correlation) hypothesis is rejected at the  $20\sigma$  level.

# CMB-N-body Lensing



# CMB angular power spectrum

Acoustic oscillations



Primordial power

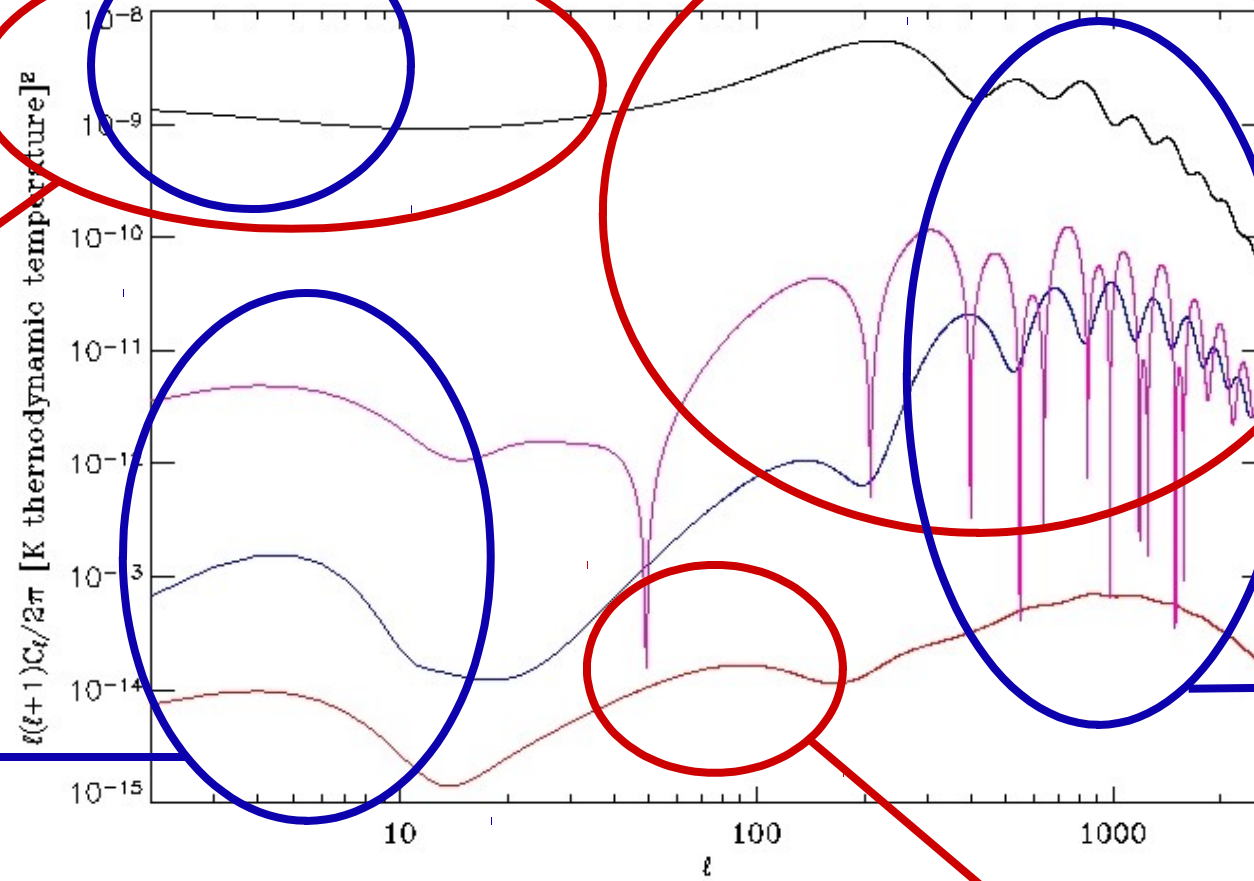
Angle  $\approx 200/\ell$  degrees

Gravitational waves

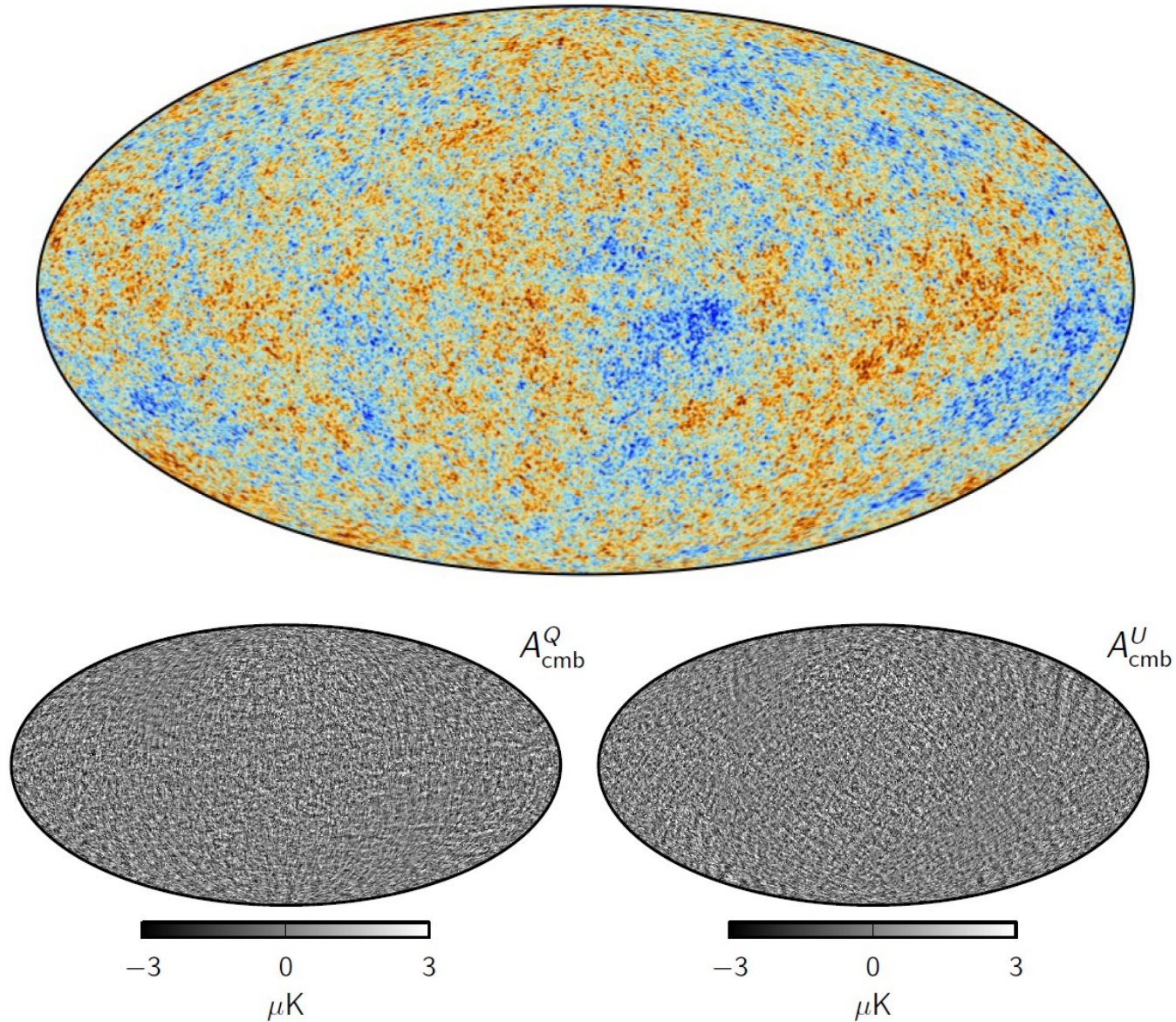
# CMB angular power spectrum

ISW

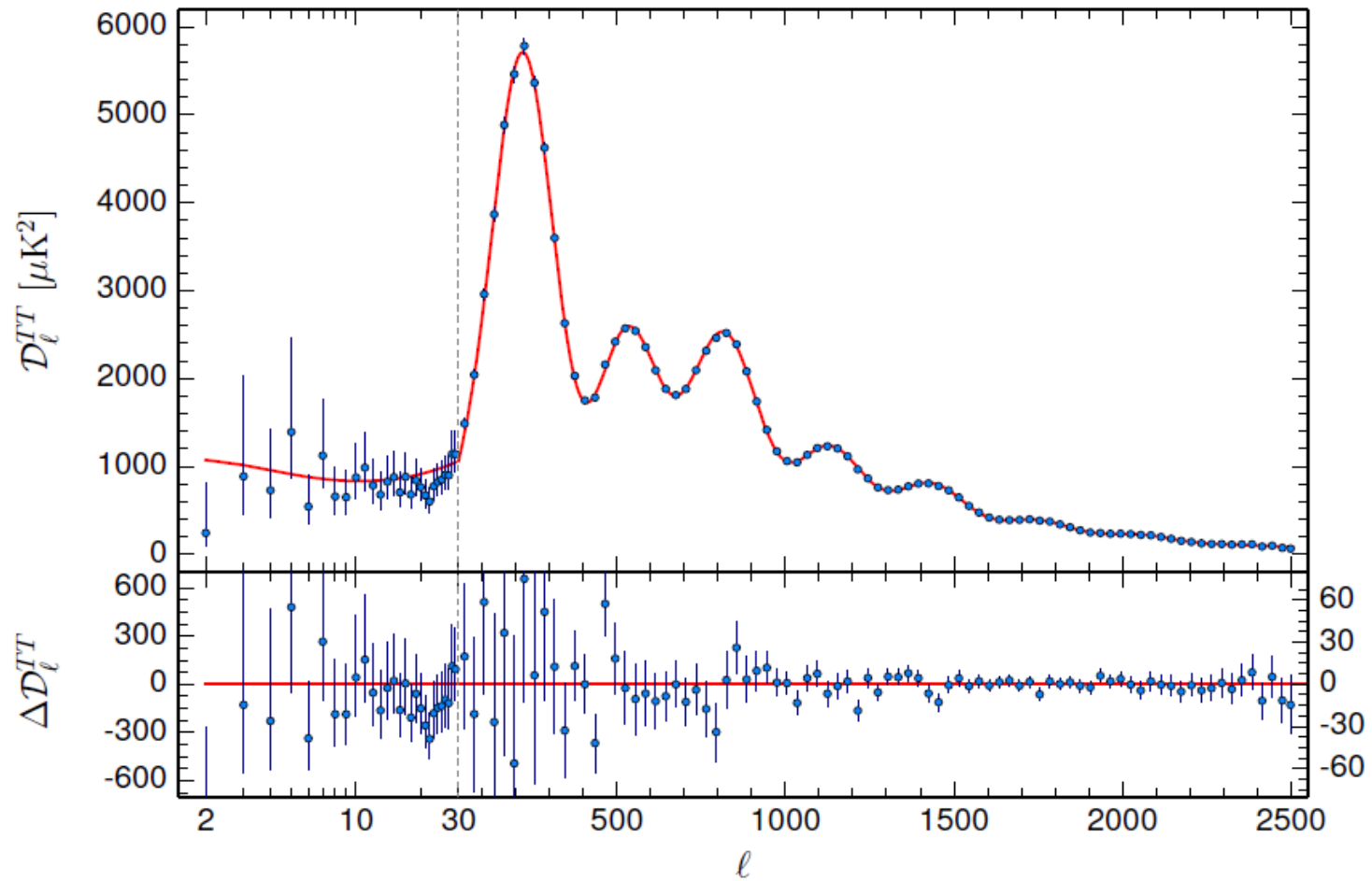
Acoustic oscillations



# Planck 2015 CMB

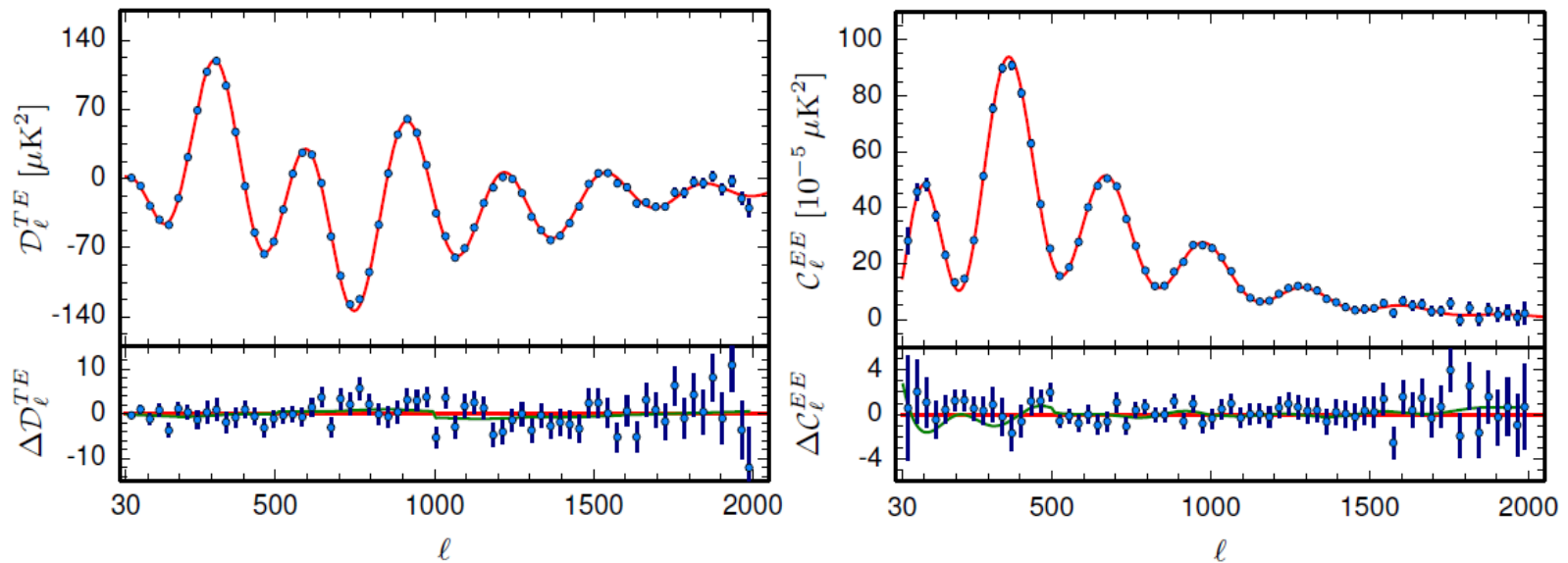


# Planck 2015 TT power spectrum

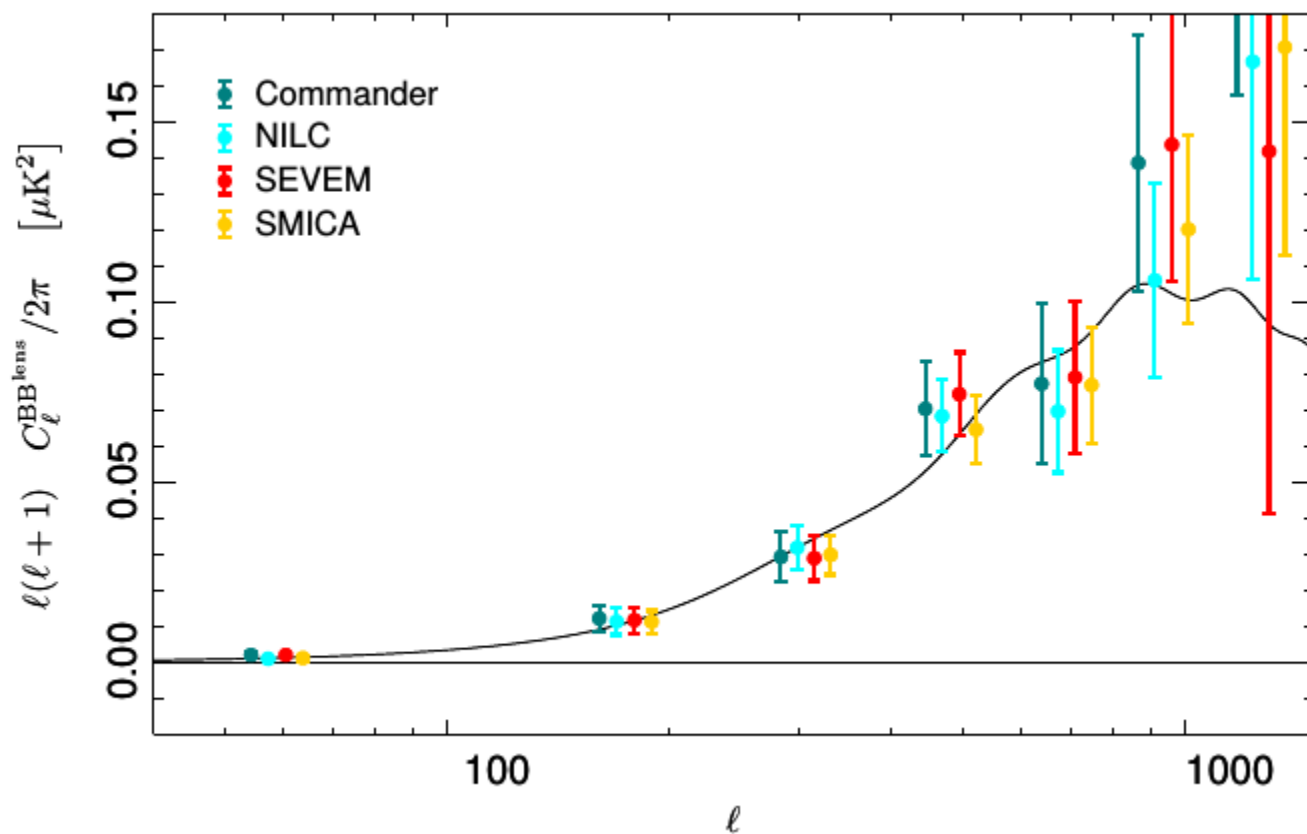


# Planck 2015

## TE and EE power spectra



# Planck 2015 BB power spectra



# Concluding remarks, I

- Effects from the Early Universe and Large Scale Structure affect large parts of CMB anisotropies, and directly compete for detection, for example in the B-modes
- Measurements of the abundance of primordial Gravitational Waves, and constraints on the Dark Energy dynamics must proceed jointly in the analysis of data from operating and future high resolution and sensitivity probes, LSS surveys
- Cross-correlation studies gathering relevance, large scale simulations CMB LSS experiments being implemented
- ...

The scientific results that we present today are a product of the Planck Collaboration, including individuals from more than 100 scientific institutes in Europe, the USA and Canada

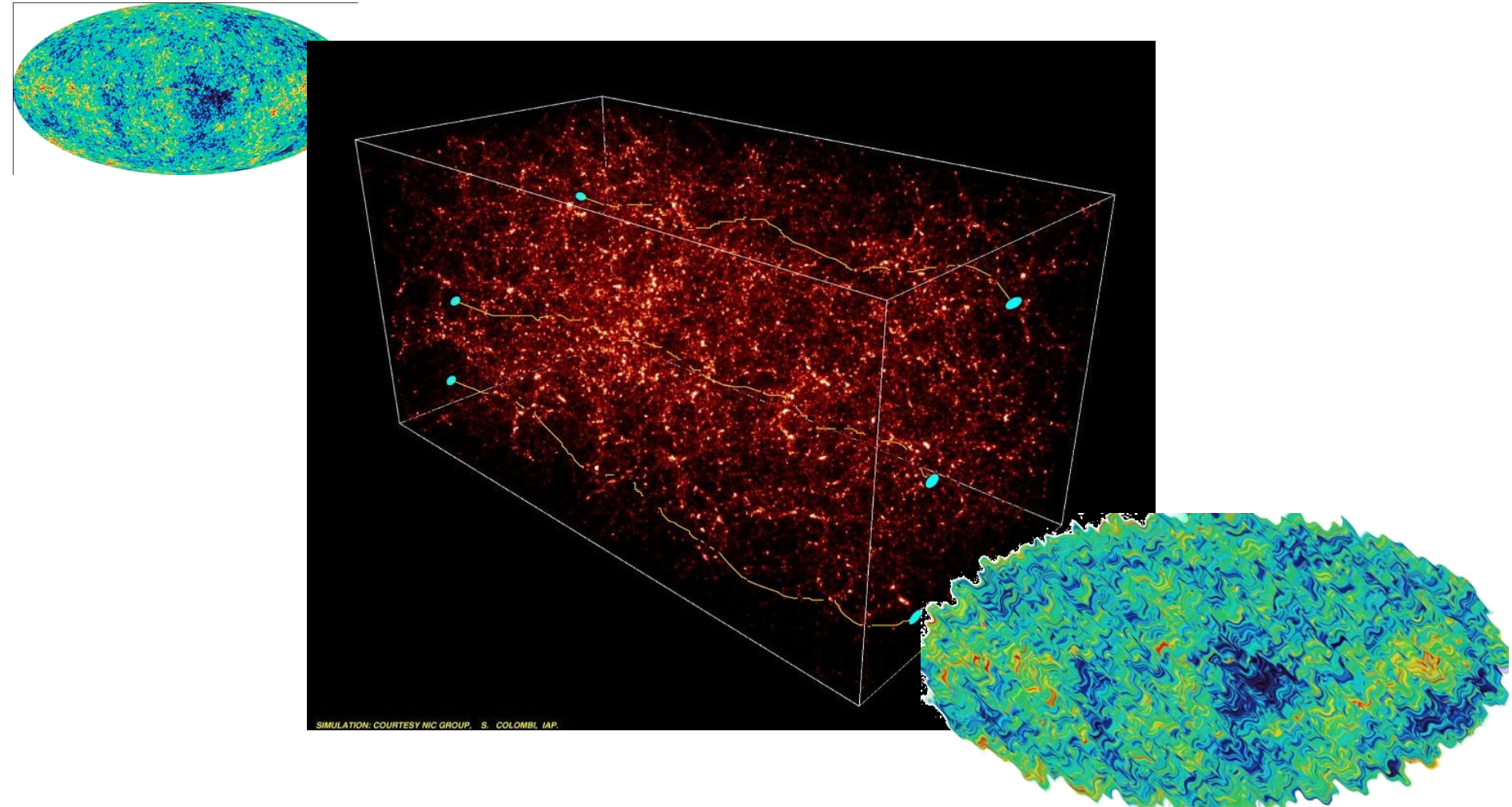


Planck is a project of the European Space Agency, with instruments provided by two scientific Consortia funded by ESA member states (in particular the lead countries: France and Italy) with contributions from NASA (USA), and telescope reflectors provided in a collaboration between ESA and a scientific Consortium led and funded by Denmark.

Backup slides

# Extension I: CMB-N-body lensing

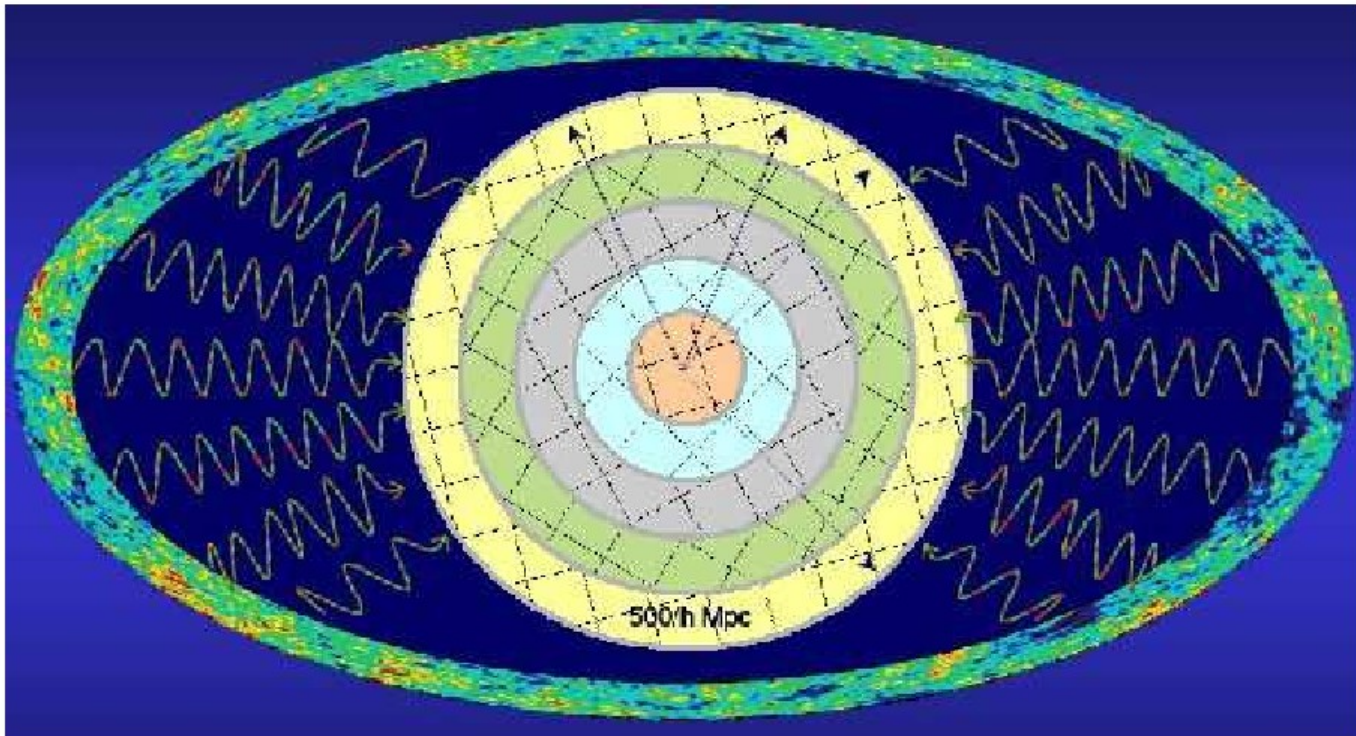
# CMB-N-Body lensing



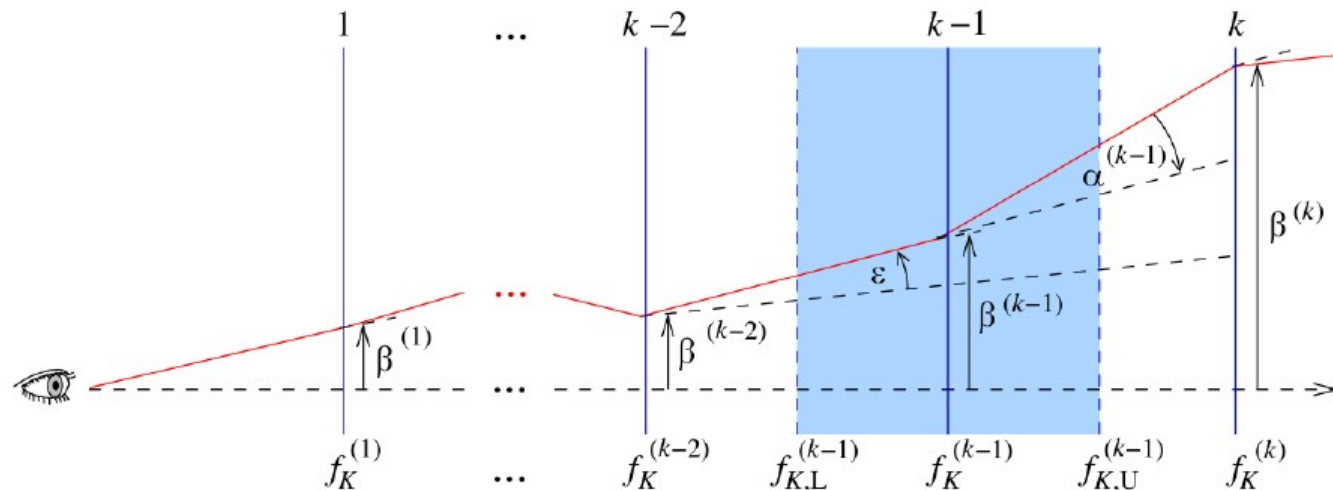
# CMB-N-body lensing: do we really need it?

- Do we have a satisfactory statistics in semi-analytic prescriptions for CMB lensing?
- Do we have an understanding of mildly and full non-linear Dark Matter clustering in Dark Energy models?
- Can we learn about simulations trying to get the signal right in particular at high resolution?
- Do we need simulations for understanding cross-correlation studies?
- ...

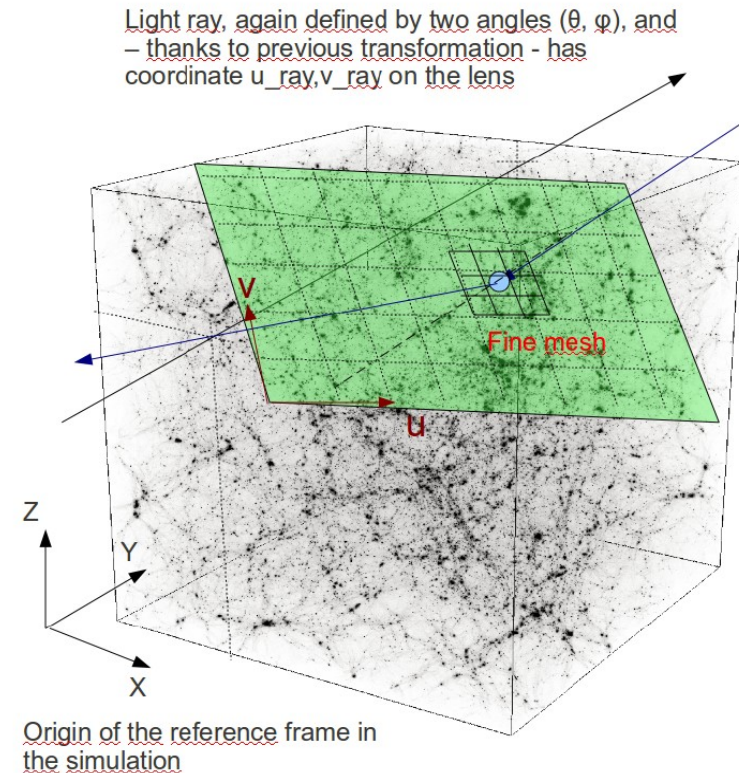
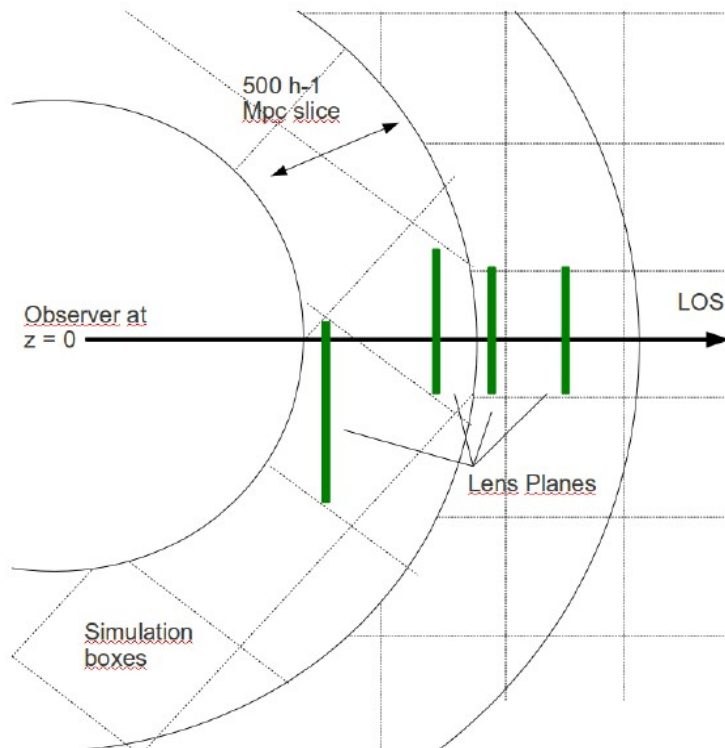
# CMB-N-body lensing



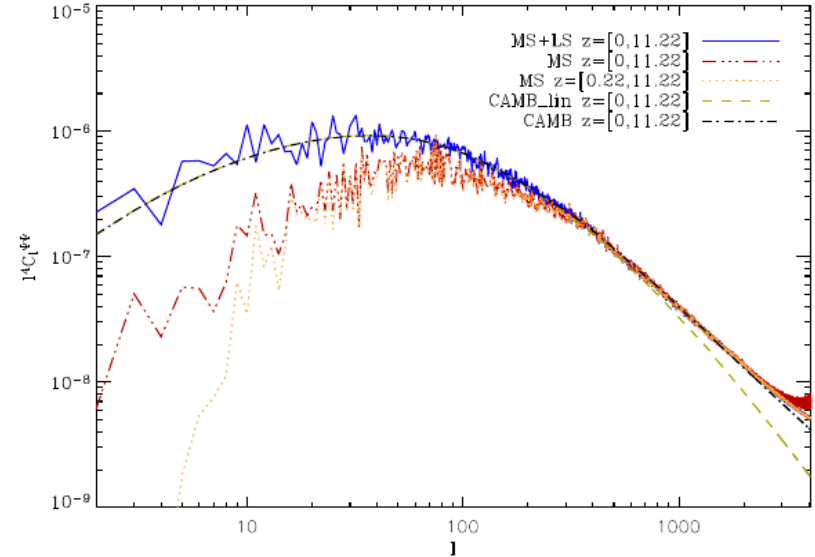
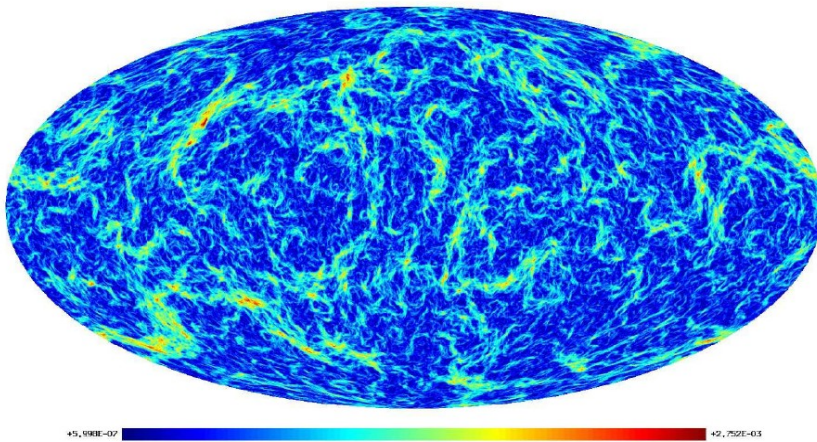
# CMB-N-body lensing ray-tracing: Born or non-Born?



# CMB-N-body lensing ray tracing

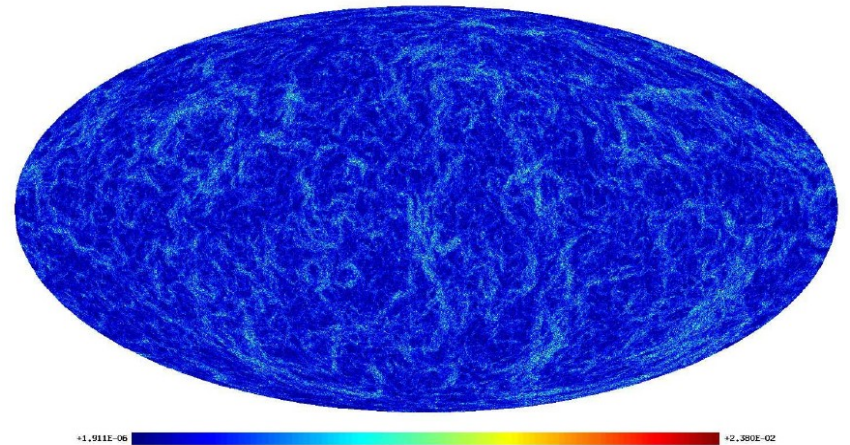
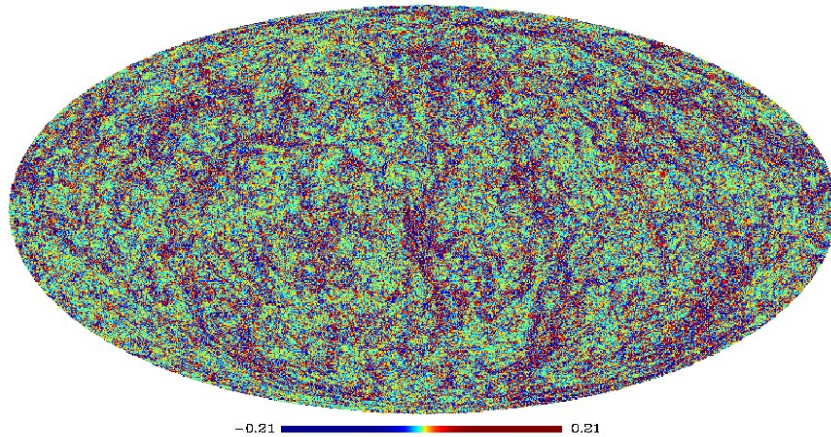


# CMB-N-body lensing

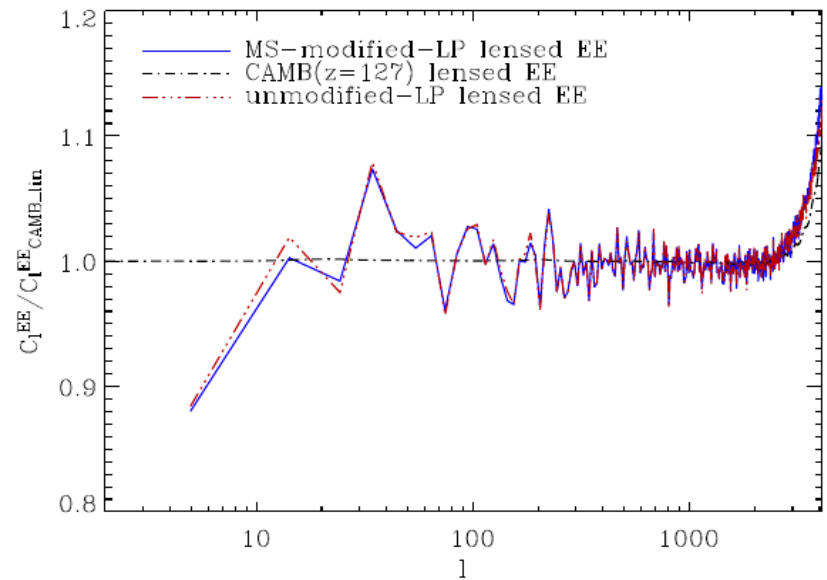
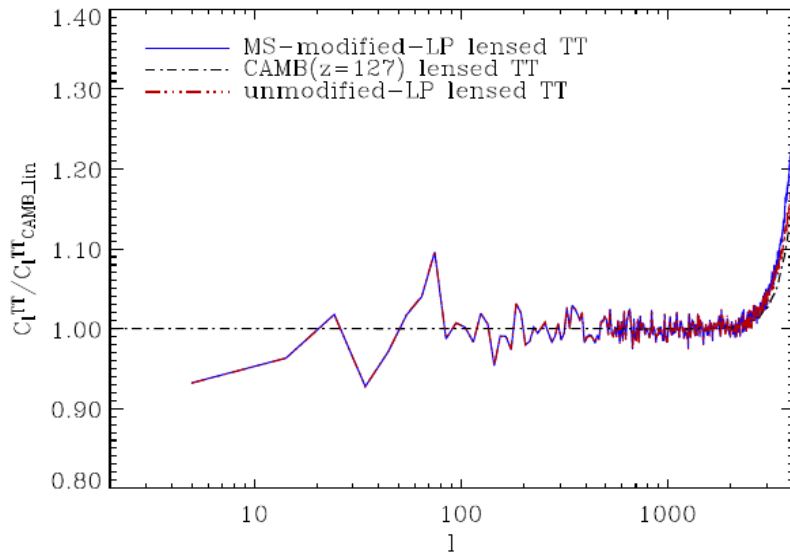


# CMB-N-body lensing

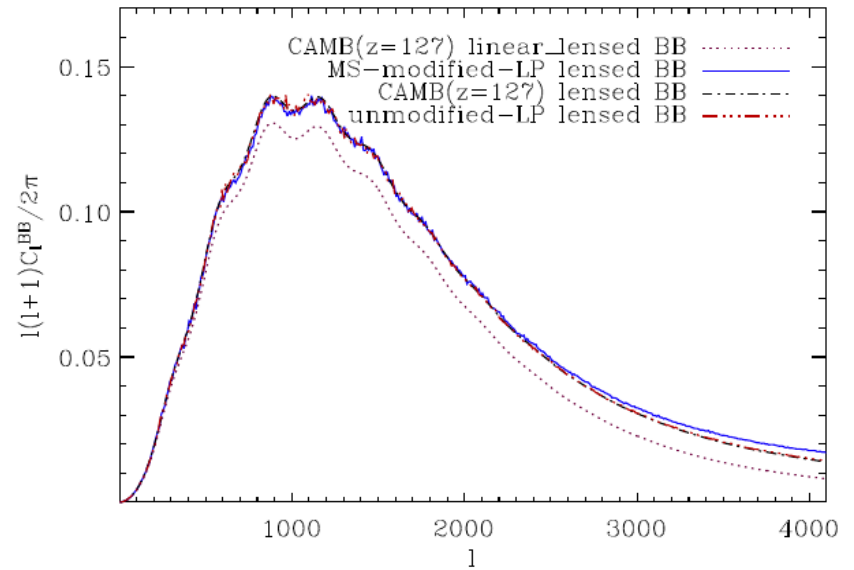
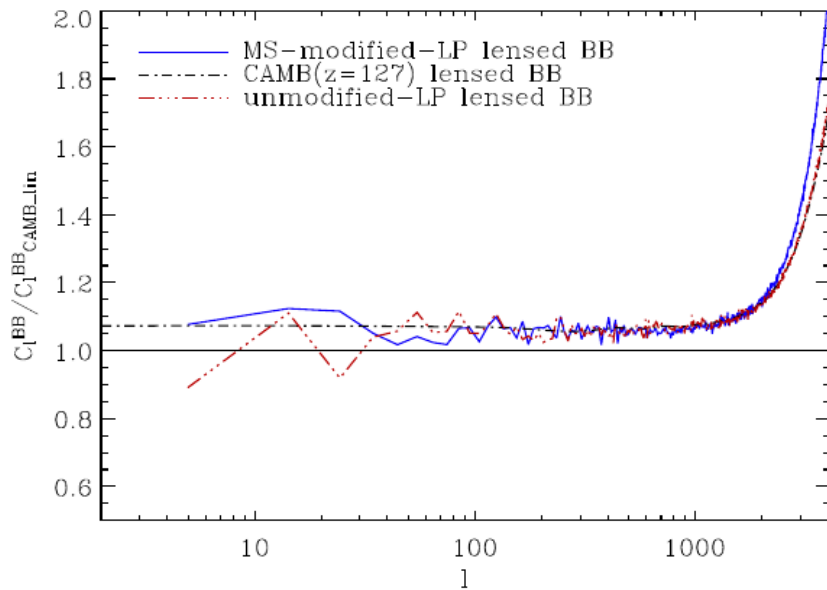
MS T difference map



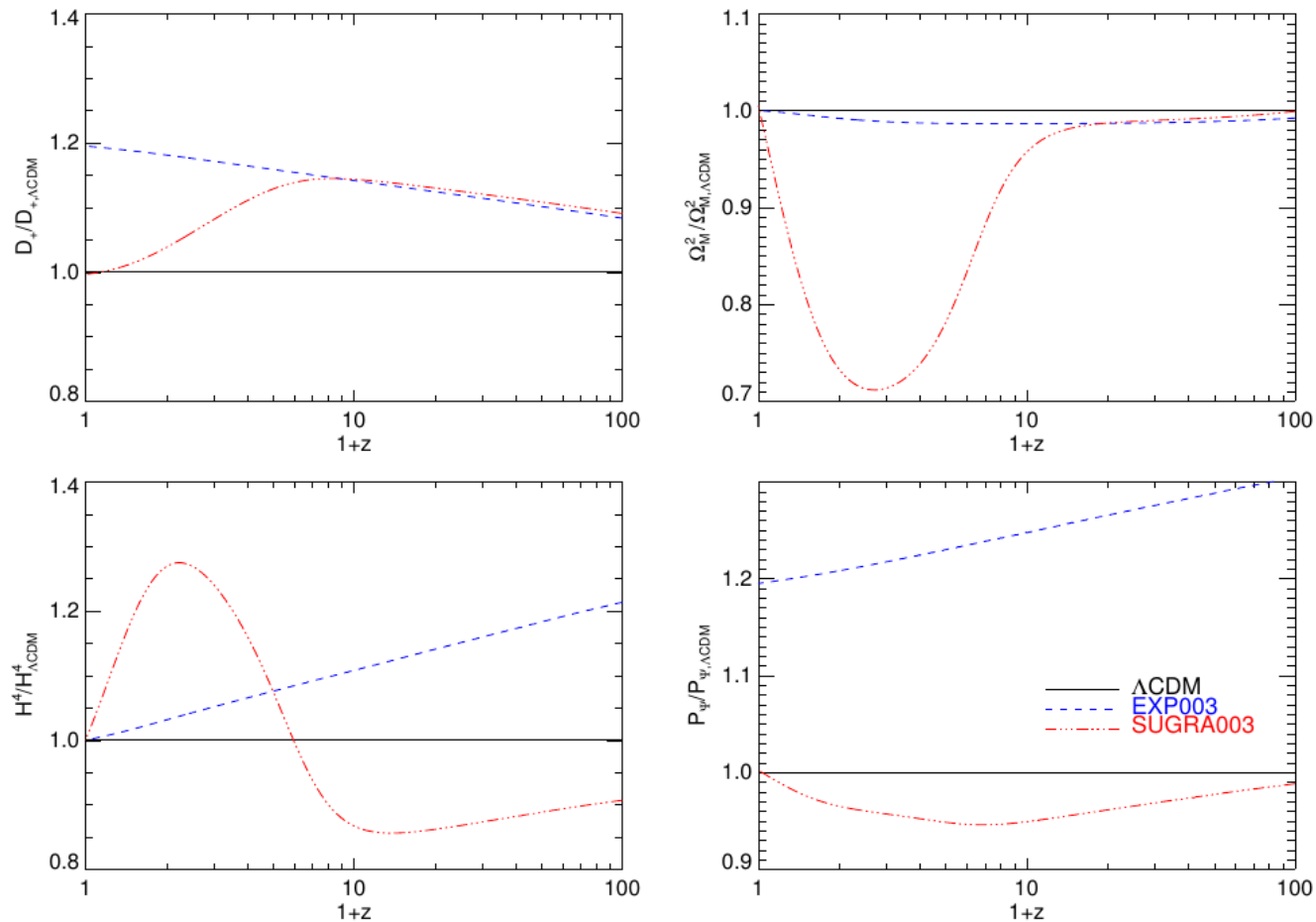
# CMB-N-body lensing



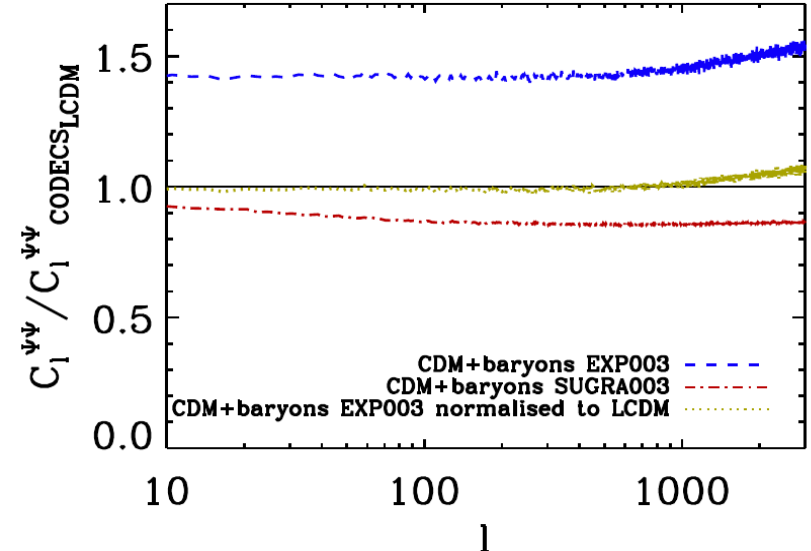
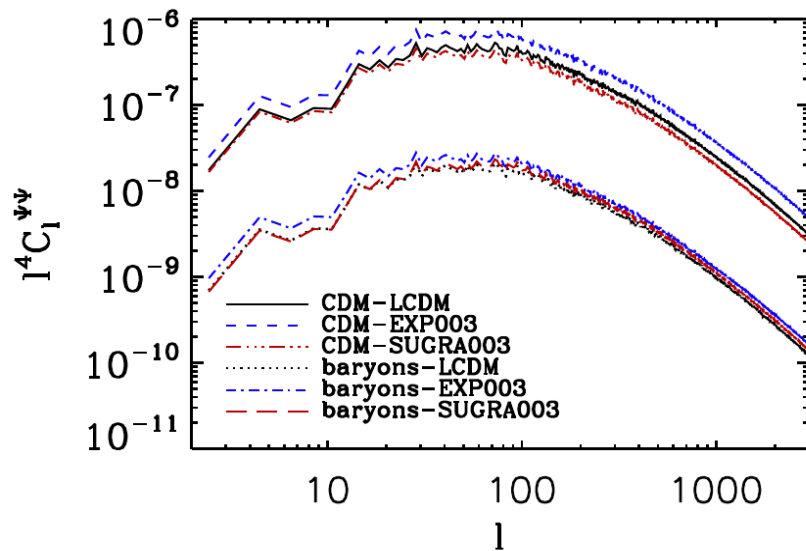
# CMB-N-body lensing



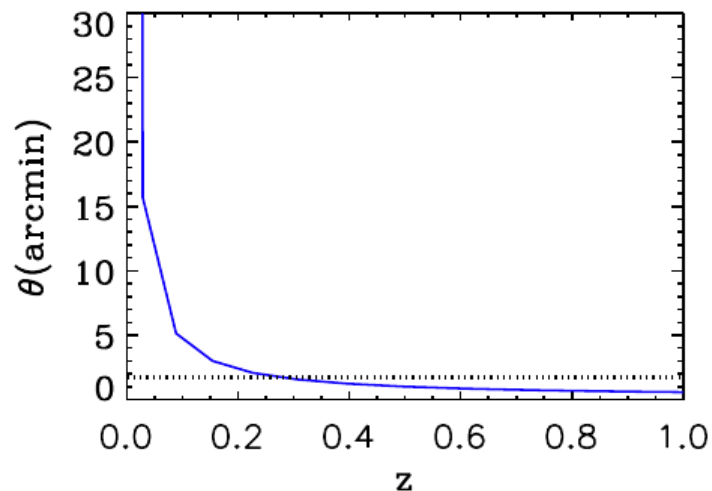
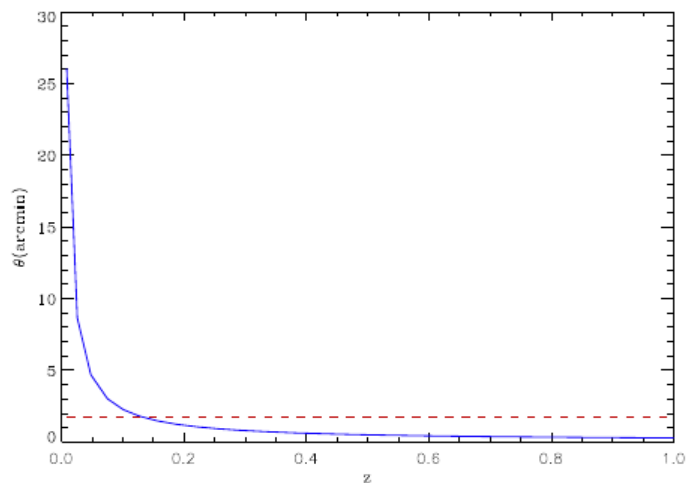
- CMB-N-body lensing: first inspection of non- $\Lambda$ CDM cosmologies



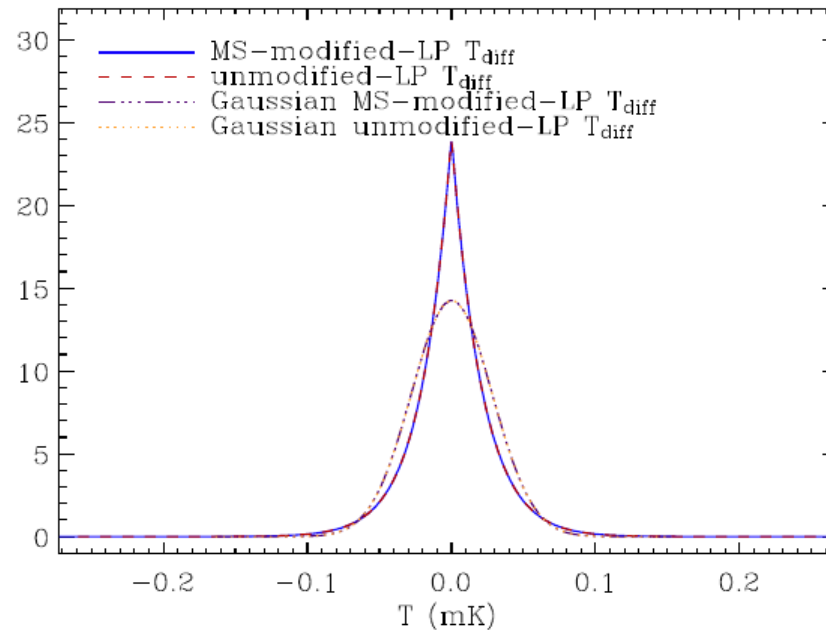
# CMB-N-body lensing: first inspection of non-LCDM cosmologies



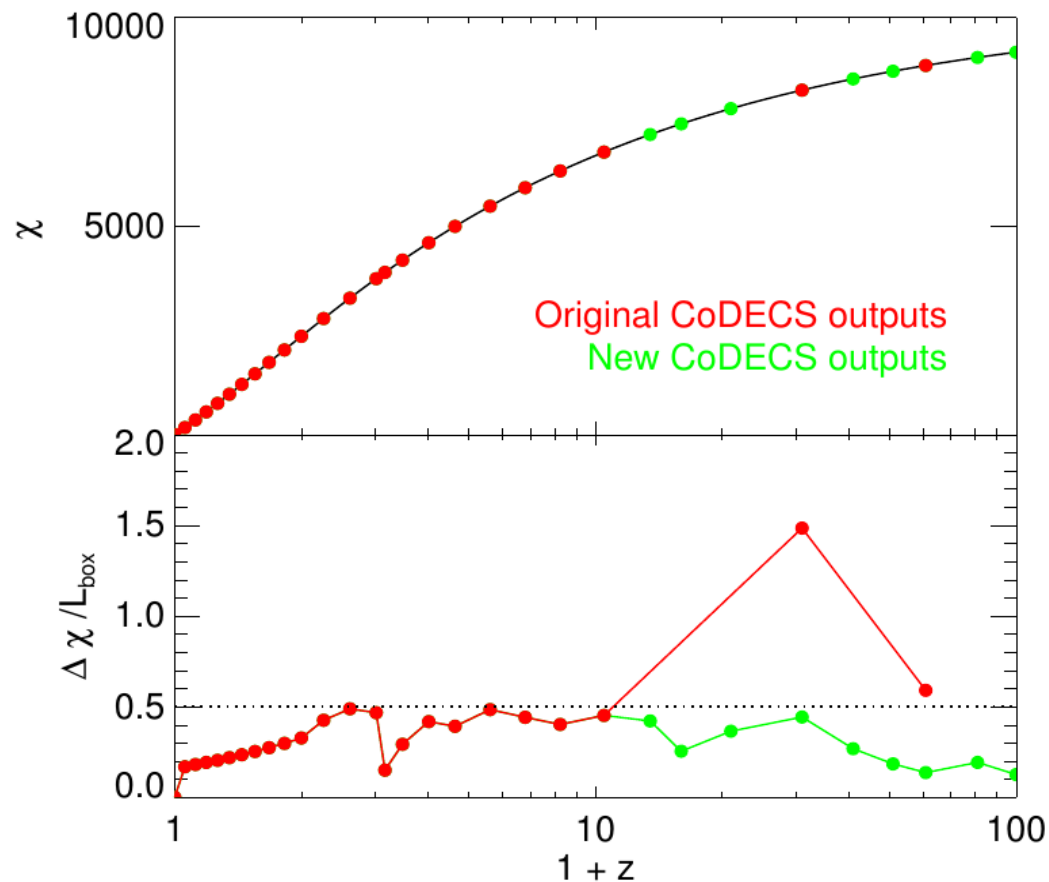
# CMB-N-body lensing: gridding issues



# CMB-N-body lensing: skewed statistics in the CMB lensed/non-lensed maps



# CMB-N-body lensing: learning lessons on N-body



# CMB lensing and sub-orbital CMB probes: extraction

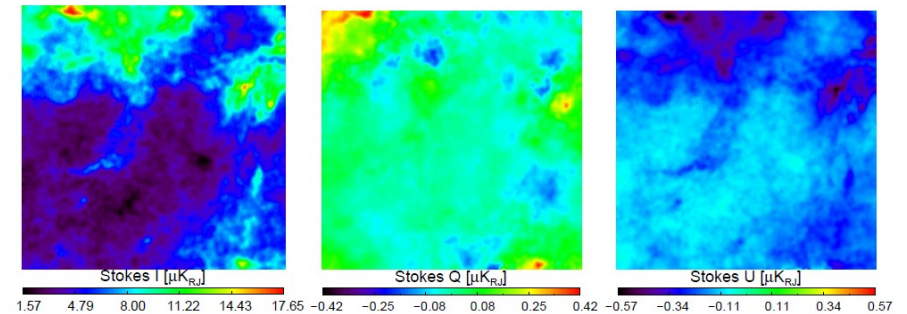
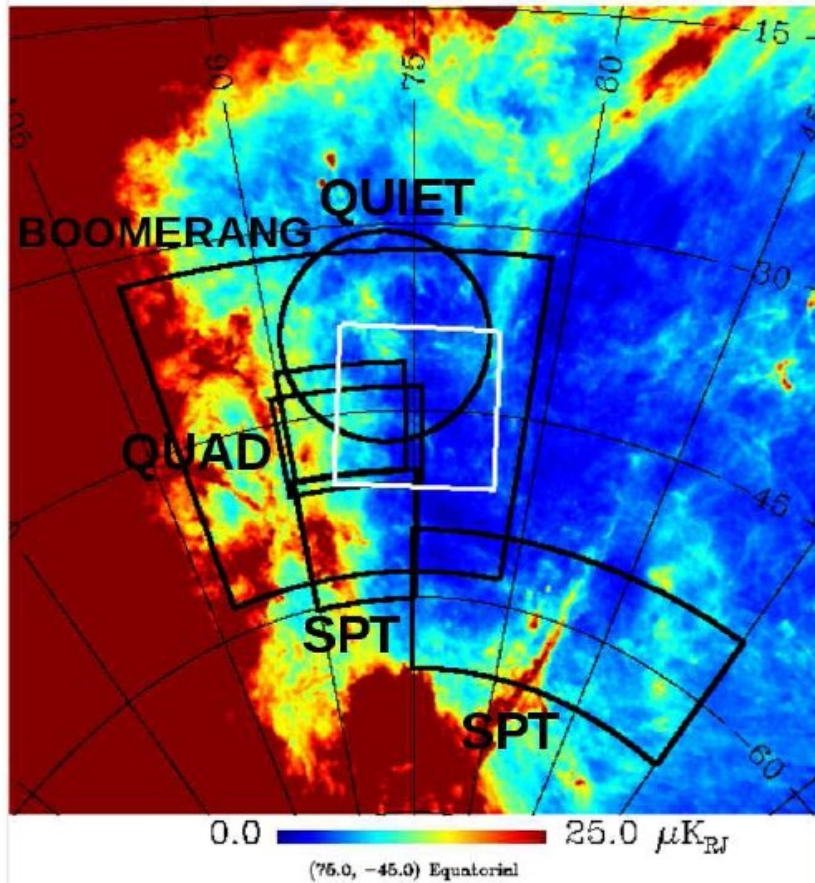


Figure 2. Dust simulation (Stokes  $I$ ,  $Q$ , and  $U$ ) at 150 GHz, with a polarization fraction  $p = 0.036$ , on a  $13^\circ \times 13^\circ$  patch centred on  $(l, b) = (250^\circ, -38^\circ)$ .

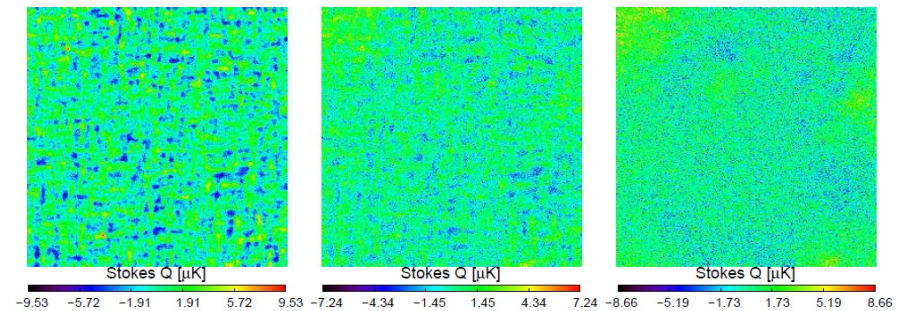
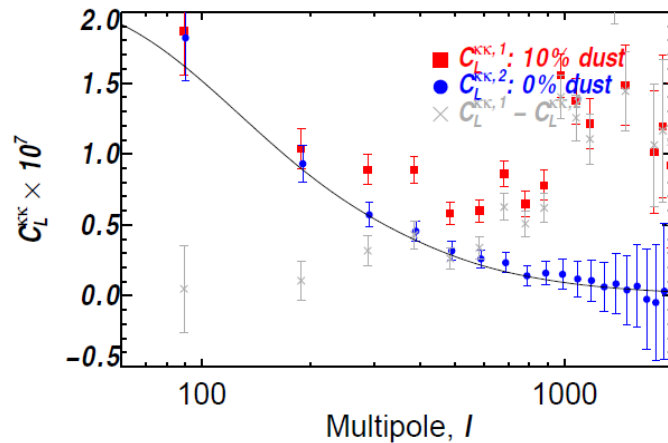
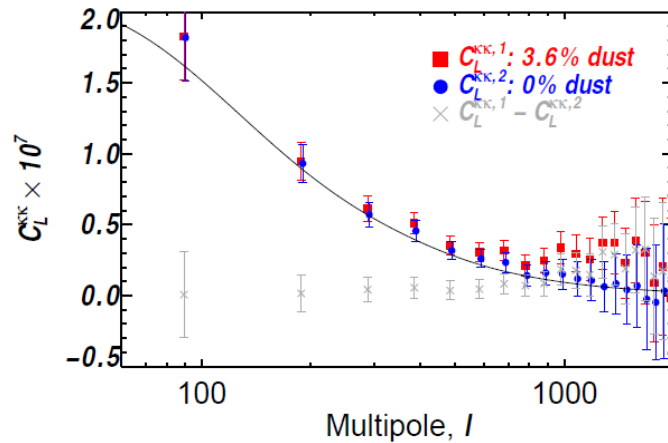
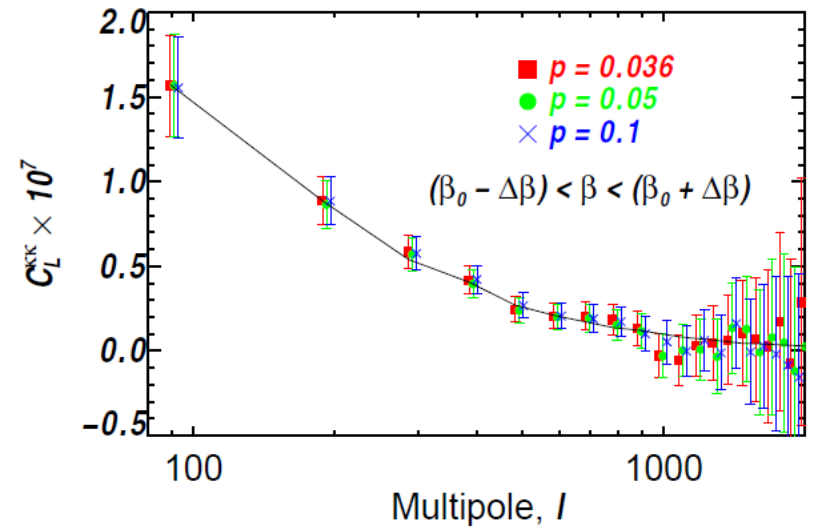


Figure 3. From left to right, we show our CMB+dust+noise simulations at 150, 250 and 410 GHz for a dust polarization fraction  $p = 0.1$ . The 150 GHz channel is CMB dominated, while the 410 GHz channel is dust dominated. The 250 GHz channel provides information about the dust spectral index.

# CMB lensing and sub-orbital CMB probes: extraction

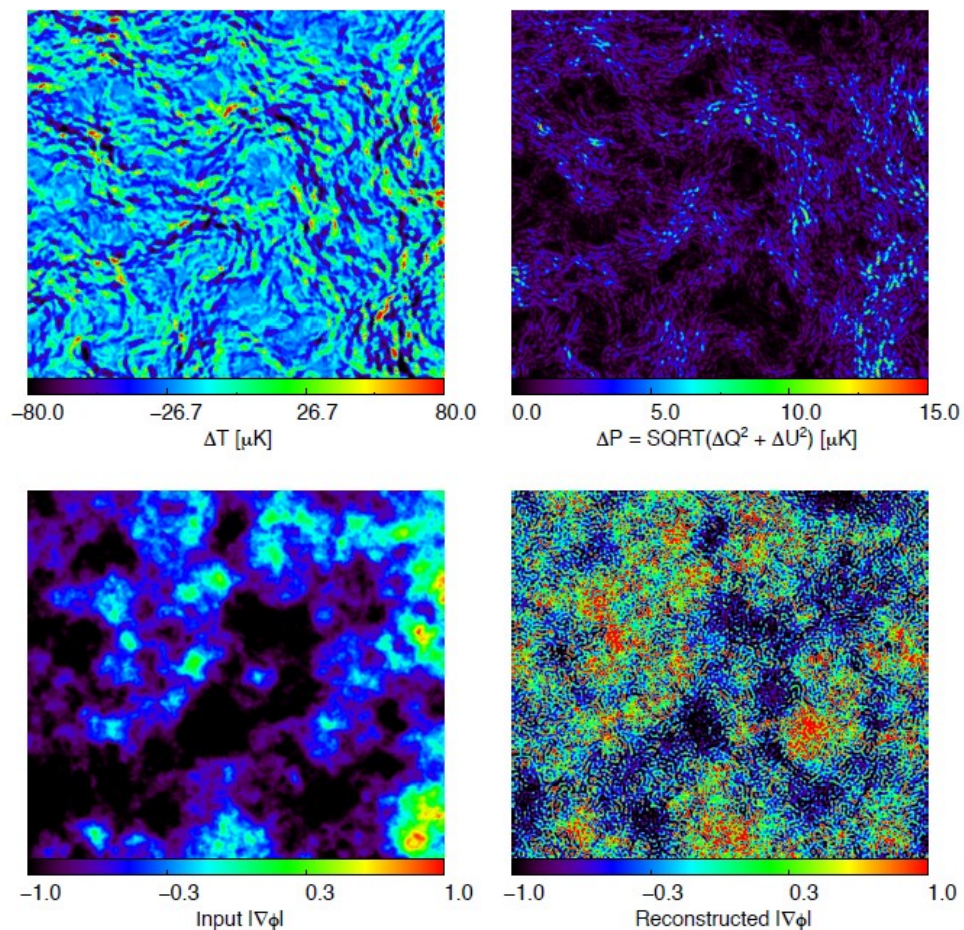


Convergence estimation  
Without foreground cleaning

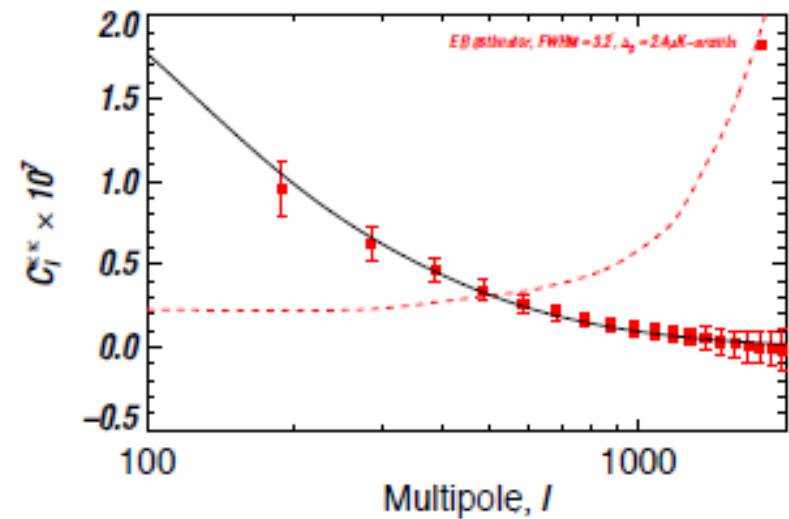
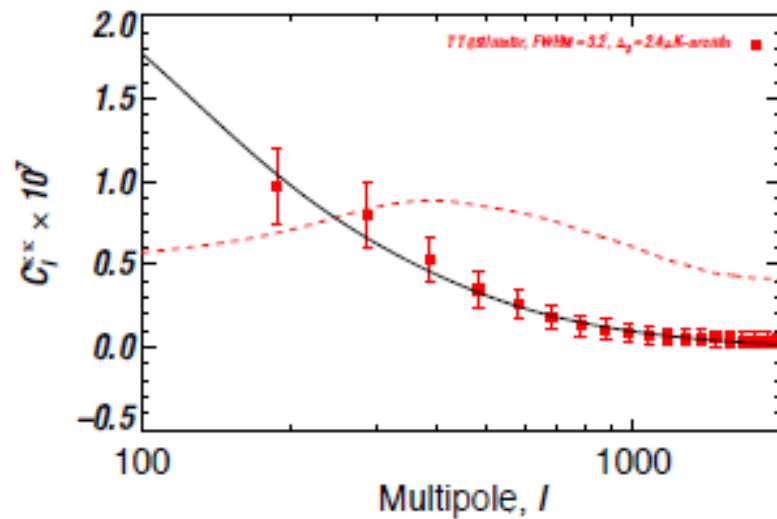


Convergence estimation  
With foreground cleaning

# CMB-N-body lensing reconstruction



# CMB-N-body lensing reconstruction



# Concluding remarks

- Is lensing relevant for Dark Energy?
- Does the signal keep its promise?
- Do we need CMB-N-body lensing?
- Do we have the technology for simulating it?
- Do we have a complete understanding of the effects of details in the simulations in the signal we see?
- Can we start simulating cross-correlation studies on the way to Euclid?
- ...

# Concluding remarks

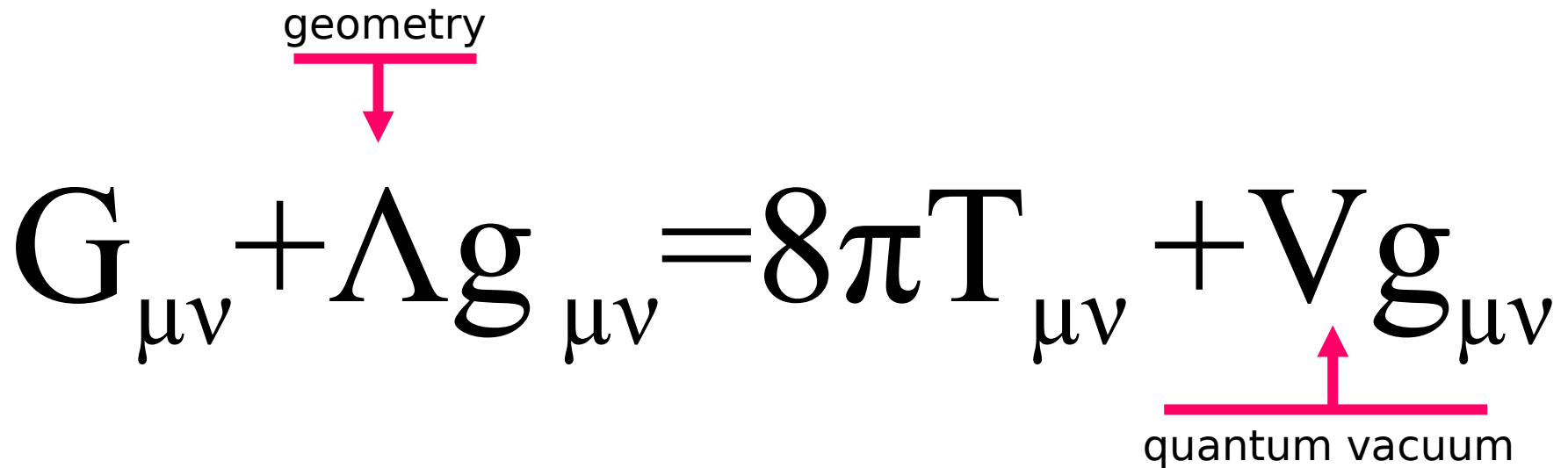
- Is lensing relevant for Dark Energy? **Yes, it is.**
- Does the signal keep its promise? **Yes, it does.**
- Do we need CMB-N-body lensing? **Well...yes.**
- Do we have the technology for simulating it? **Yes...still limited by present simulations, but it works.**
- Do we have a complete understanding of the effects of details in the simulations in the signal we see? **No. We still need to distinguish between physical effects/N-body details/gridding issues at high resolution.**
- Can we start simulating cross-correlation studies on the way to Euclid? **Oh yes, we should...**
- ...

# Extension II: A mini-Dark Energy Course

# Fighting the cosmological constant

$$G_{\mu\nu} = 8\pi T_{\mu\nu}$$

# Fighting the cosmological constant



The diagram illustrates the Einstein field equation,  $G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi T_{\mu\nu} + V g_{\mu\nu}$ , with annotations. A pink arrow points from the word "geometry" to the left-hand side of the equation. Another pink arrow points from the word "quantum vacuum" to the term  $V g_{\mu\nu}$  on the right-hand side.

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi T_{\mu\nu} + V g_{\mu\nu}$$

geometry

quantum vacuum

# Fighting the cosmological constant

$\Lambda$ :???

# Fighting the cosmological constant

$$V:M^4_{\text{Planck}} ???$$

# Fighting the cosmological constant

$\Lambda: ???$

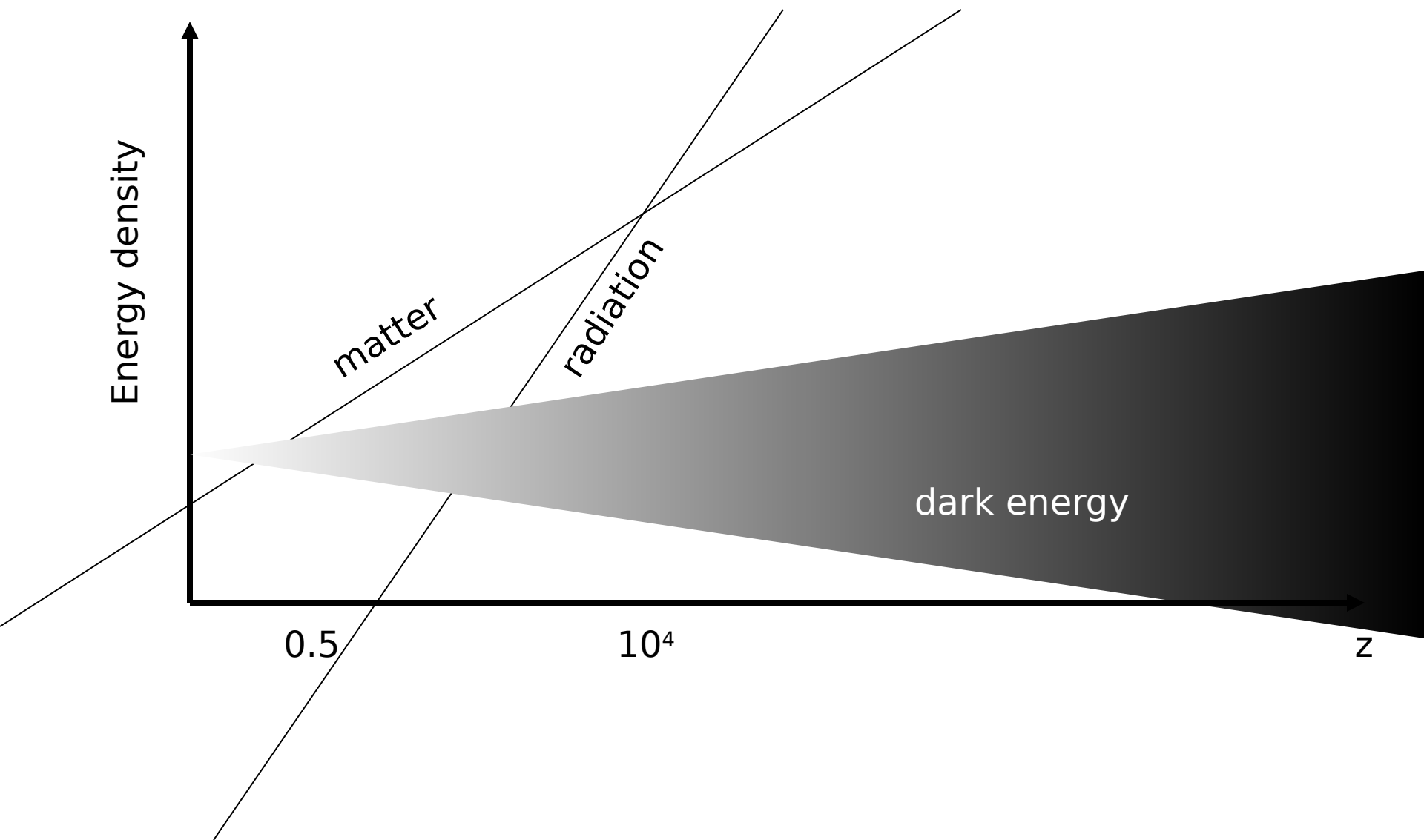
$$|\Lambda - V|/M_{\text{Planck}}^4 = 10^{-123}$$

$V/M_{\text{Planck}}^4 ???$

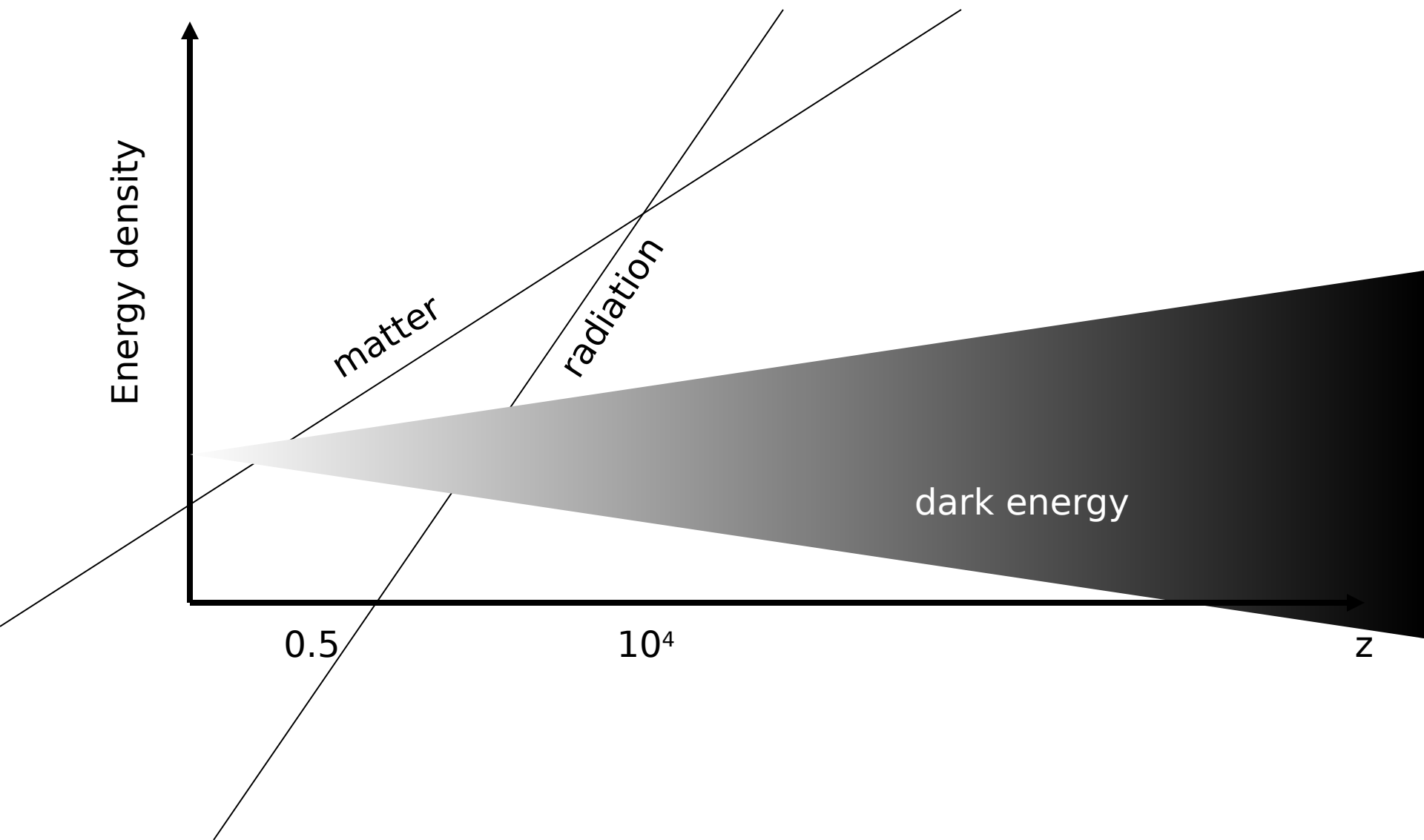
$(\text{Boh?})^2$

- Why so small with respect to any other known energy scale in physics?
- Why comparable to the matter energy density today?

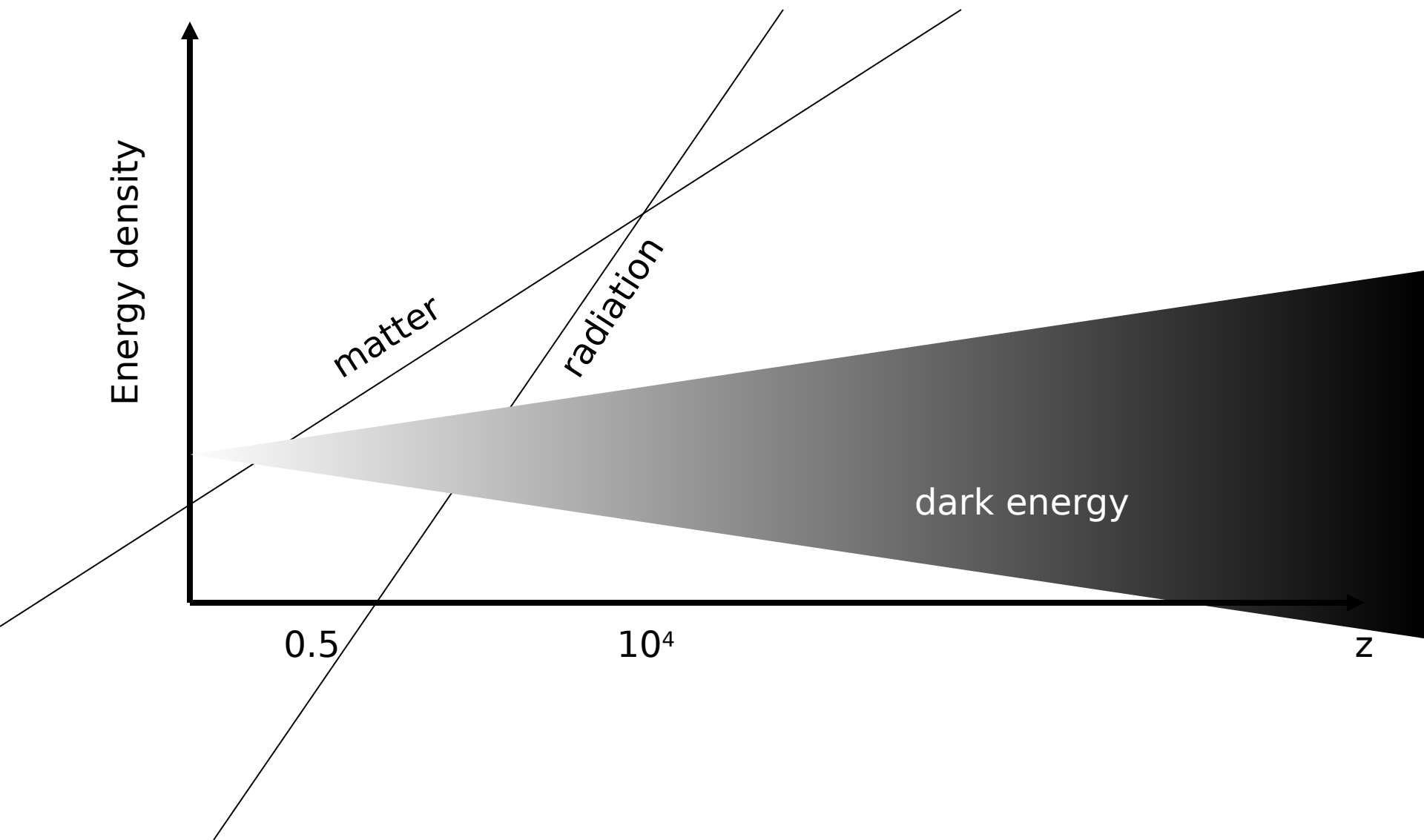
# Dark Energy



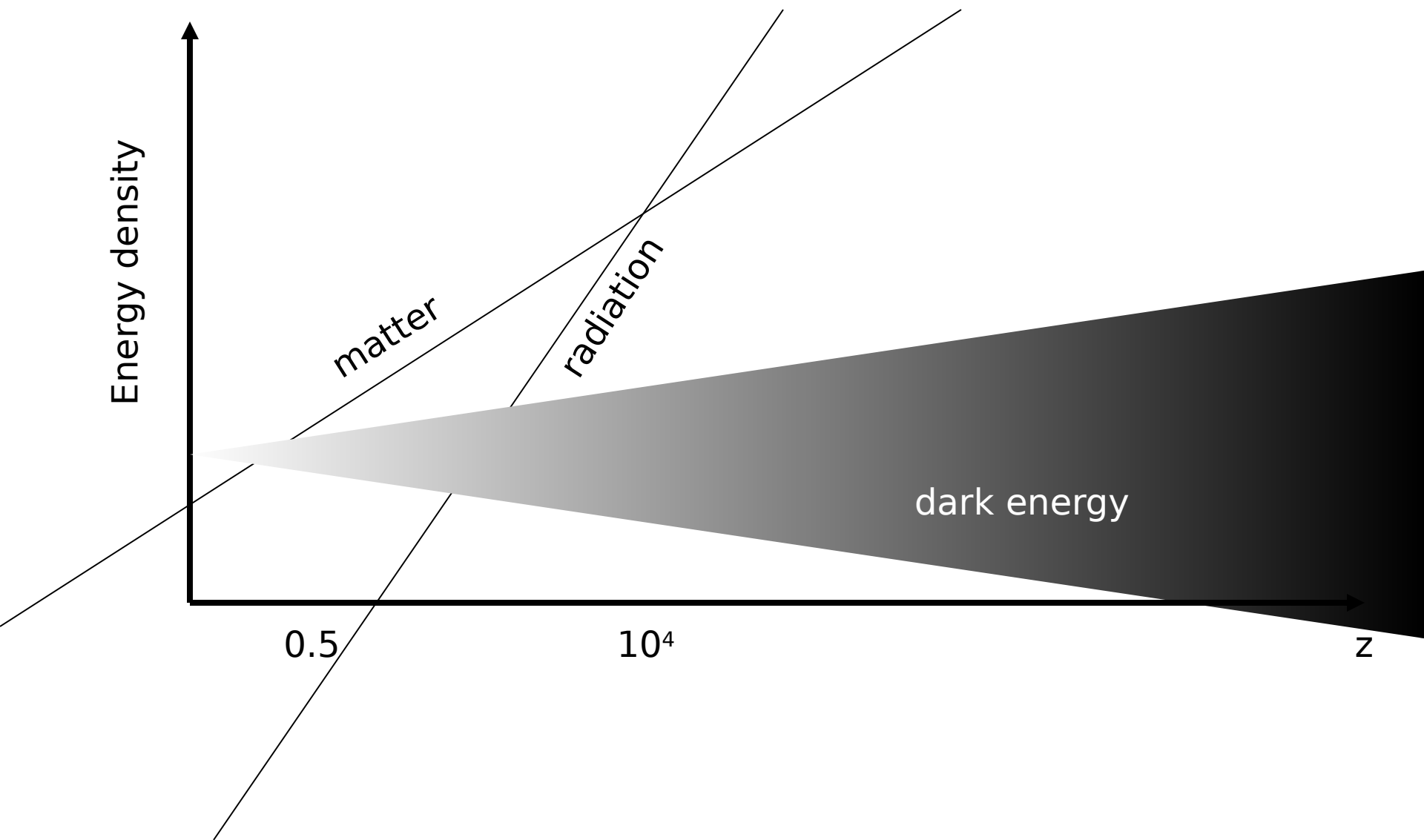
# Dark Energy



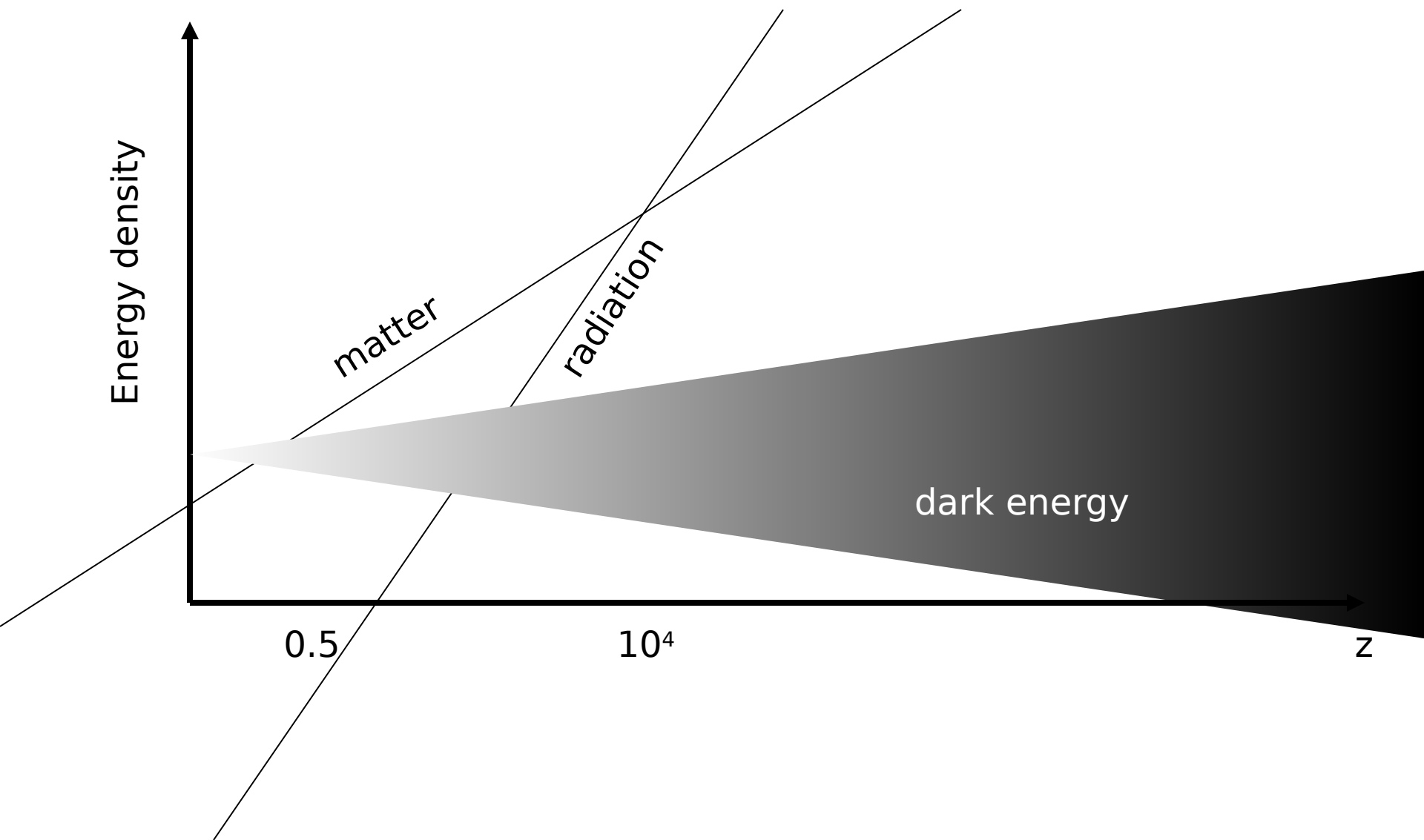
# Dark Energy



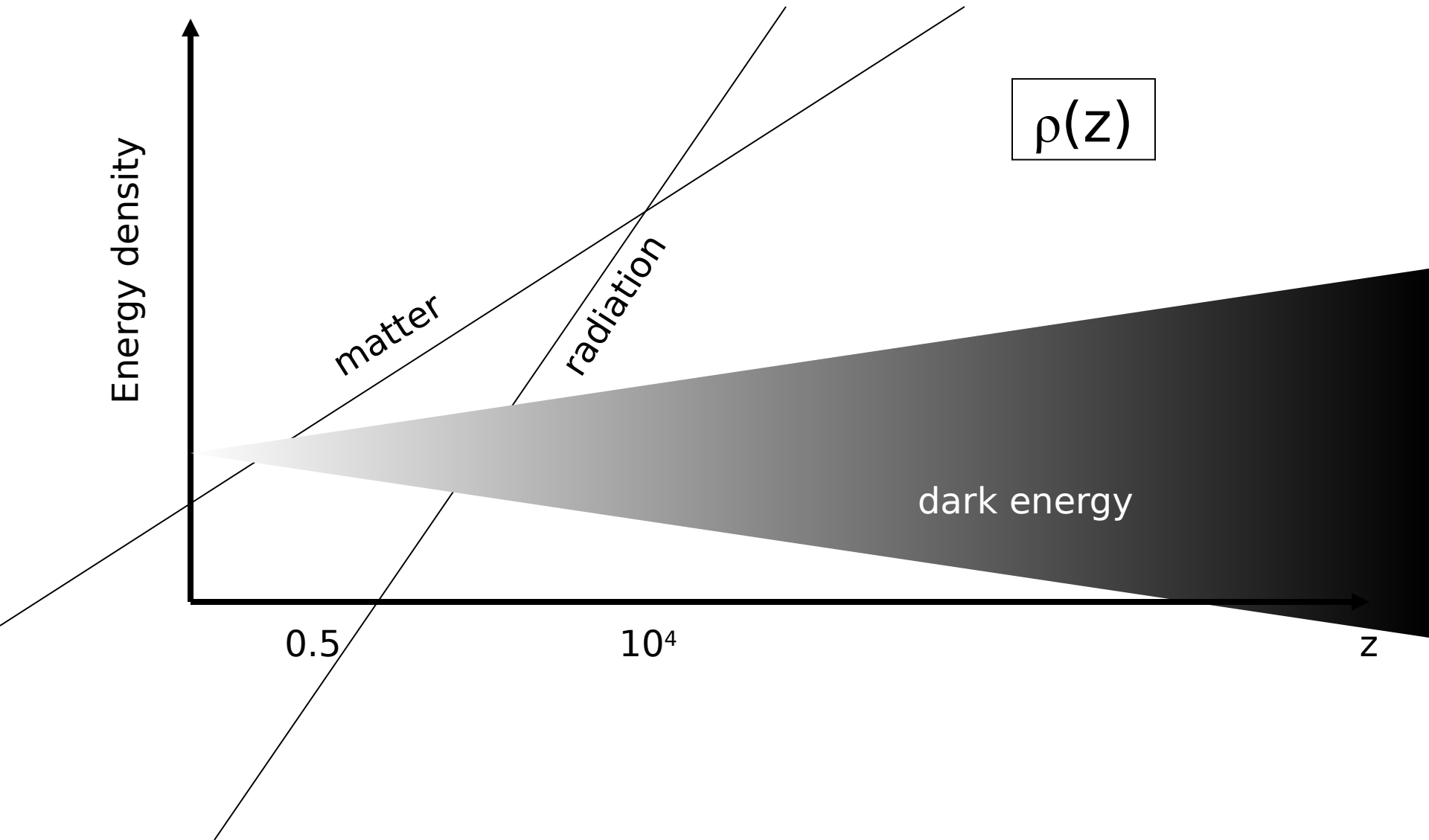
# Dark Energy



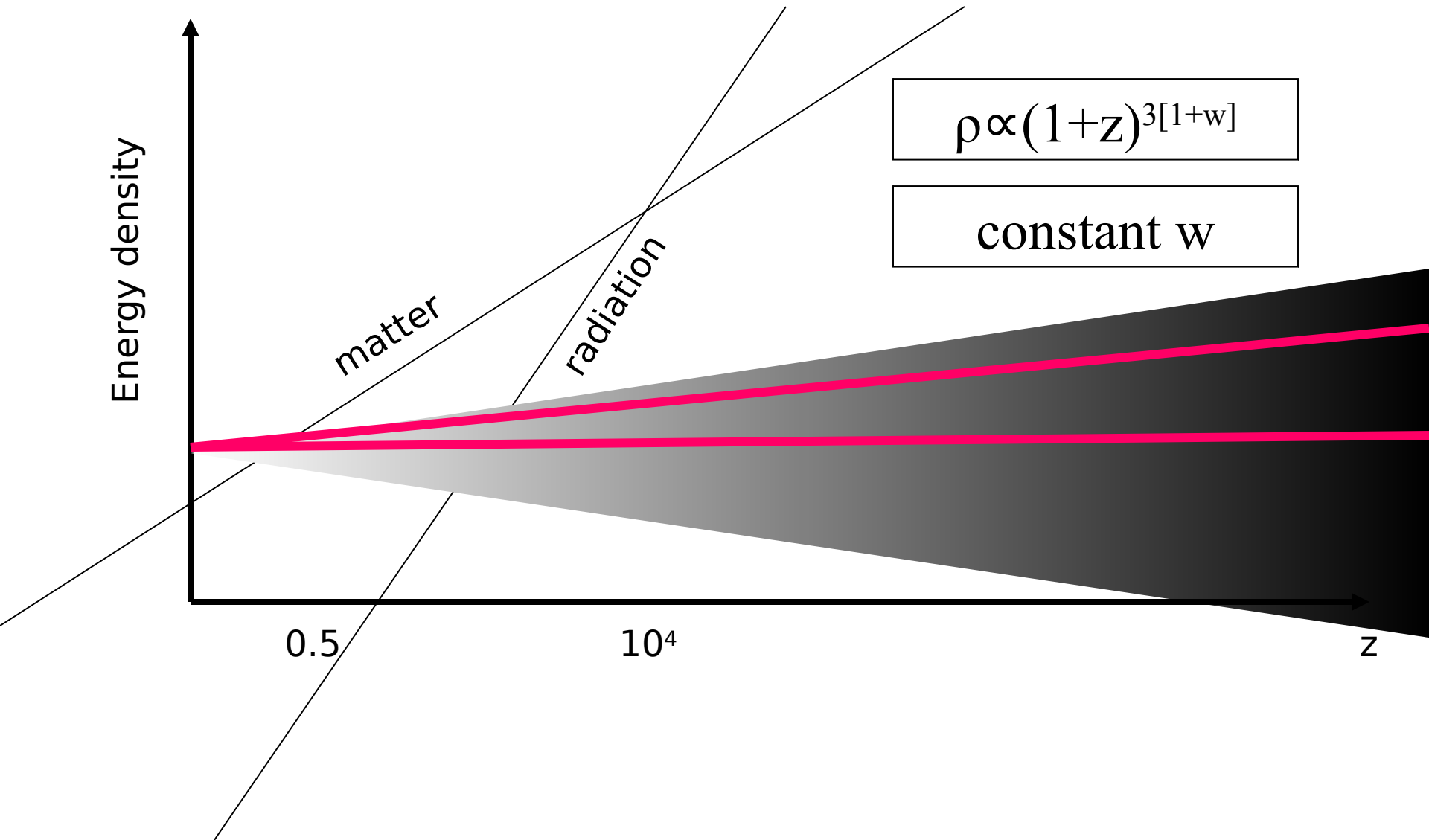
# Dark Energy



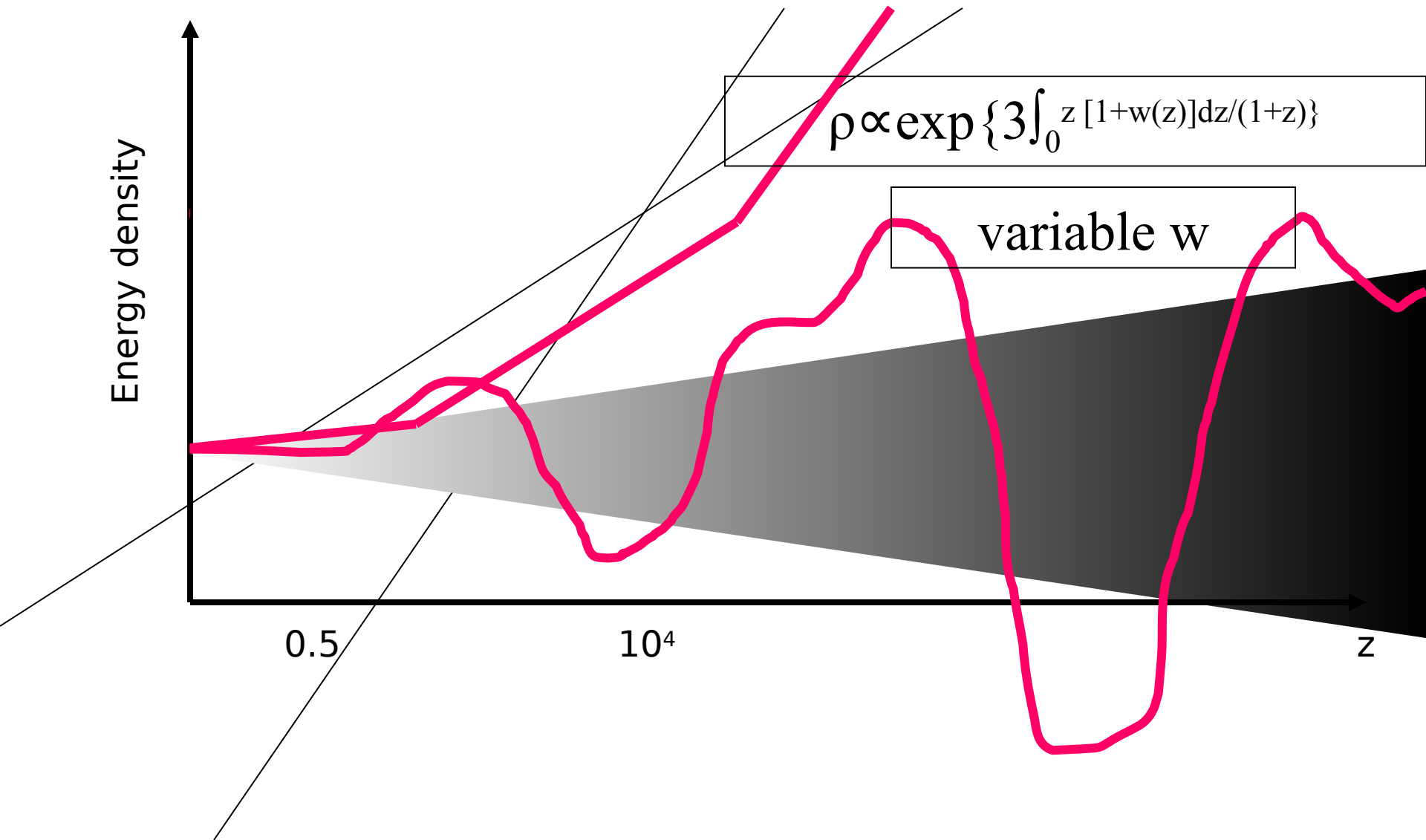
# Parametrizing cosmic acceleration is ...



# ...parametrizing cosmic density

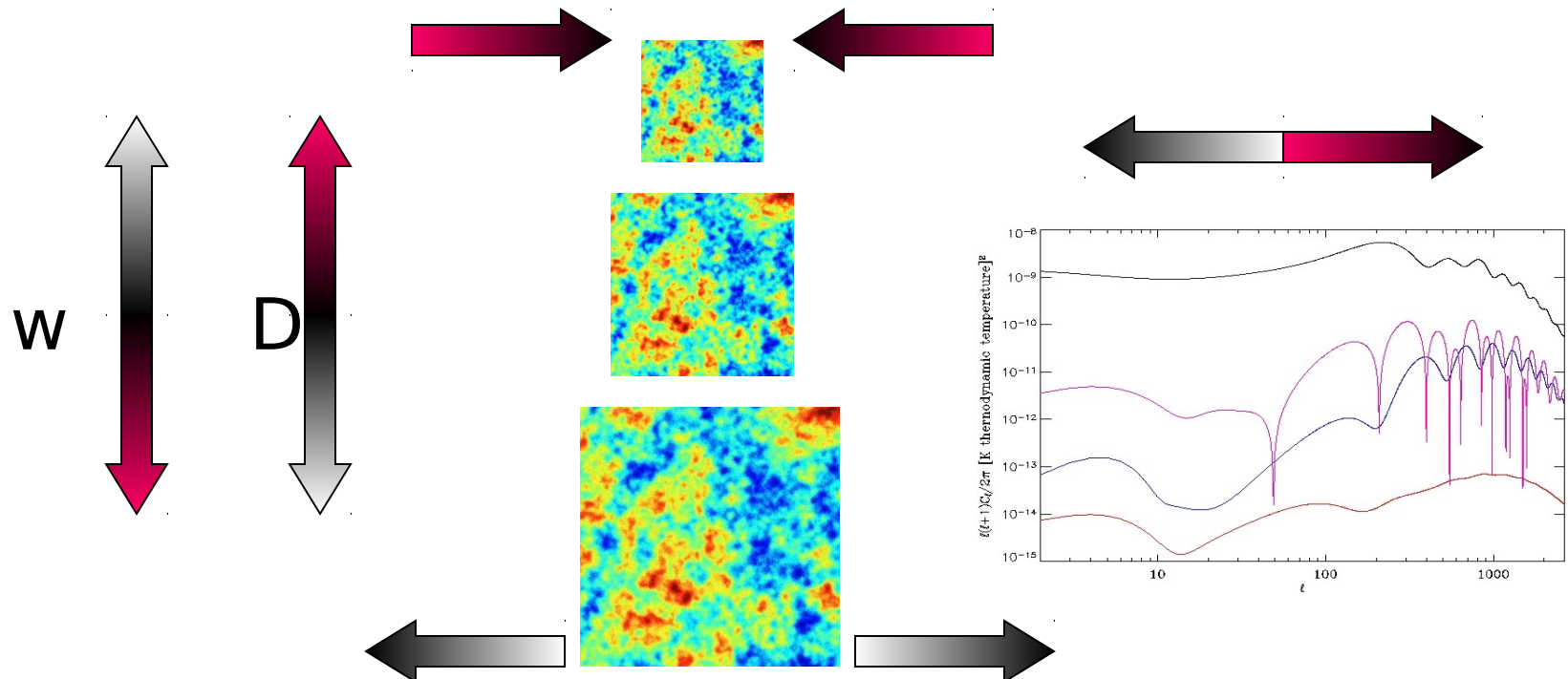


# Parametrizing cosmic density



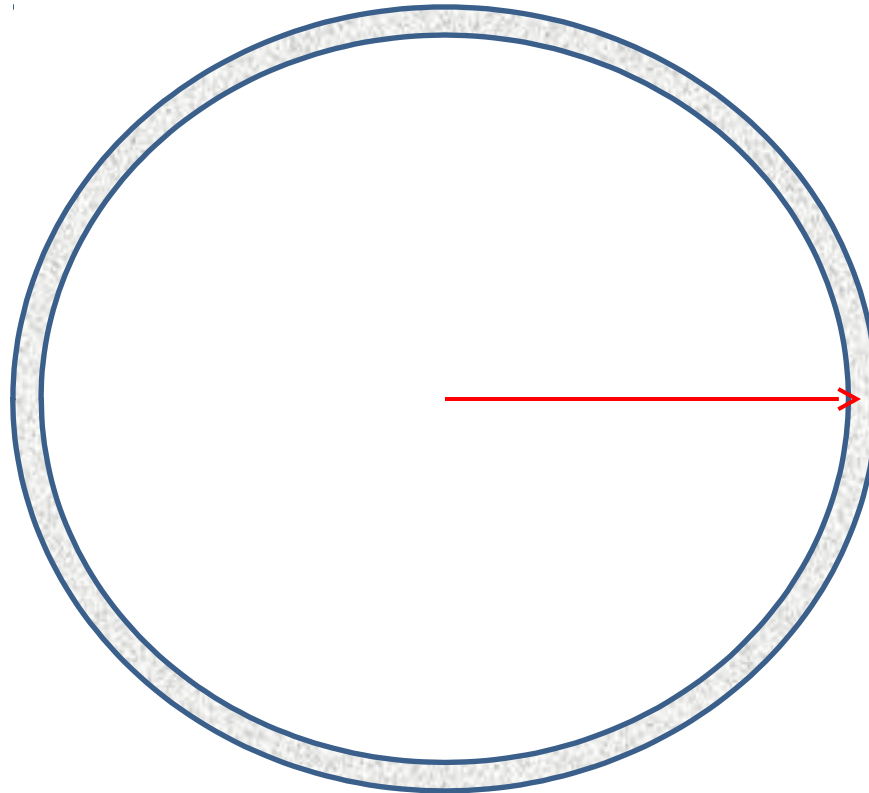
# Cosmological expansion and CMB: projection

$$D = H_0^{-1} \int_0^z \frac{dz}{[\sum_i \Omega_i (1+z)^{3(1+w_i)}]^{1/2}}$$



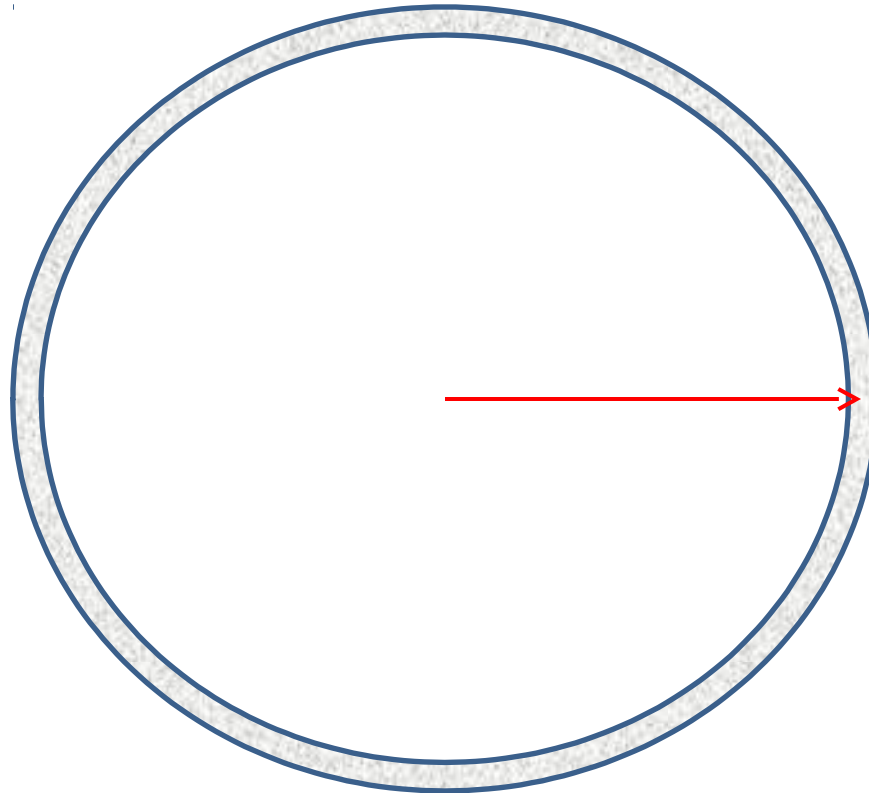
# Constraining the expansion with distances

$$D = H_0^{-1} \int_0^z \frac{dz}{[\sum_i \Omega_i (1+z)^{3(1+w_i)}]^{1/2}}$$



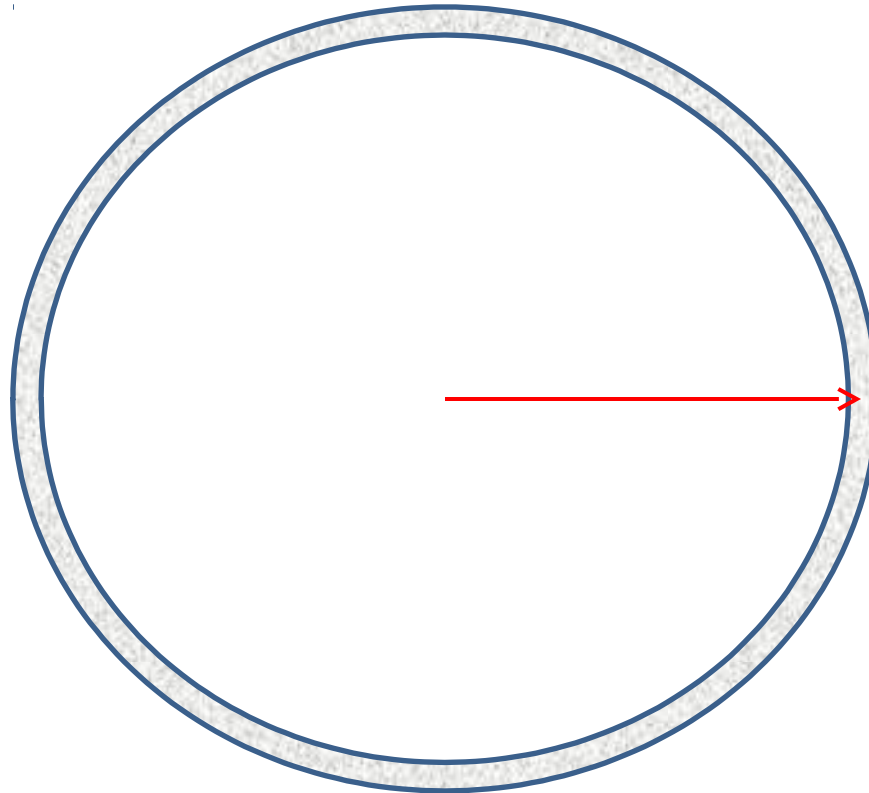
# Constraining the expansion with distances

$$D = H_0^{-1} \int_0^z \frac{dz}{[\sum_i \Omega_i (1+z)^{3(1+w_i)}]^{1/2}}$$

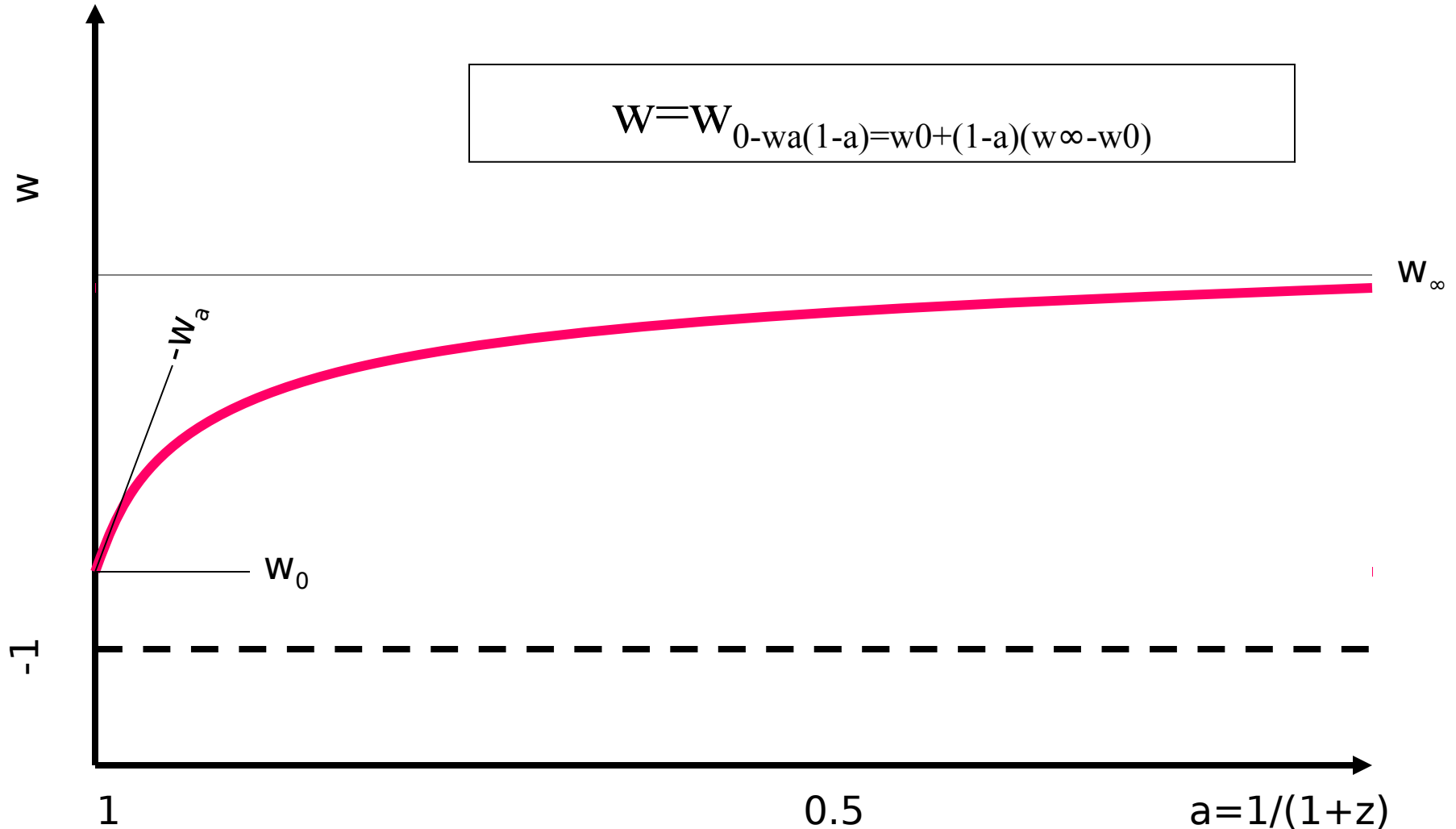


# Constraining the expansion with distances

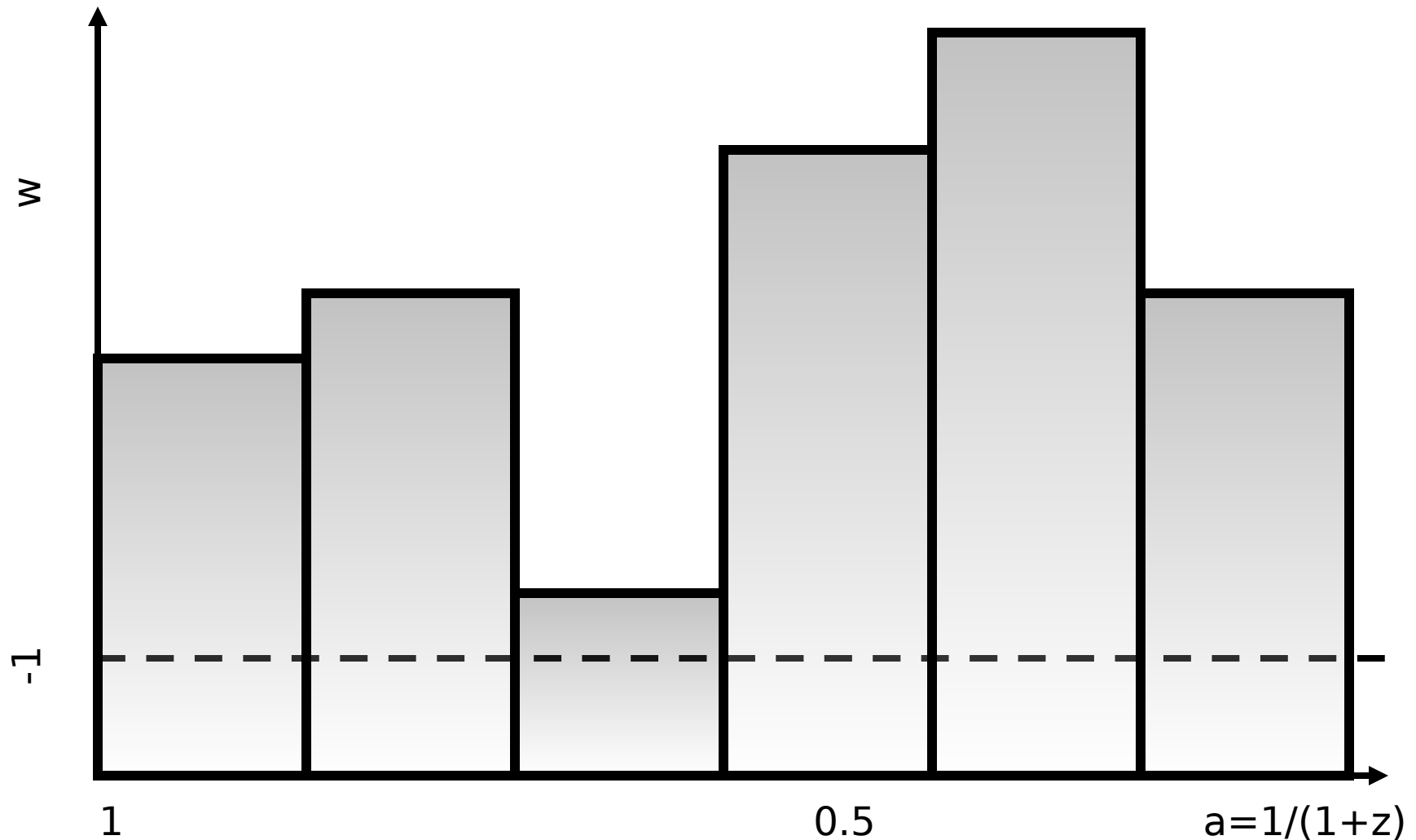
$$D = H_0^{-1} \int_0^z \frac{dz}{[\sum_i \Omega_i (1+z)^{3(1+w_i)}]^{1/2}}$$



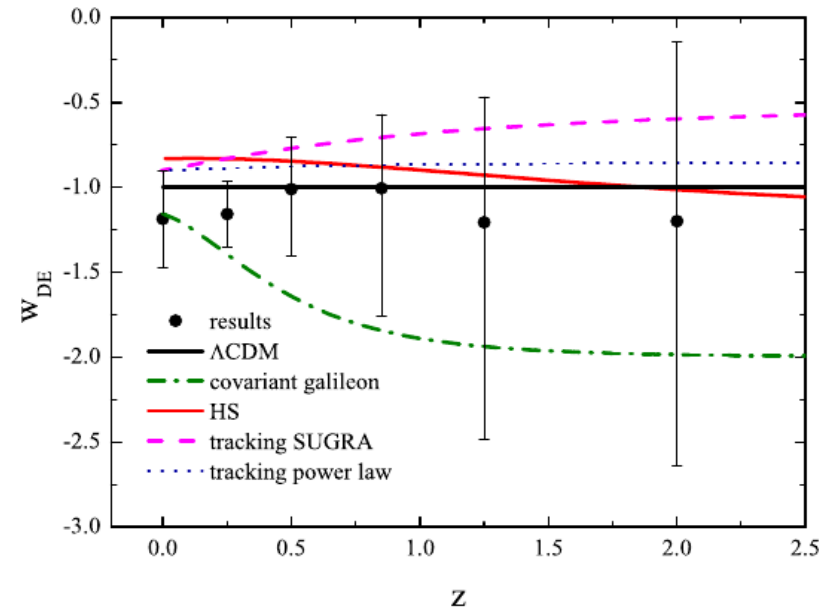
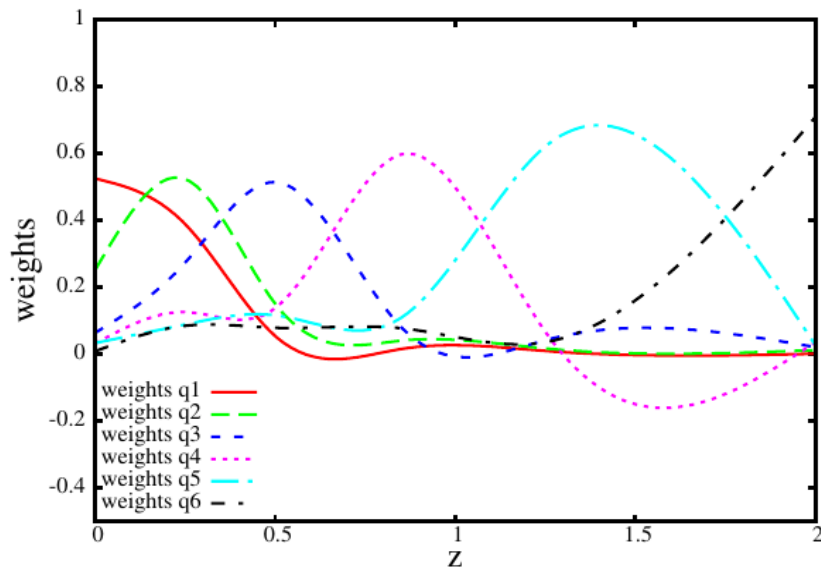
# Constraining Dark Energy: models



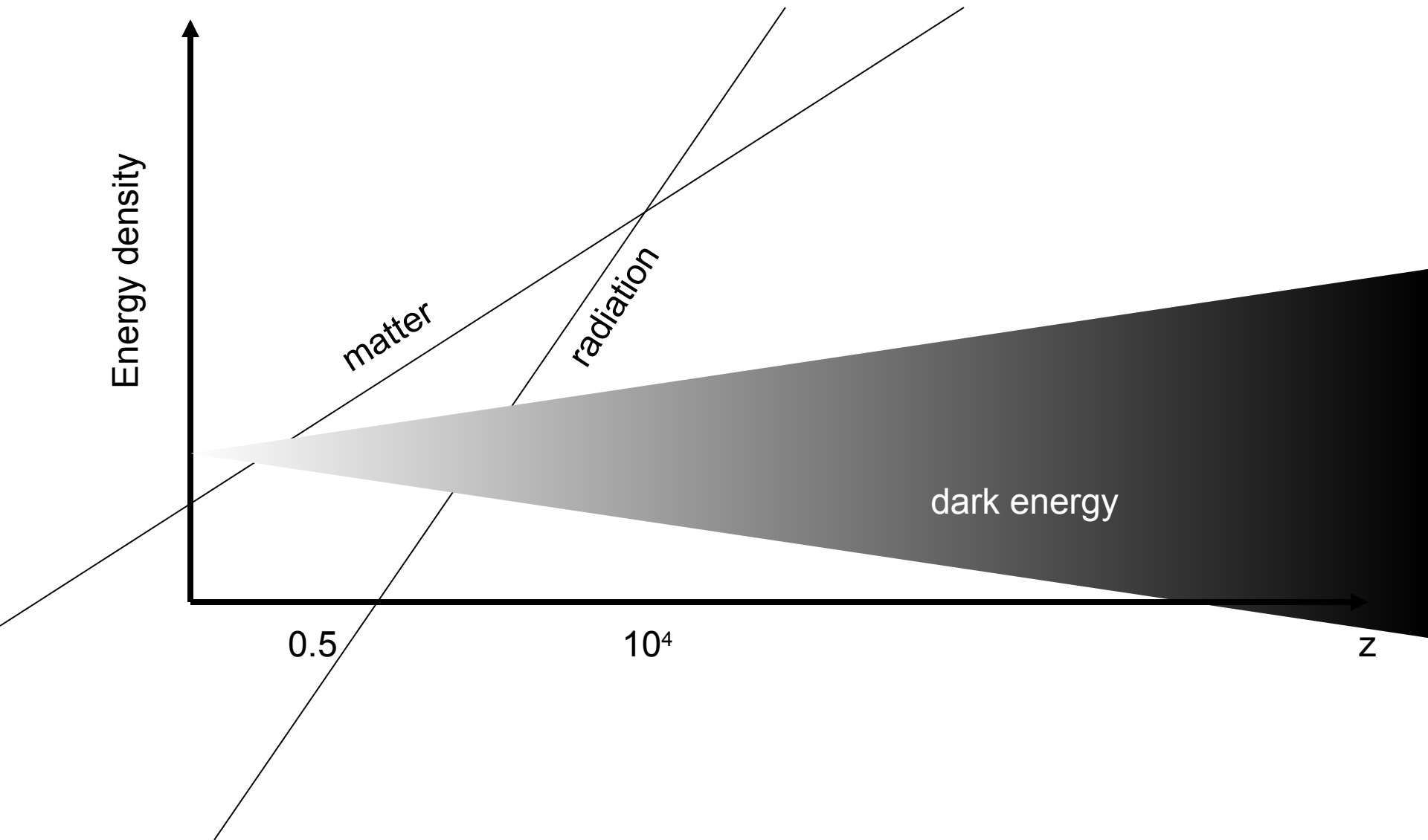
# Constraining Dark Energy: binning



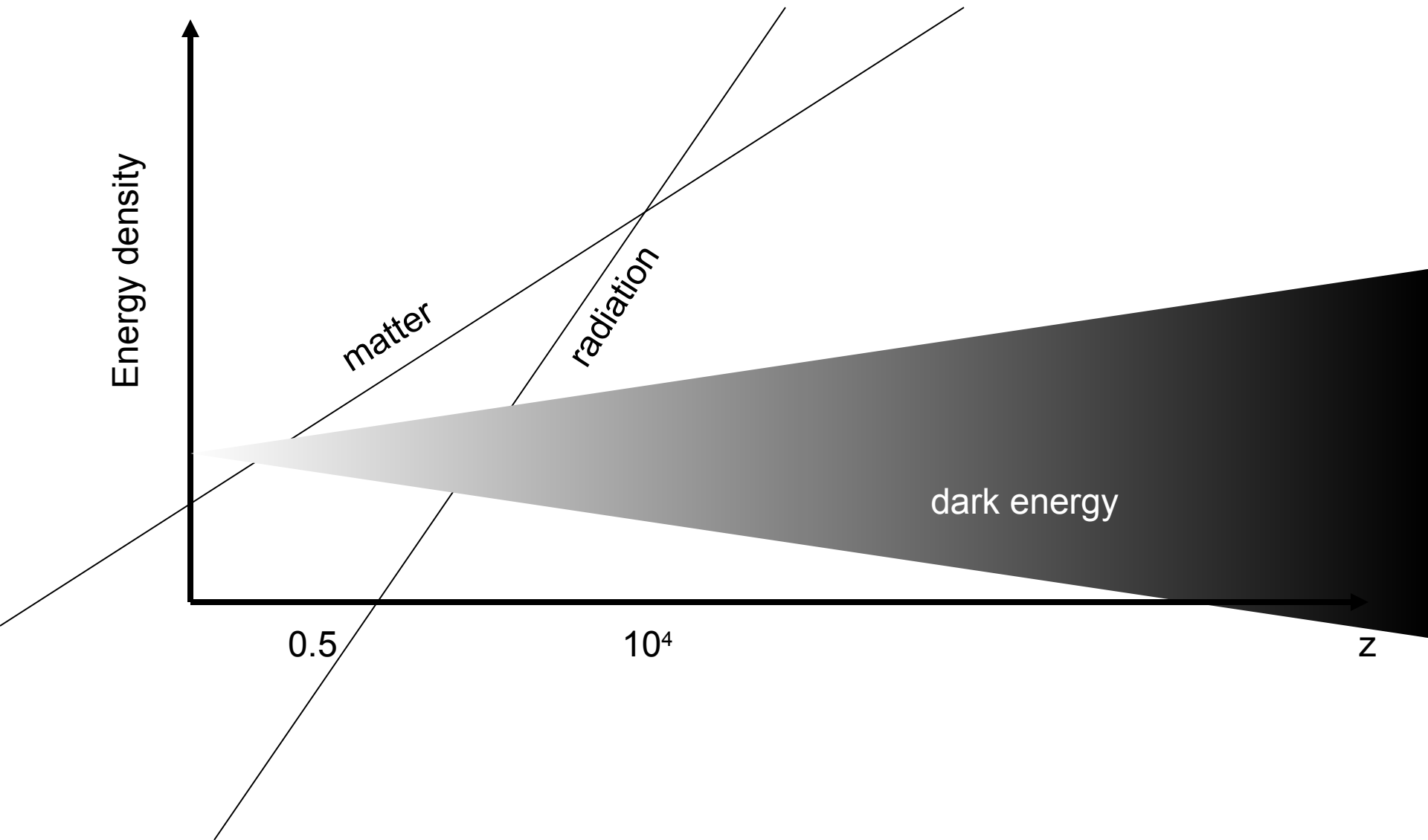
# Constraining Dark Energy: latest pre-Planck constraints



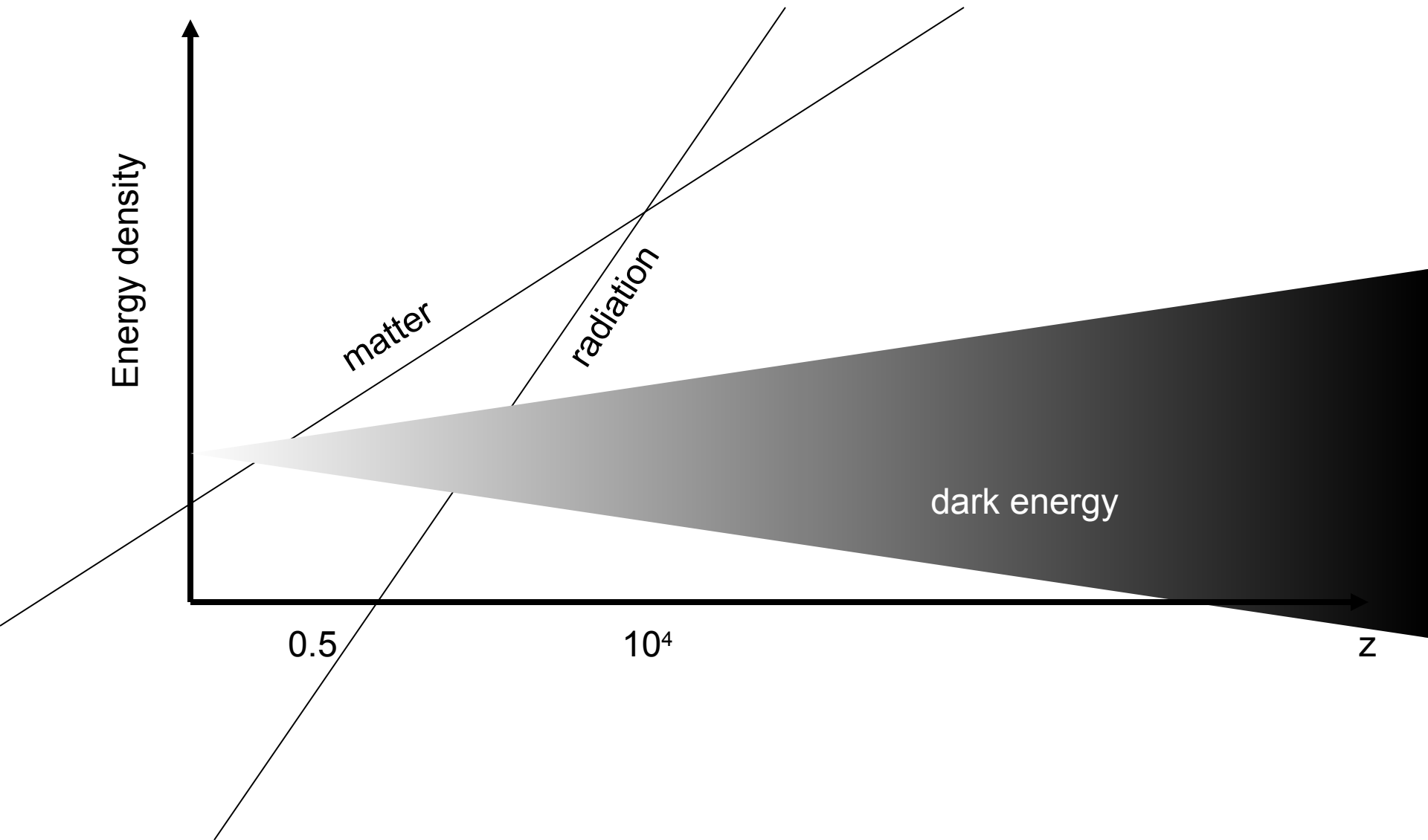
# The promise of CMB lensing



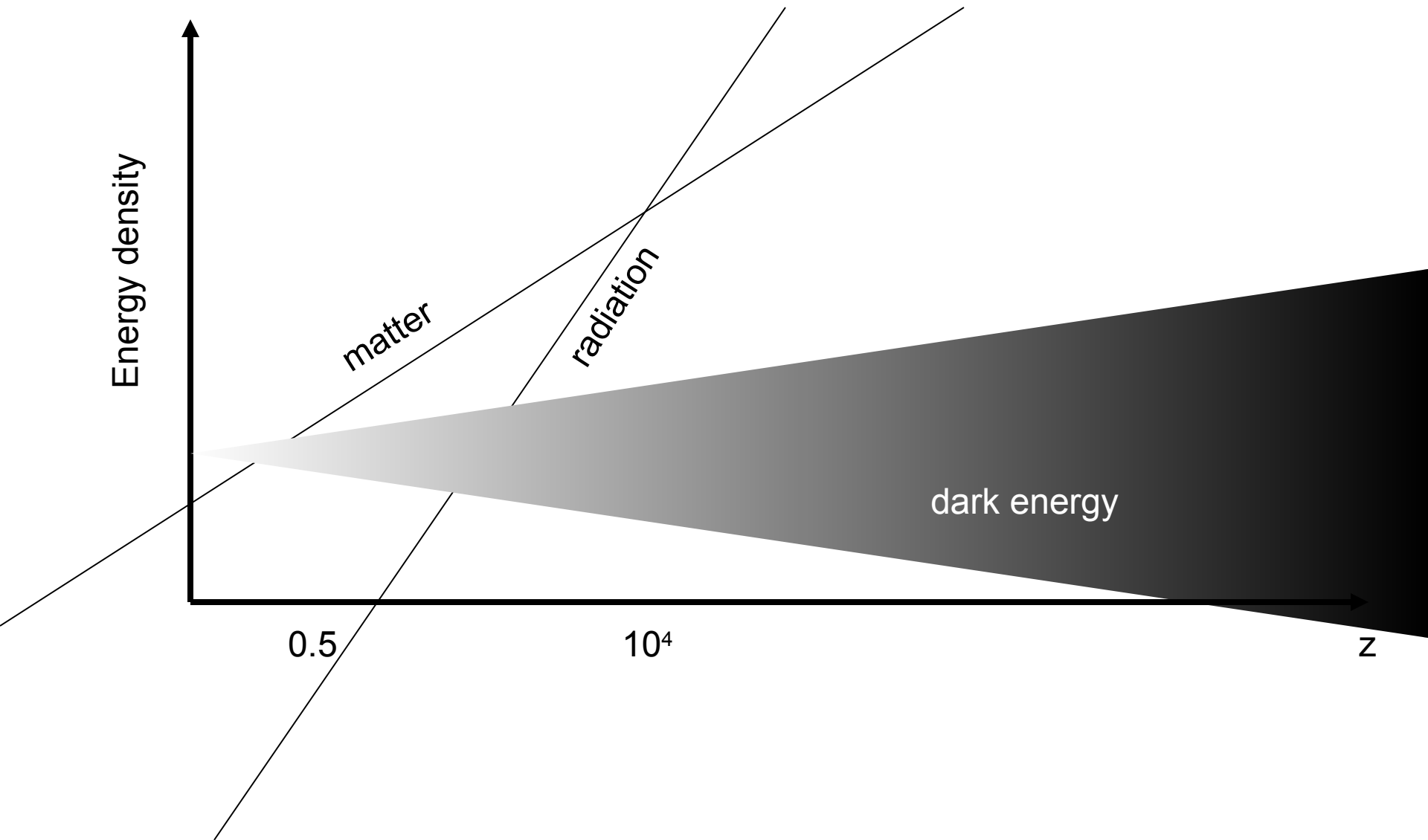
# The promise of CMB lensing



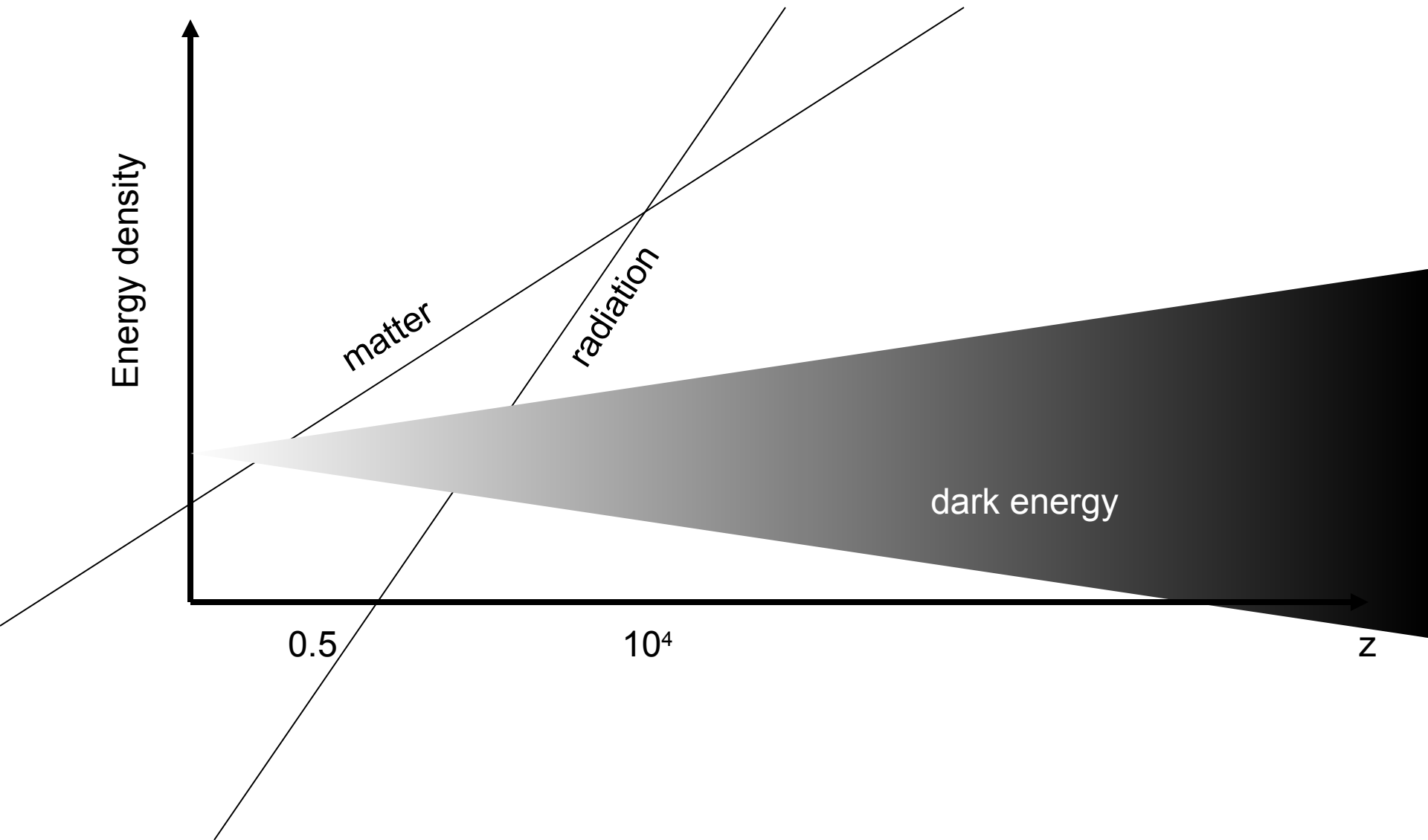
# The promise of CMB lensing



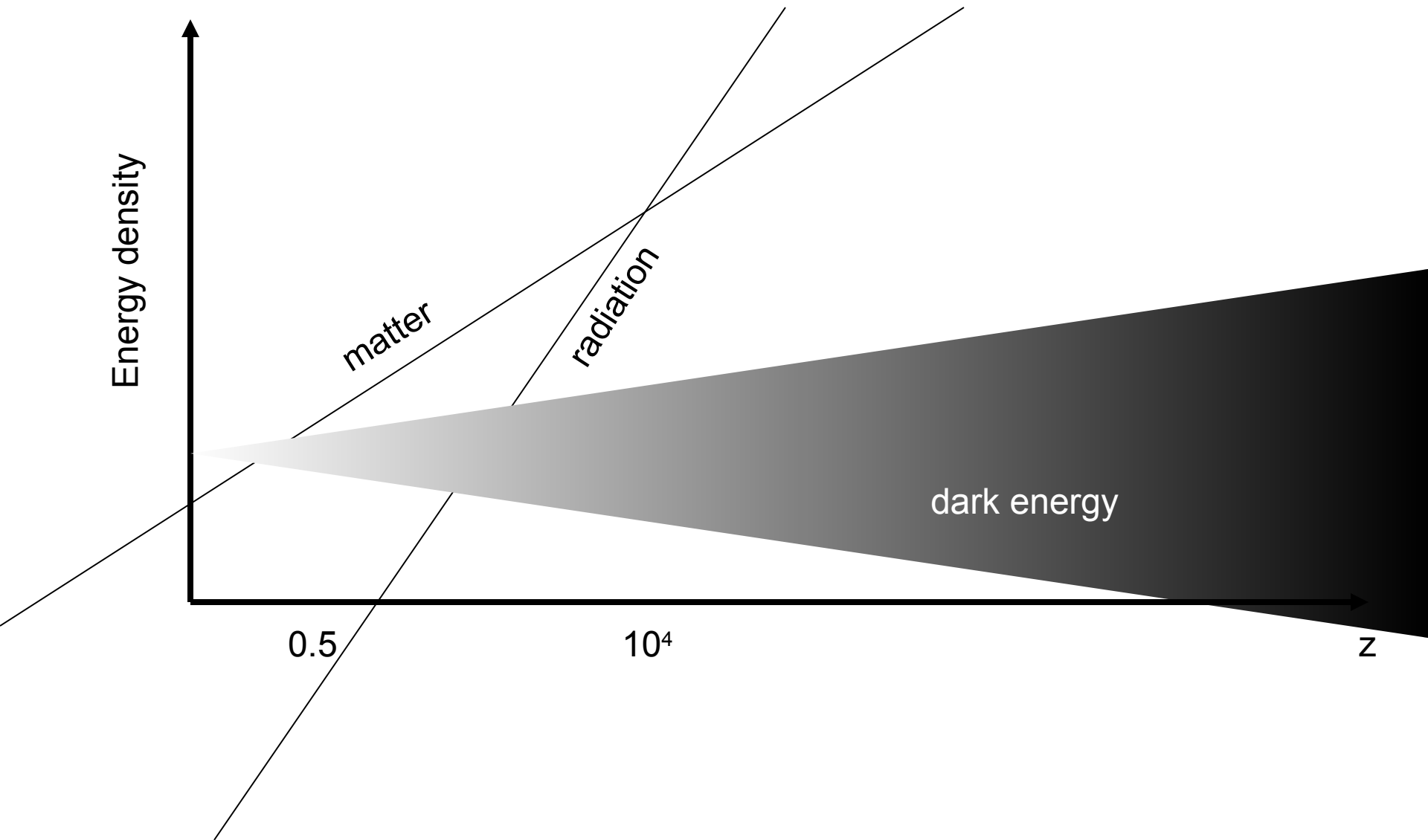
# The promise of CMB lensing



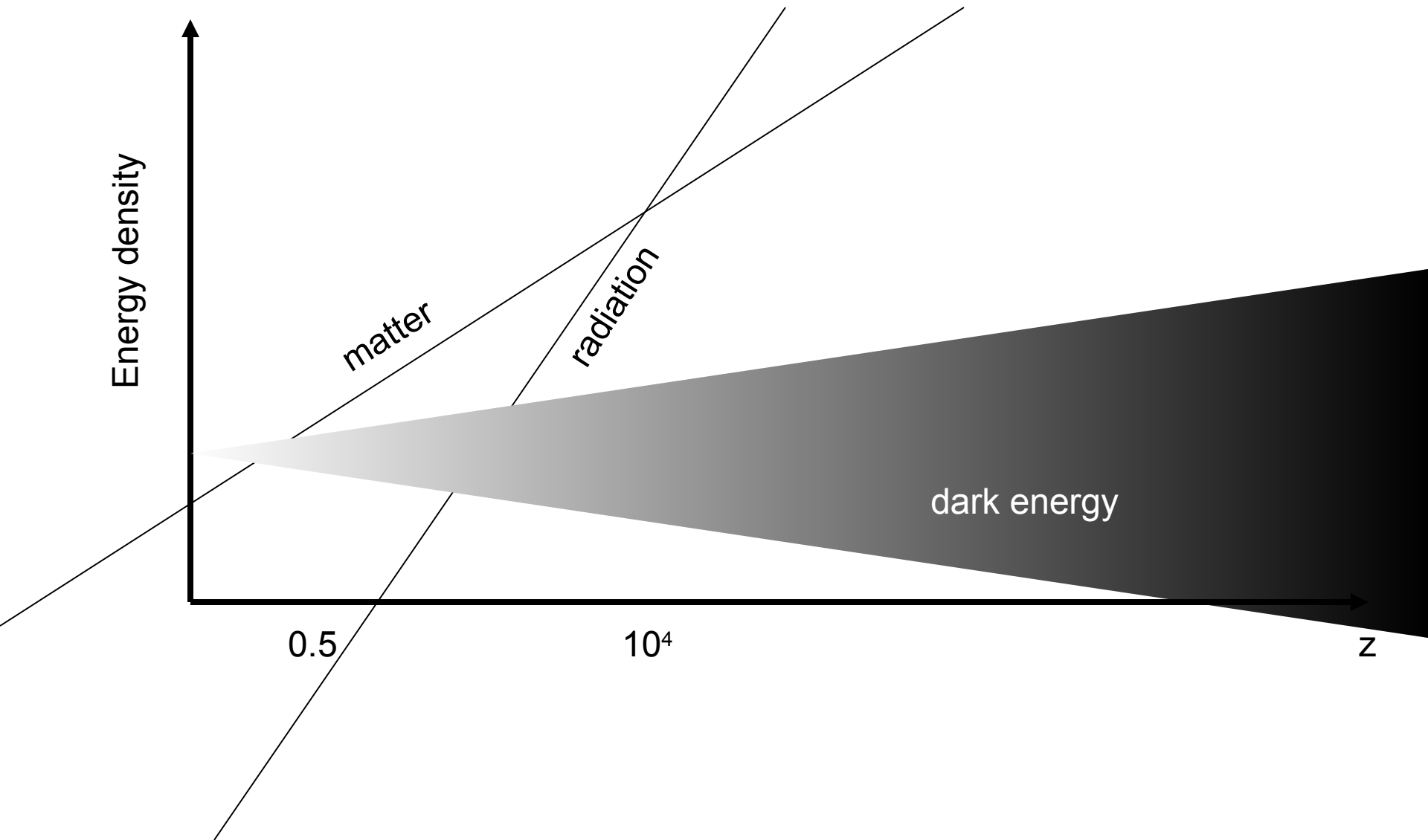
# The promise of CMB lensing



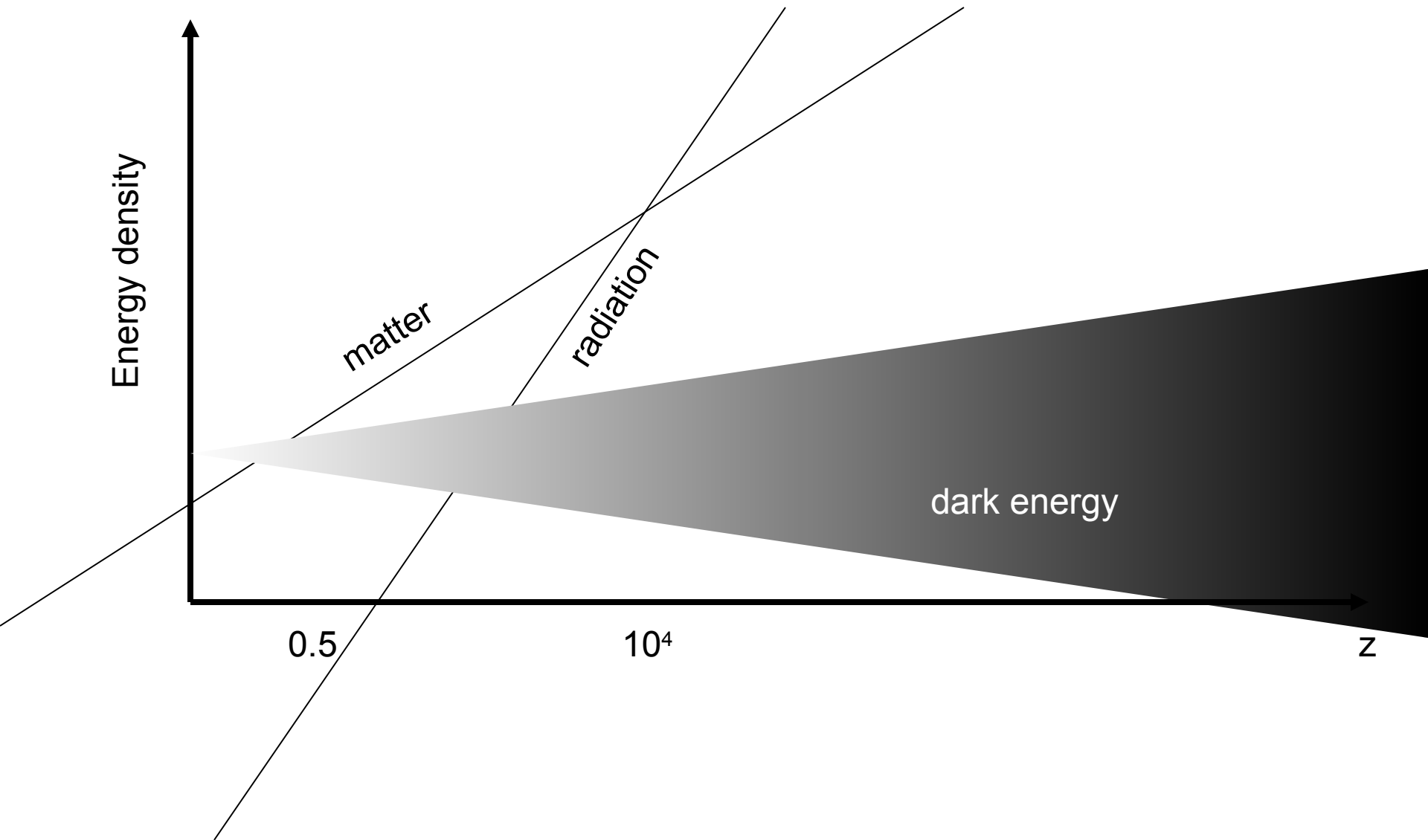
# The promise of CMB lensing



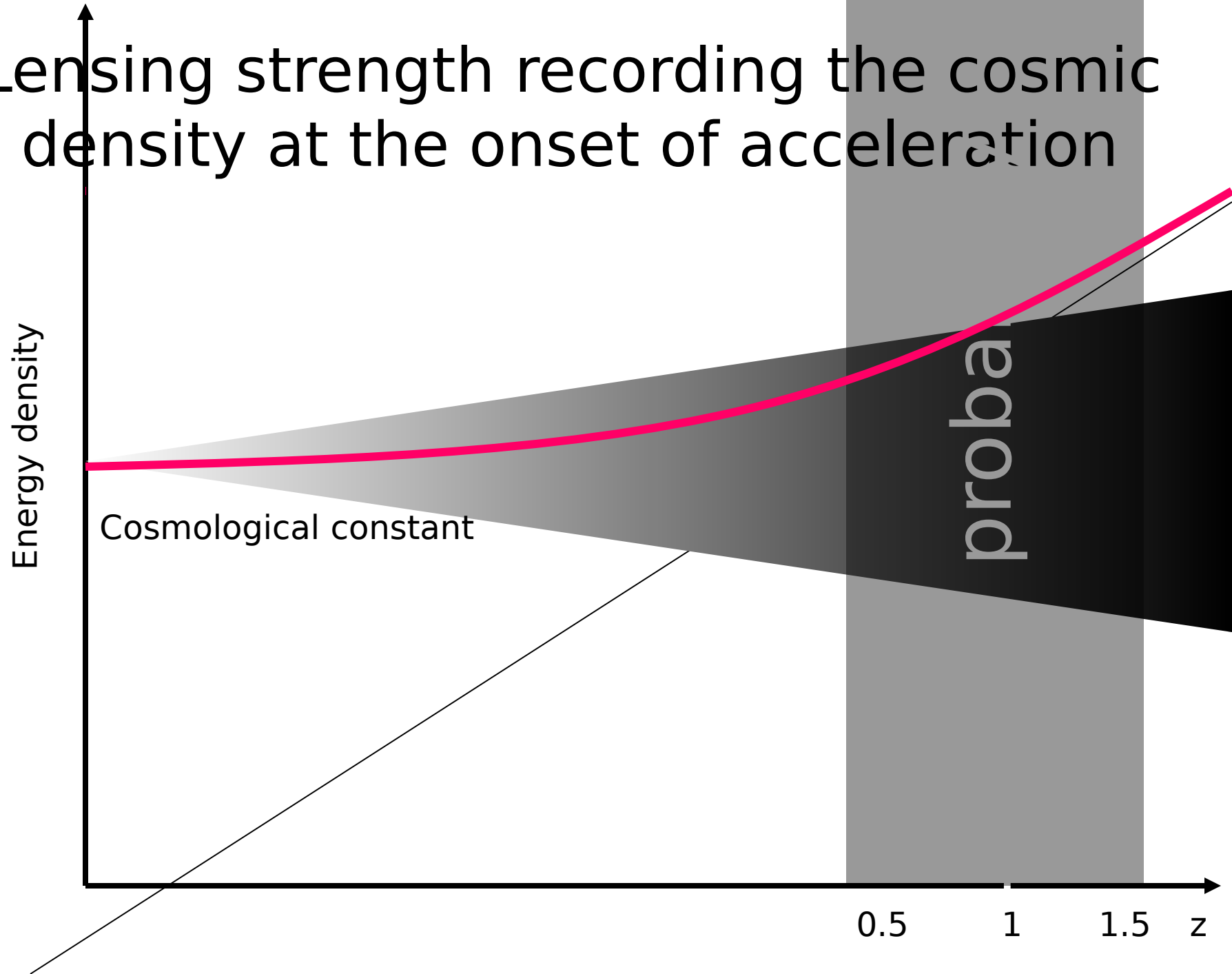
# The promise of CMB lensing



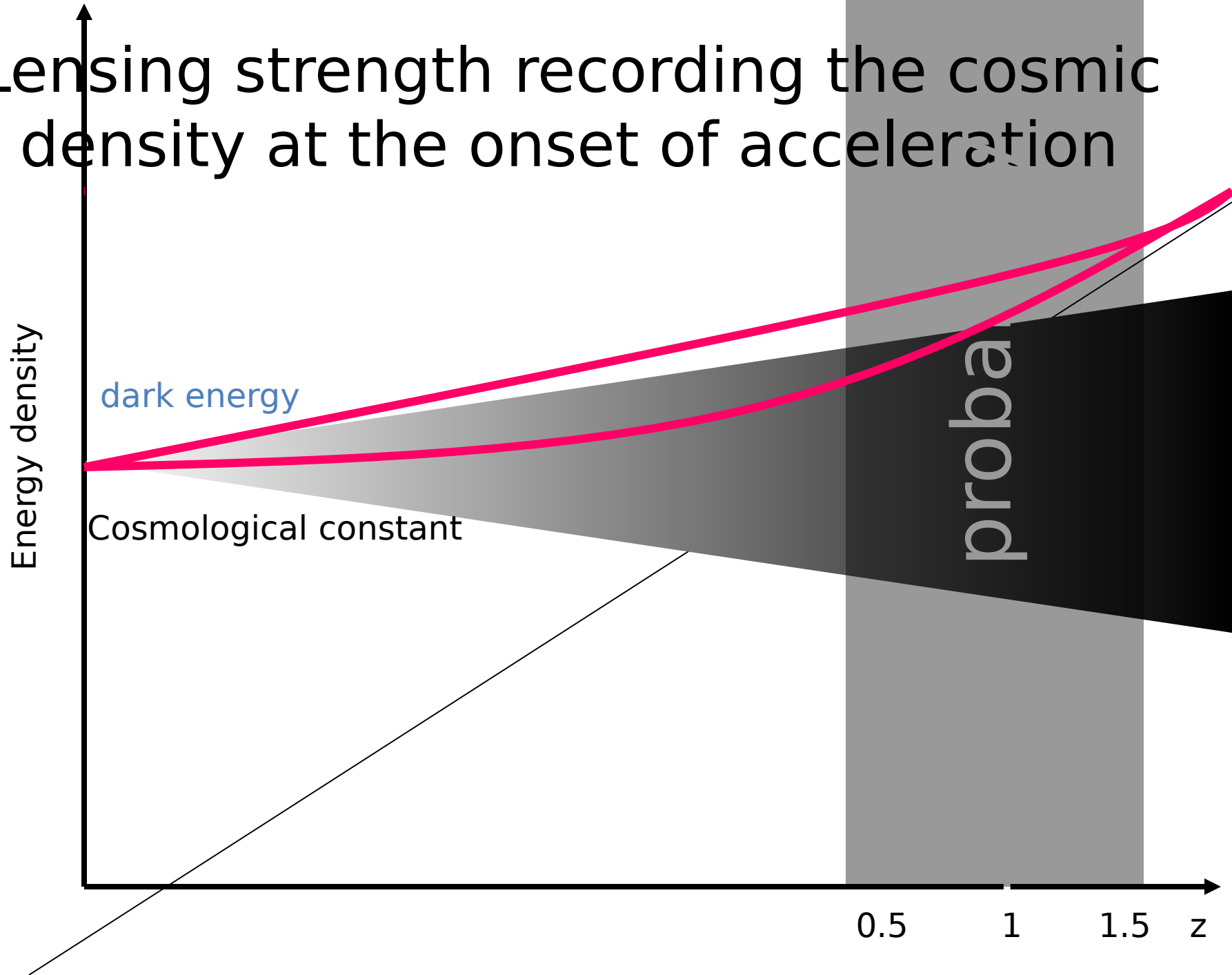
# The promise of CMB lensing



# Lensing strength recording the cosmic density at the onset of acceleration

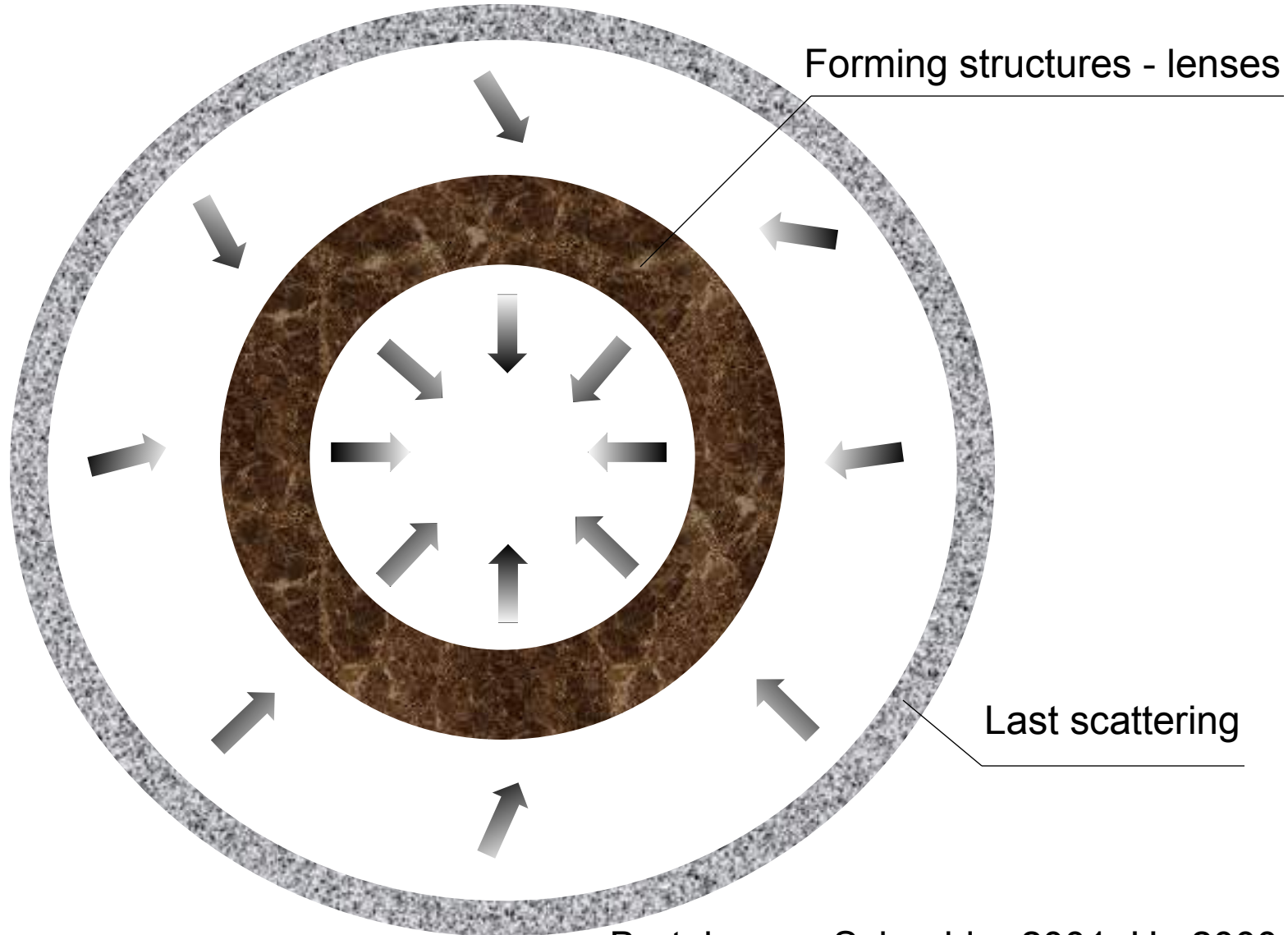


# Lensing strength recording the cosmic density at the onset of acceleration



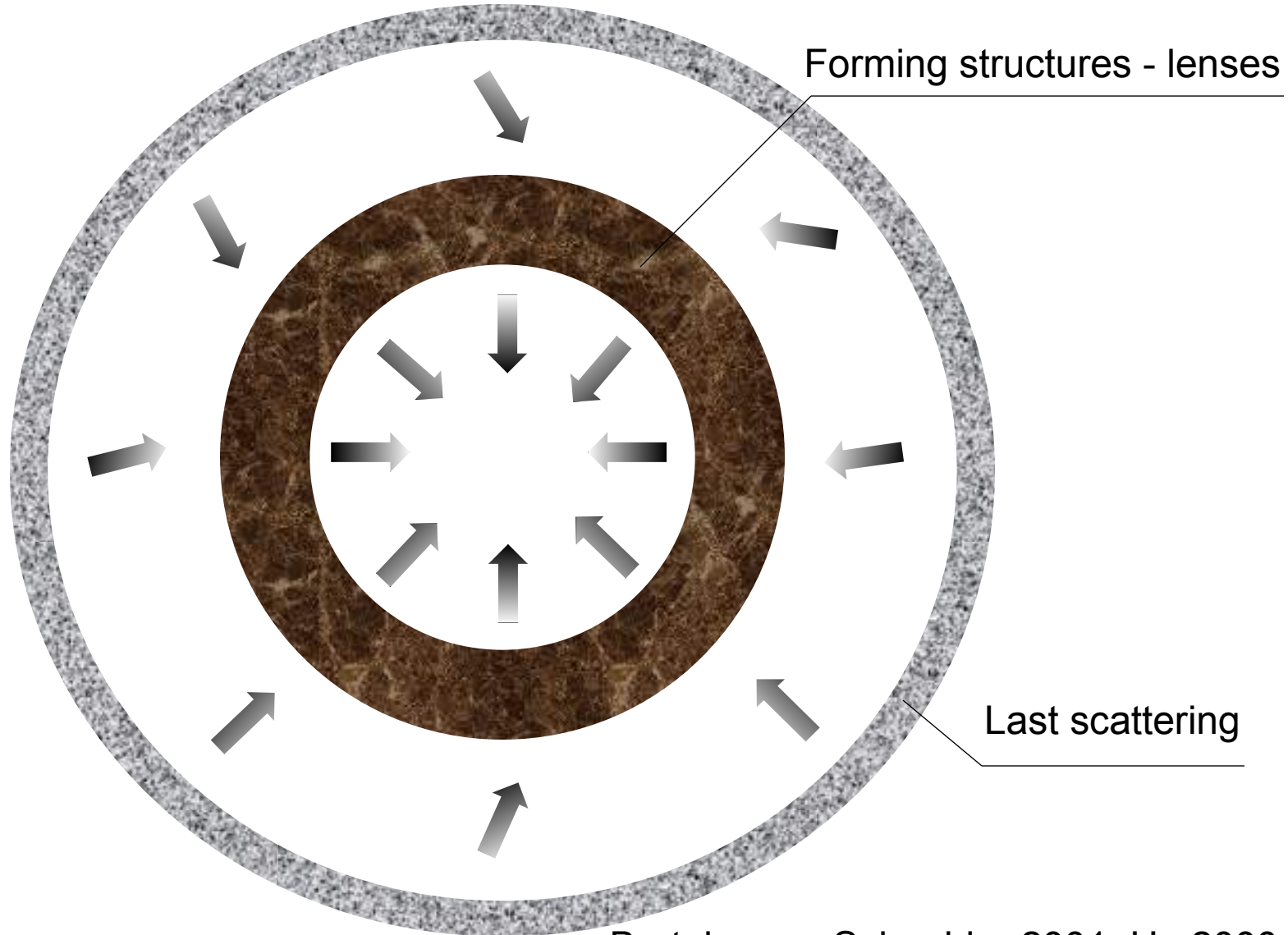
# CMB lensing

T



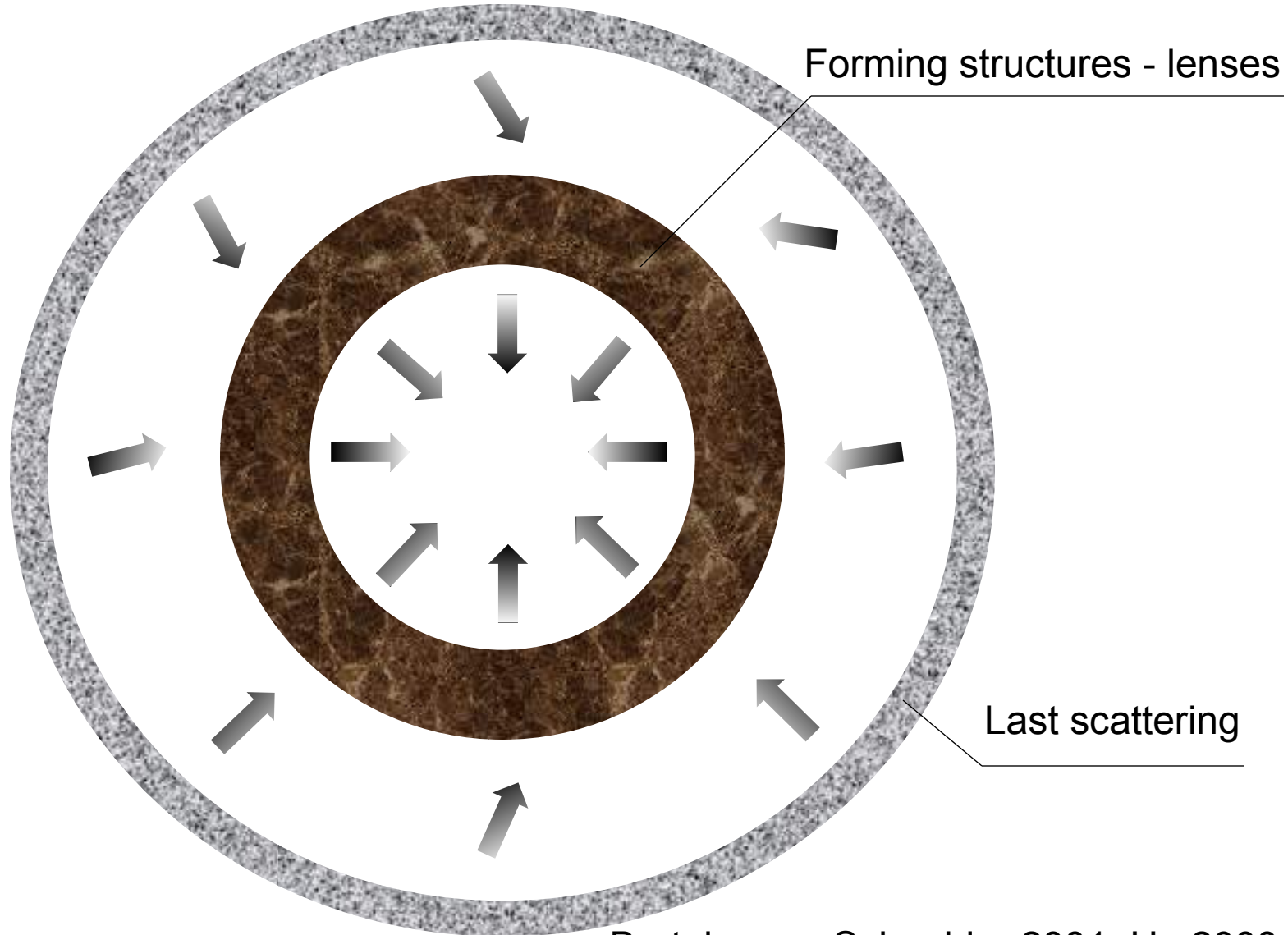
# CMB lensing

T



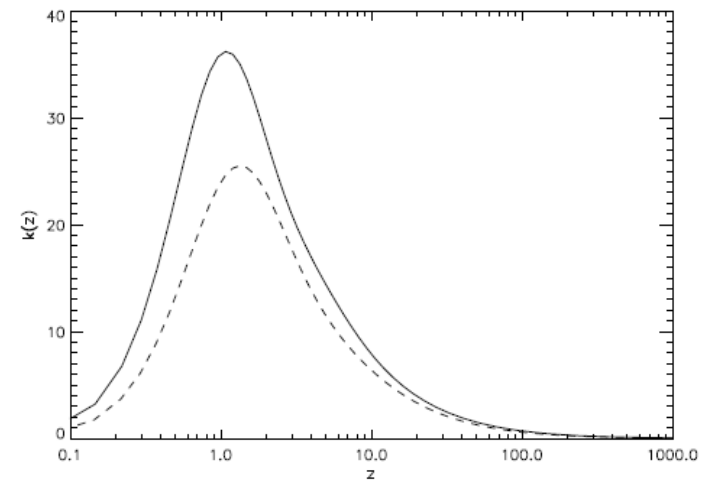
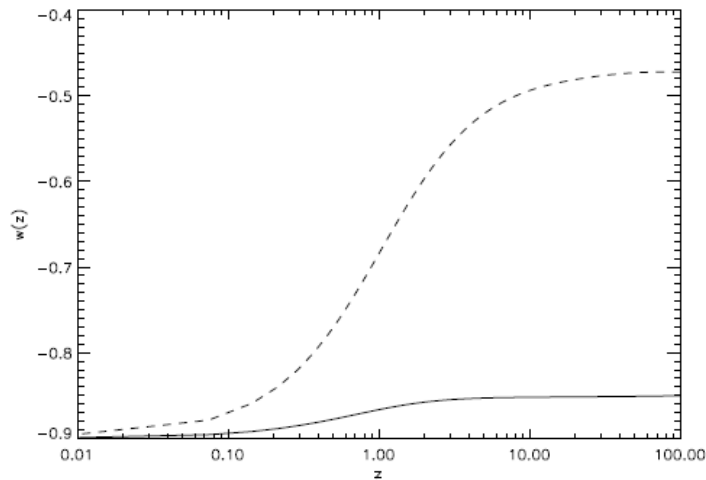
# CMB lensing

T

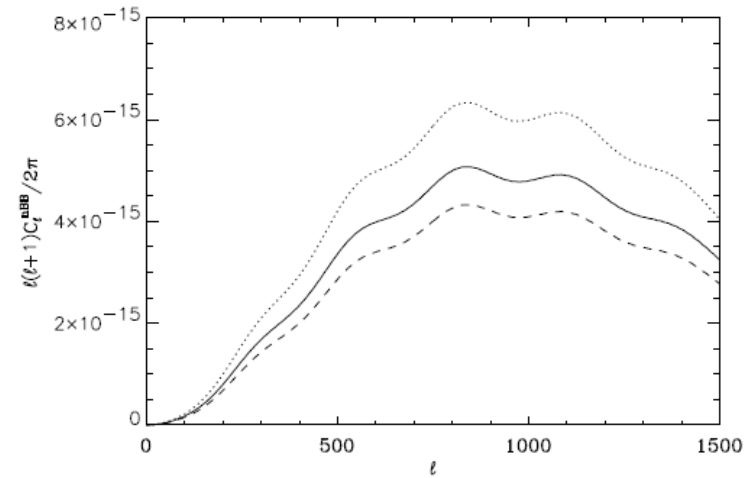
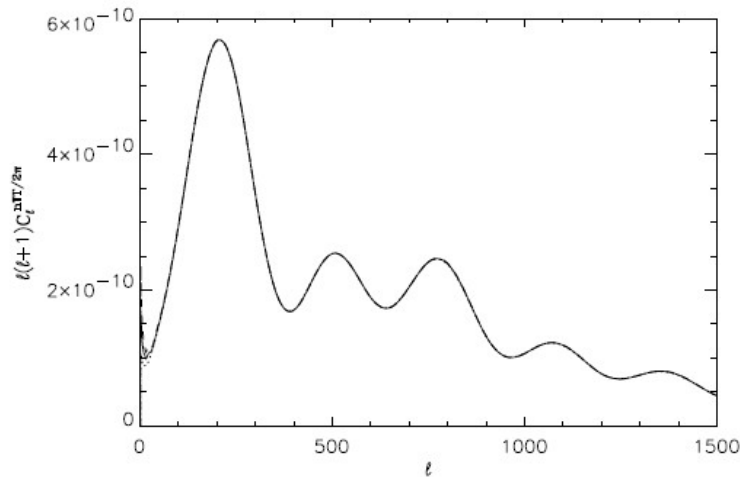


Bartelmann, Schneider 2001, Hu 2000

# Early dark energy and CMB lensing



# CMBreaking projection degeneracy



# Expanding universe => CMB

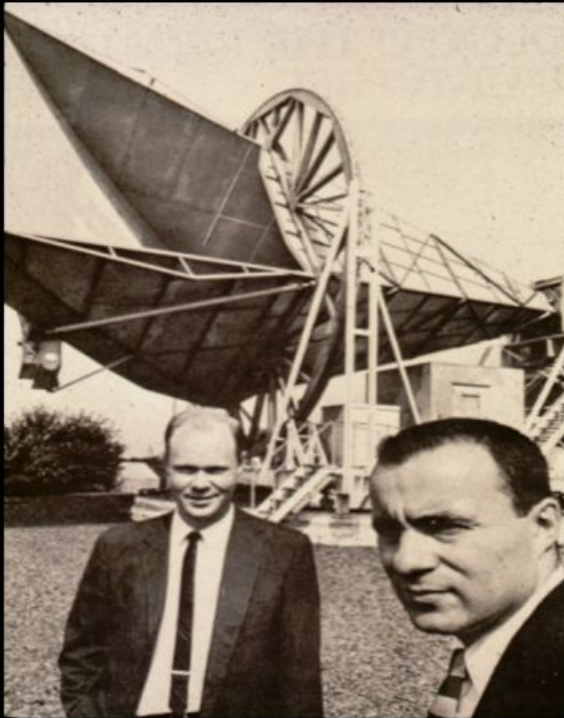
- compression in the early stages of an expanding universe causes lots of radiation arising from thermonuclear explosions
- Reactions are rapid enough to achieve thermalization and a black body spectrum
- It is possible to compute the rarefaction caused by the expansion since that epoch
- The relic radiation is predicted to peak in microwaves, temperature of a few Kelvin, known today as the Cosmic Microwave Background (CMB, Gamow et al. 1948)



George Gamow, three years old in Odessa, Ukraine, 1907

# Discovery

## Arno Penzias and Robert Wilson



Early 1960s - Penzias and Wilson are hired by Bell Labs to evaluate the performance of the new radio telescope to be used in trans-Atlantic telephone communications.

They find a small, unexplained signal regardless of the direction the telescope is pointed. It is not enough to be a problem, but they are curious.

1964 - They become aware that the noise in their telescope is the cosmic background radiation predicted by the Big Bang theory.

# COsmic Background Explorer

